

Binary black hole coalescence: Analysis of the plunge

Sourabh Nampalliwar^{1,2}

with Richard Price^{2,3} and Gaurav Khanna⁴

¹Fudan University, Shanghai

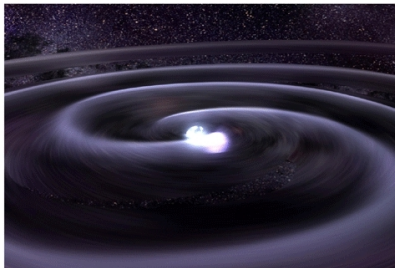
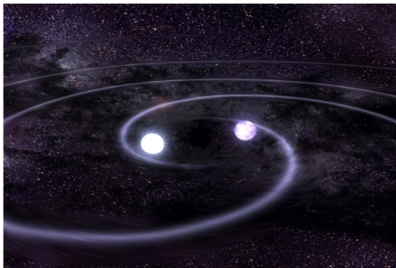
²U. Texas, Brownsville

³Massachusetts Institute of Technology

⁴U. Mass., Dartmouth

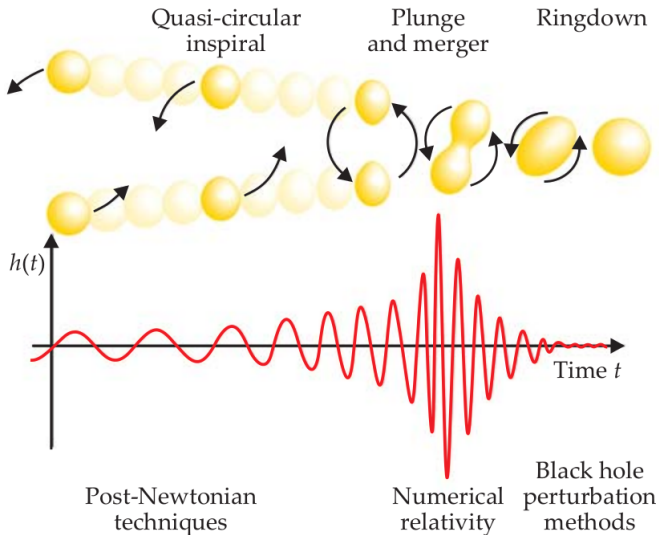
August 30, 2016

Gravitational wave sources



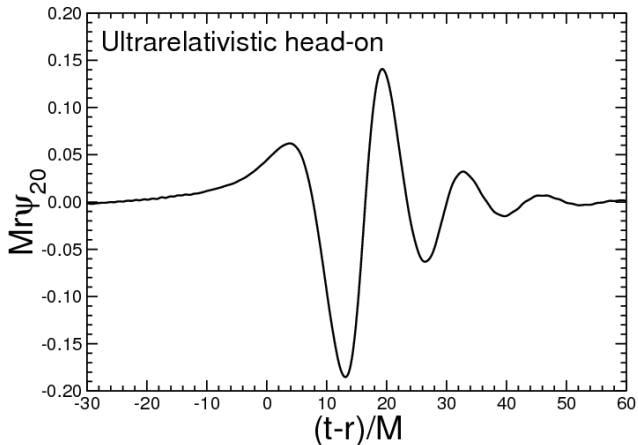
via LIGO

Binary black hole coalescence - the standard picture



via T. W. Baumgarte, S. L. Shapiro, Numerical Relativity, Cambridge U. Press, New York (2010)

Quasinormal ringing during the ringdown



via Berti et al., CQG 26 (2009) 163001



Binary black hole coalescence

- Goal: Understand plunge, in particular the excitation of QNR.
- Approach: Get insights, test insights.
- Principle: Simplify.

Complex frequencies \rightarrow FDGF \rightarrow Linearized Einstein equations

Complex frequencies \rightarrow FDGF \rightarrow Linearized Einstein equations

$$\Phi = \frac{1}{r} \sum_{l,m} \Psi_{lm}(r, t) Y_{lm}(\theta, \phi) \quad (1)$$

Complex frequencies \rightarrow FDGF \rightarrow Linearized Einstein equations

$$\Phi = \frac{1}{r} \sum_{l,m} \Psi_{lm}(r, t) Y_{lm}(\theta, \phi) \quad (1)$$

$$-\frac{\partial^2 \Psi_{lm}}{\partial t^2} + \frac{\partial^2 \Psi_{lm}}{\partial x^2} - V_\ell(x) \Psi_{lm} = S_{lm}(x, t) \quad (2)$$

