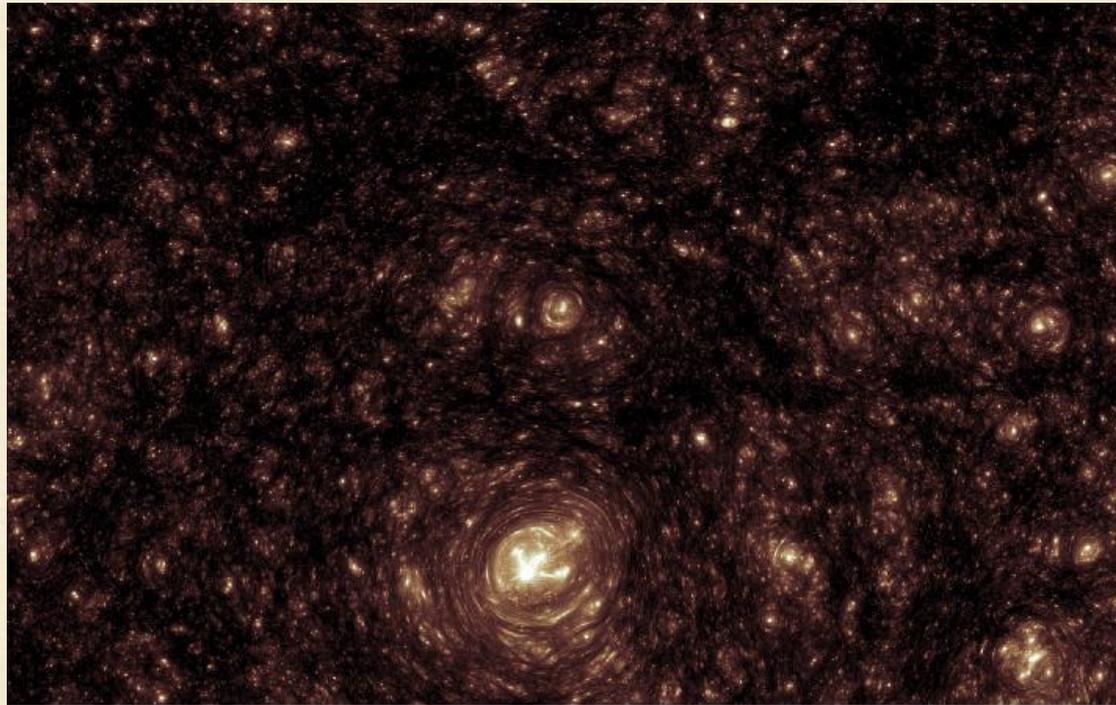


# Prediction of Gravitational lensing signal through Horizon-AGN light