$$r \rightarrow r^* \rightarrow x \quad (3)$$

Complex frequencies \rightarrow FDGF \rightarrow Linearized Einstein equations

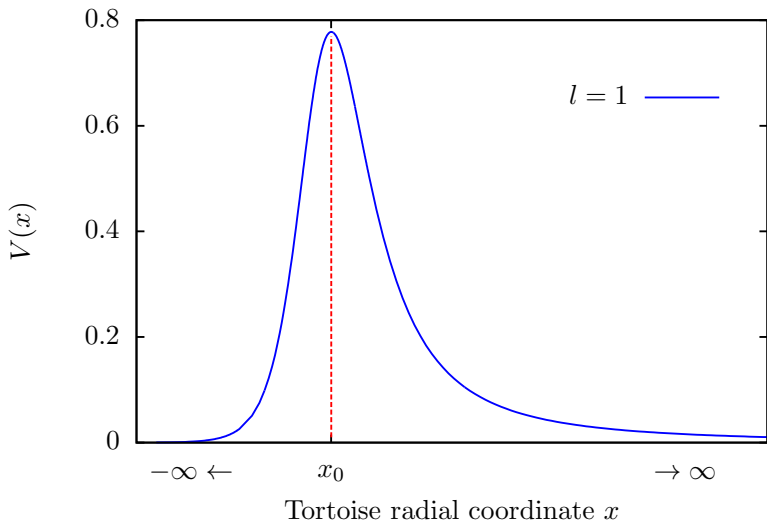
$$\Phi = \frac{1}{r} \sum_{l,m} \Psi_{lm}(r, t) Y_{lm}(\theta, \phi) \quad (1)$$

$$-\frac{\partial^2 \Psi_{lm}}{\partial t^2} + \frac{\partial^2 \Psi_{lm}}{\partial x^2} - V_\ell(x) \Psi_{lm} = S_{lm}(x, t) \quad (2)$$

$$r \rightarrow r^* \rightarrow x \quad (3)$$

$$\Psi(t, x) = \text{Integral over Green function} \quad (4)$$

Comparison - TDP vs Schwarzschild

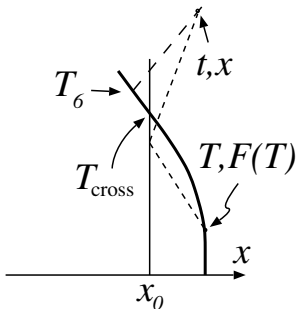


$$\begin{aligned}\Psi(t-x) = & \frac{x_0}{2} \int_{T_{\text{cross}}}^{T_6} e^{-\omega_d \theta} \left[\frac{1}{x_0} (\sin \omega_o \theta - \cos \omega_o \theta) - \frac{2}{x} \sin \omega_o \theta \right] dT \\ & - \frac{1}{2} \int_{-\infty}^{T_{\text{cross}}} e^{-\omega_d \xi} \left[\left(\frac{2x_0}{F(T)} - 1 \right) \cos \omega_o \xi - \sin \omega_o \xi \right. \\ & \left. - \frac{2x_0^2}{x F(T)} (\cos \omega_o \xi - \sin \omega_o \xi) \right] dT. \quad (5)\end{aligned}$$

where

$$\theta = u - T + F(T), \quad \xi = u - T - F(T) + 2x_0.$$

Analysis - QNR from radial infall in TDP



$$\begin{aligned}\Psi(u) = & \frac{x_0}{2} \int_{T_{\text{cross}}}^{T_6} e^{-\omega_d \theta} \left[\frac{1}{x_0} (\sin \omega_o \theta - \cos \omega_o \theta) - \frac{2}{x} \sin \omega_o \theta \right] dT \\ & - \frac{1}{2} \int_{-\infty}^{T_{\text{cross}}} e^{-\omega_d \xi} \left[\left(\frac{2x_0}{F(T)} - 1 \right) \cos \omega_o \xi - \sin \omega_o \xi \right. \\ & \left. - \frac{2x_0^2}{xF(T)} (\cos \omega_o \xi - \sin \omega_o \xi) \right] dT .\end{aligned}\quad (6)$$

where

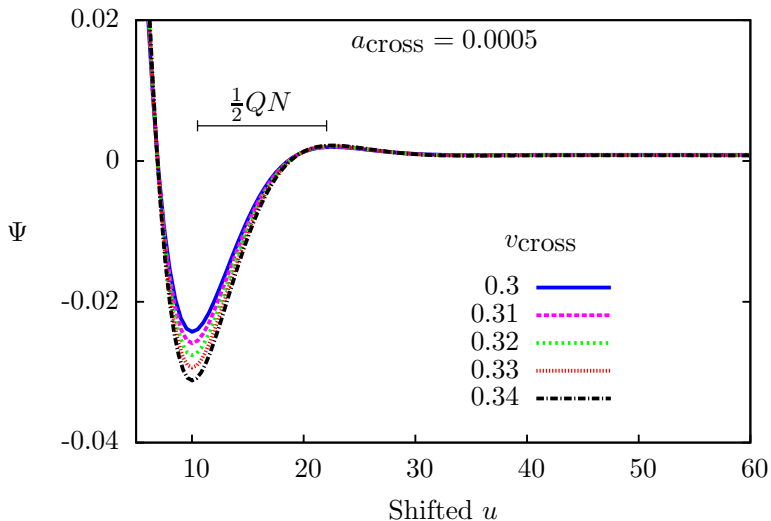
$$\theta = u - T + F(T), \quad \xi = u - T - F(T) + 2x_0$$

$$\begin{aligned}\Psi(u) = & \frac{x_0}{2} \int_{T_{\text{cross}}}^{T_6} e^{-\omega_d \theta} \left[\frac{1}{x_0} (\sin \omega_o \theta - \cos \omega_o \theta) - \frac{2}{x} \sin \omega_o \theta \right] dT \\ & - \frac{1}{2} \int_{-\infty}^{T_{\text{cross}}} e^{-\omega_d \xi} \left[\left(\frac{2x_0}{F(T)} - 1 \right) \cos \omega_o \xi - \sin \omega_o \xi \right. \\ & \quad \left. - \frac{2x_0^2}{xF(T)} (\cos \omega_o \xi - \sin \omega_o \xi) \right] dT.\end{aligned}\quad (7)$$