cone



C. Gouin, R. Gavazzi, Y. Dubois and C. Pichon

# 1. Introduction to Gravitational Lensing

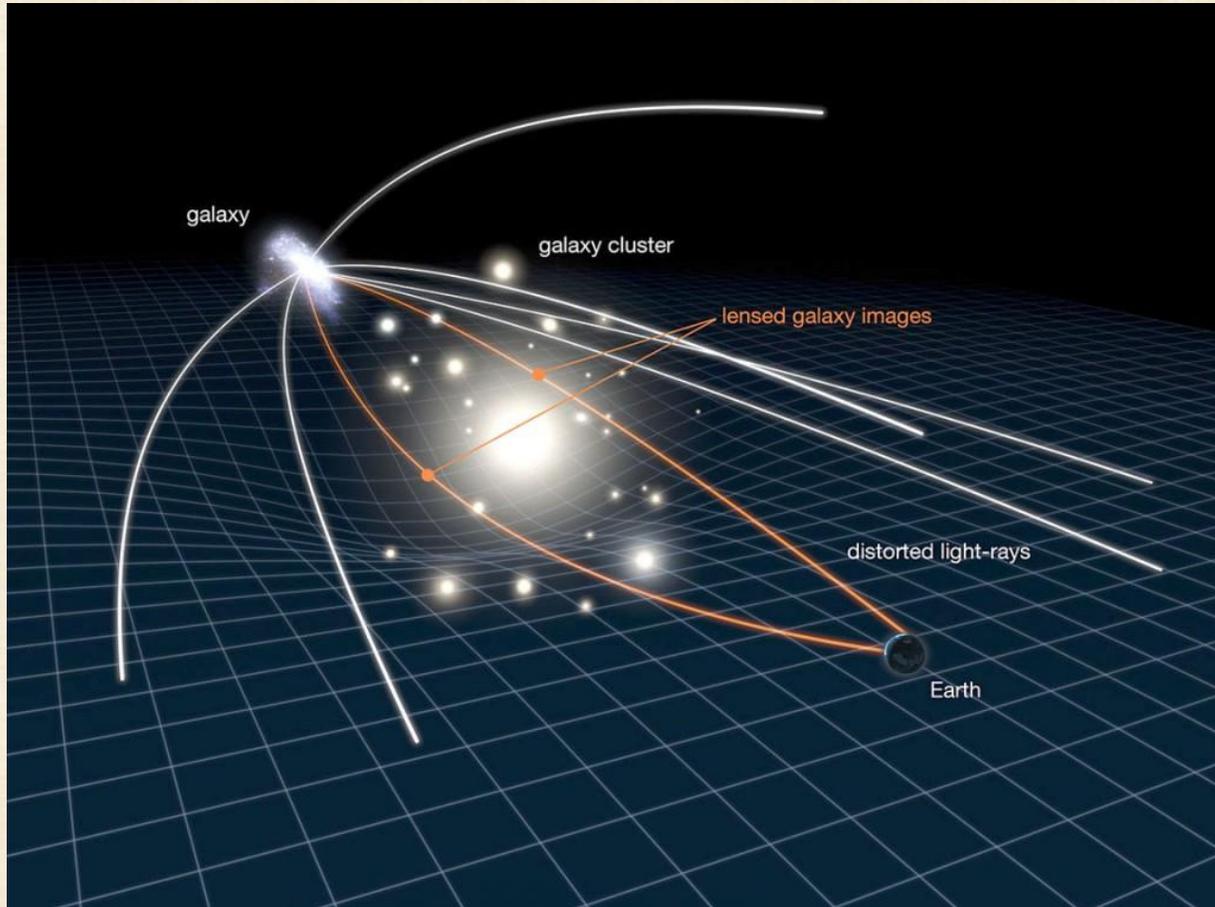
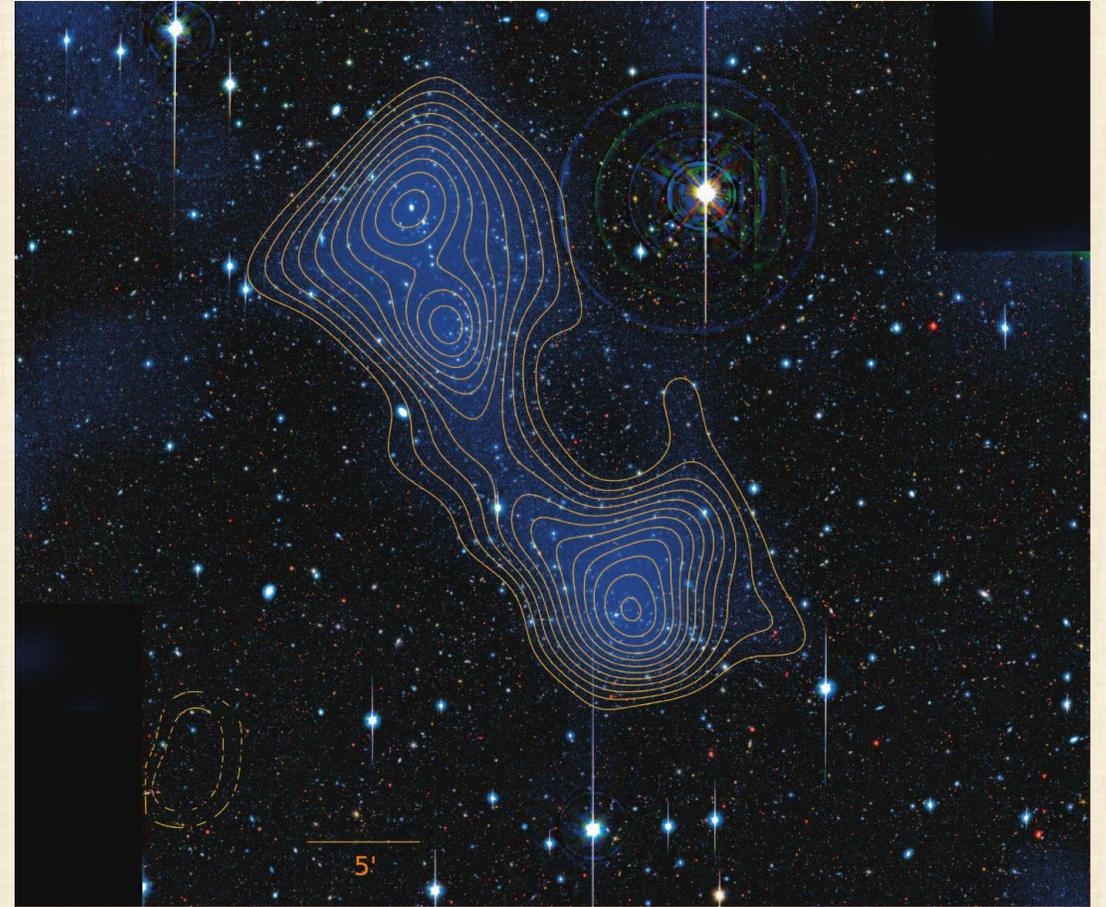
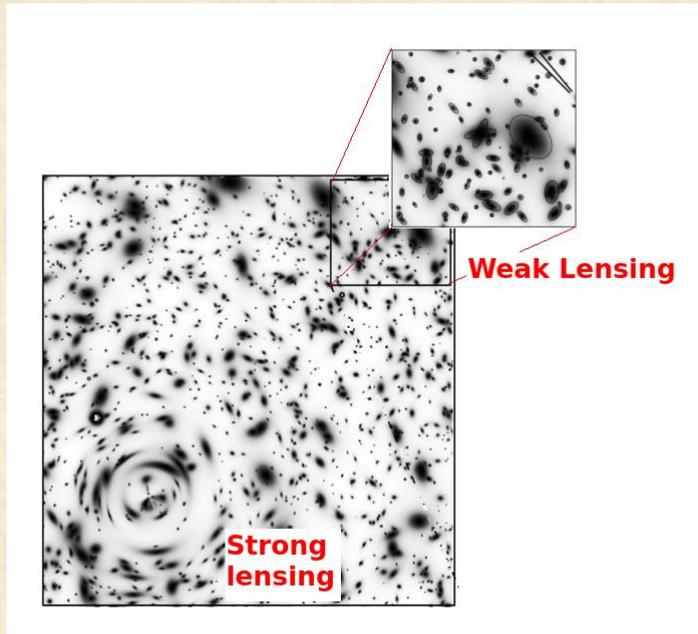