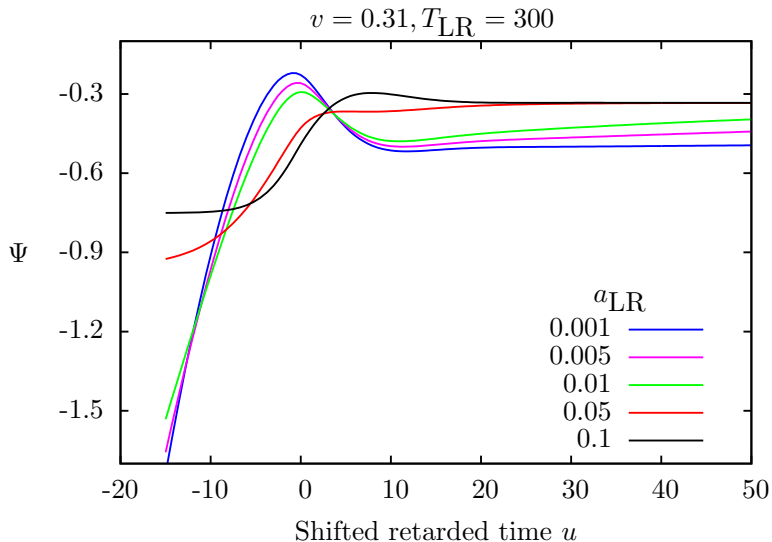
where

$$\theta = u - T + F(T), \quad \xi = u - T - F(T) + 2x_0$$

Analysis - QNR from radial infall in TDP

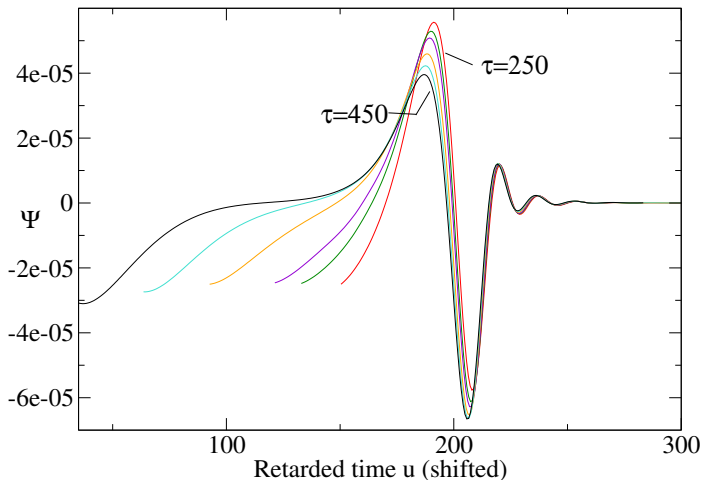


Analysis - QNR from radial infall in TDP



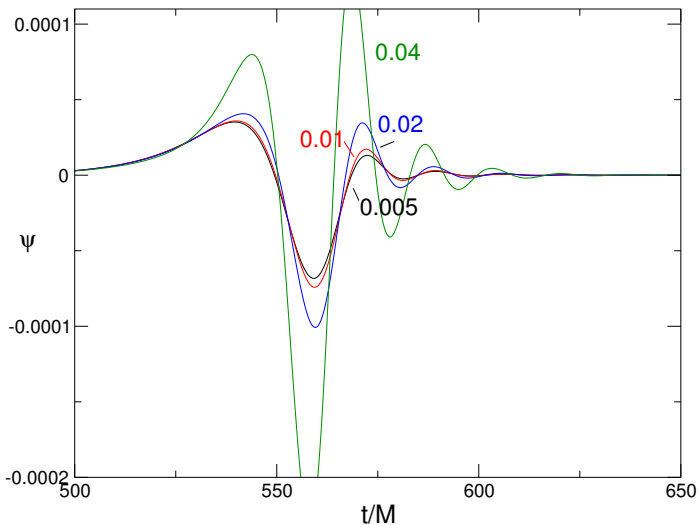
$$-\frac{\partial^2 \Psi_{\ell m}}{\partial t^2} + \frac{\partial^2 \Psi_{\ell m}}{\partial x^2} - V_\ell(x) \Psi_{\ell m} = S_{\ell m}(x, t) \quad (8)$$

Analysis - QNR from radial infall in Schwarzschild



Data courtesy Gaurav Khanna

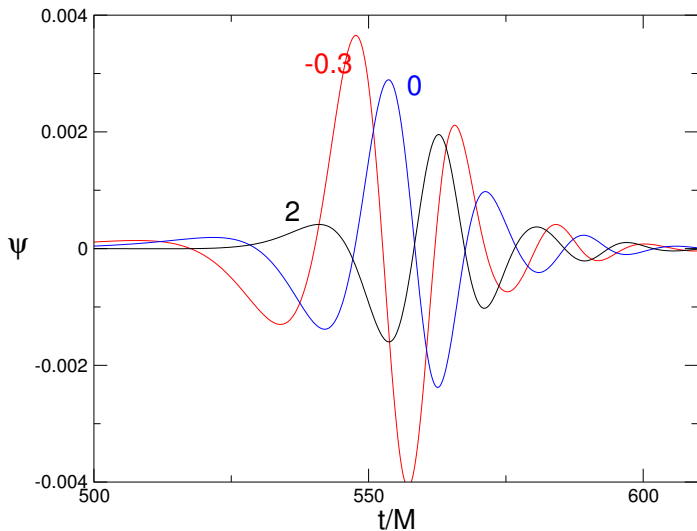
Analysis - QNR from orbital infall in Schwarzschild



Data courtesy Gaurav Khanna



Analysis - QNR from orbital infall in Schwarzschild



Data courtesy Gaurav Khanna



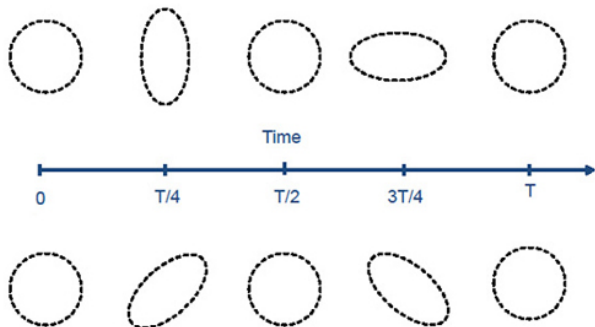
During a binary black hole coalescence, the transition from the inspiral phase to the ringdown phase is completely determined by the local properties at the photon orbit.

Thank you!

BONUS!

Gravitational waves

- Travel at the speed of light.
- Carry energy.

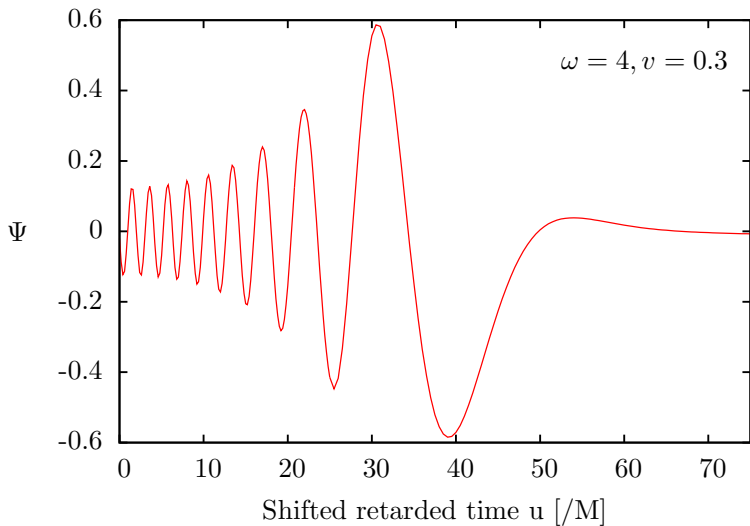


via learner.org

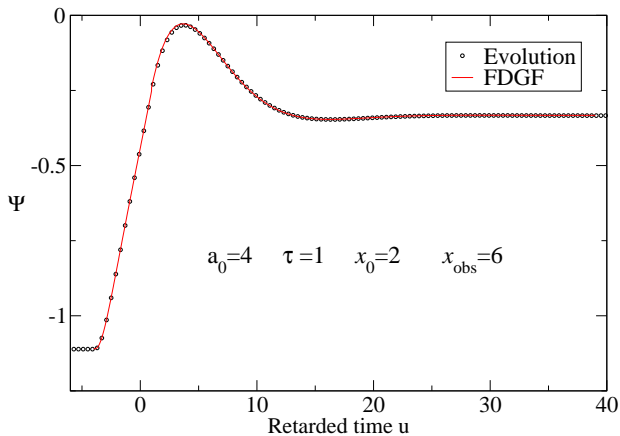
Quasinormal ringing



Analysis - QNR from orbital infall in TDP



Comparison - Evolution vs FDGF



Evolution data courtesy Gaurav Khanna

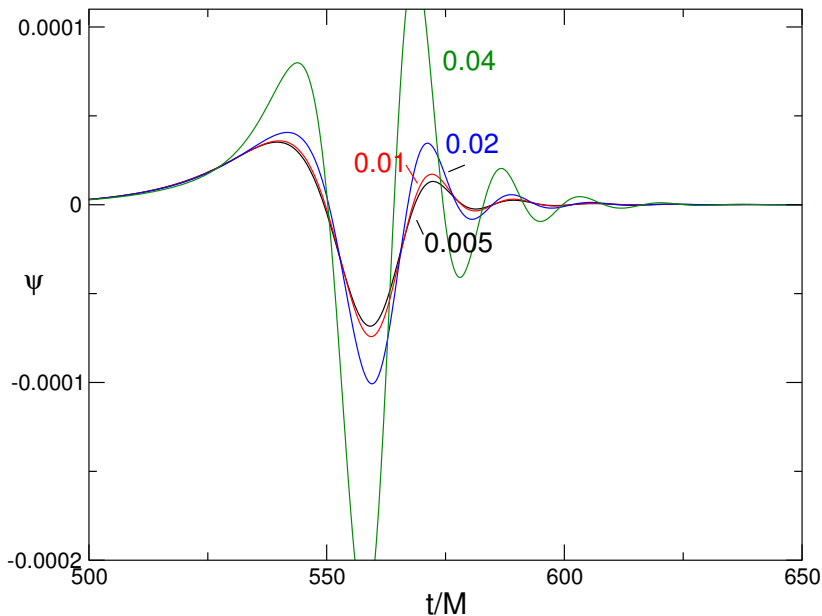


Assumption - Trajectories: Orbital

- Redshift in angular velocity.
- $d\omega/dt = 0$ at the light ring.
- Speed of light should not be breached.

$$\omega(T) = \omega_{\text{LR}} \frac{27M^2}{(1 + \sigma)r(T)^2} \left(1 - \frac{2M}{r(T)}\right) \left(1 + \frac{3\sigma M}{r(T)}\right) \quad (9)$$

Analysis - QNR from orbital motion in Schwarzschild



Interesting features of orbital trajectories

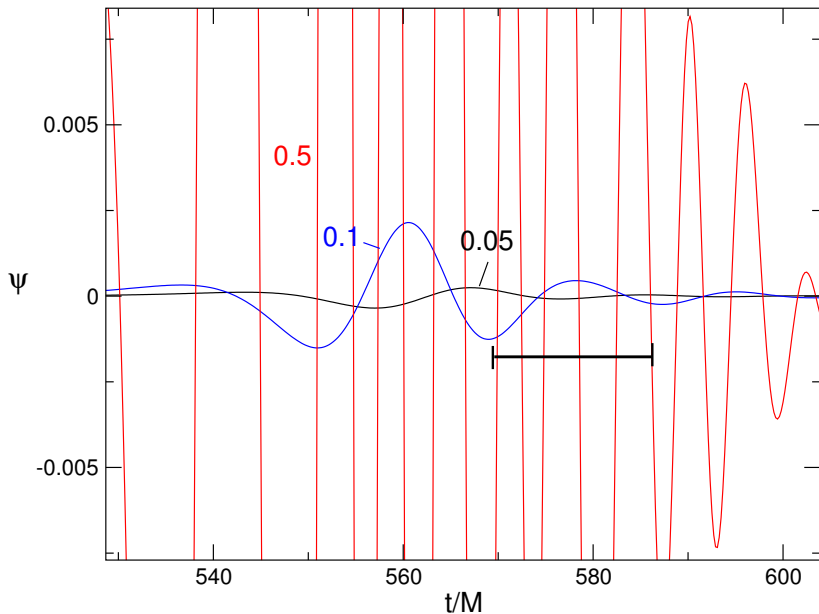
- Speed of light constraint:

$$\omega_{\text{LR}}^2 M^2 < (1 - v_{\text{LR}}^2)/27. \quad (10)$$

- Here this implies

$$\omega_{\text{LR}} < 0.1836/M. \quad (11)$$

Direct radiation



$$\frac{\partial^2 G}{\partial x^2} - \frac{\partial^2 G}{\partial t^2} - V(x)G = \delta(x - a)\delta(t - T) \quad (12)$$

$$\Psi(t, x) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} G(x, a; t - T) f(T) \delta(a - F[T]) da dT \quad (13)$$

$$G(x, a; t - T) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-i\omega(t-T)} \mathcal{G}(x, a; \omega) d\omega, \quad (14)$$

$$\Psi(t, x) = \frac{1}{2\pi} \int \int \int_{-\infty}^{\infty} e^{-i\omega(t-T)} \mathcal{G}(x, a; \omega) \delta(a - F[T]) d\omega da dT. \quad (15)$$

Assumption - Trajectories: Radial

- $dx/dt \rightarrow -1$ near the horizon.
- At least two parameters.
- Static initial data.

$$F(T) = \begin{cases} a_0 + \tau - (T^3 + \tau^3)^{1/3} & T \geq 0 \\ a_0 & T < 0. \end{cases} \quad (16)$$

Approximation - Truncated Multipole Potential

- Features of *curvature potential**

$$V_\ell = \begin{cases} \frac{\ell(\ell+1)}{x^2} & \text{for } r \gg 2M \\ (1 - \frac{2M}{r}) & \text{for } r \sim 2M \end{cases} \quad (17)$$

*Cunningham, Price & Moncrief, ApJ, 224, 643-667 (1978)

Approximation - Truncated Multipole Potential

- Features of *curvature potential**

$$V_\ell = \begin{cases} \frac{\ell(\ell+1)}{x^2} & \text{for } r \gg 2M \\ (1 - \frac{2M}{r}) & \text{for } r \sim 2M \end{cases} \quad (17)$$

- Truncated Dipole Potential (TDP)

$$V = \begin{cases} 2/x^2 & \text{for } x \geq x_0 \\ 0 & \text{for } x < x_0 \end{cases} \quad (18)$$

where x_0 is the light ring *aka* photon orbit *aka* UCO.

*Cunningham, Price & Moncrief, ApJ, 224, 643-667 (1978)