Image : NASA /ESA



Mass reconstruction of Abell 222/223 by WL  
Dietrich and al, 2012

# 1. Introduction to Gravitational Lensing



⇒ Strong Lensing probe the internal region of galaxies

⇒ Weak lensing signal probe the outskirts of galaxies

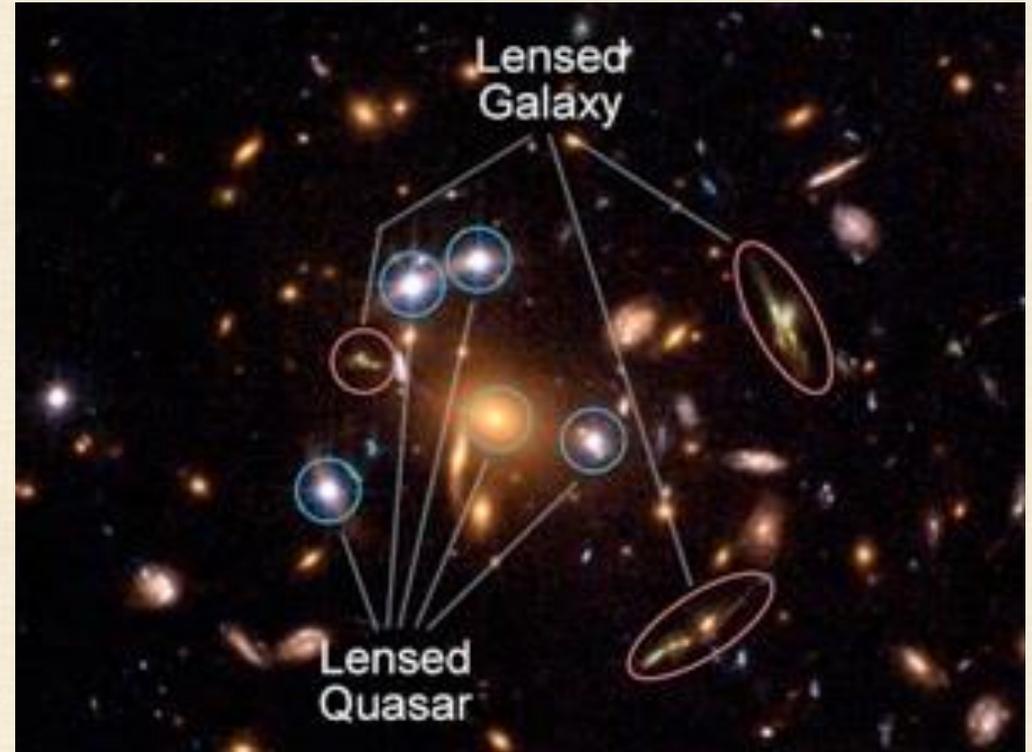


Image : NASA / ESA, K. Sharon (Tel Aviv) E. Ofek (Caltech)



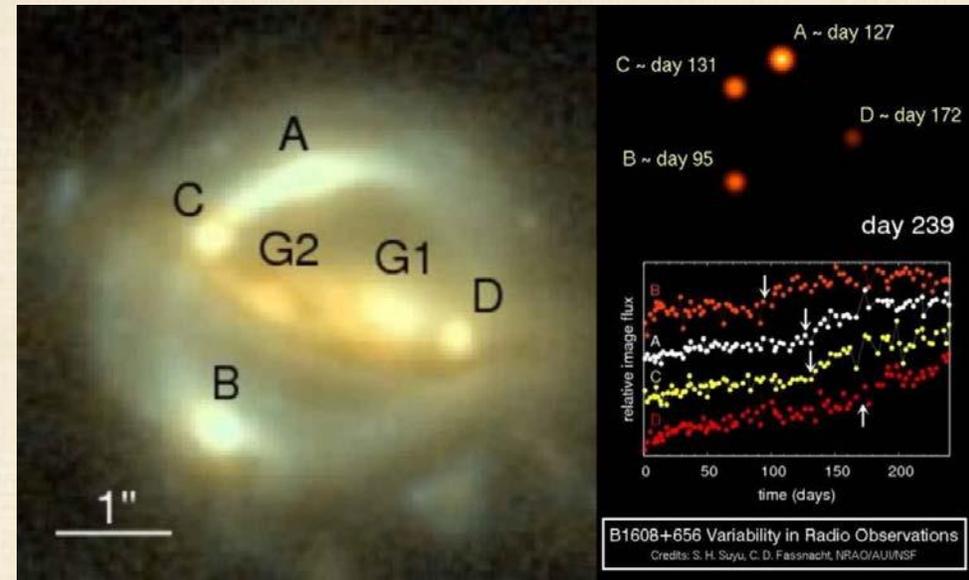
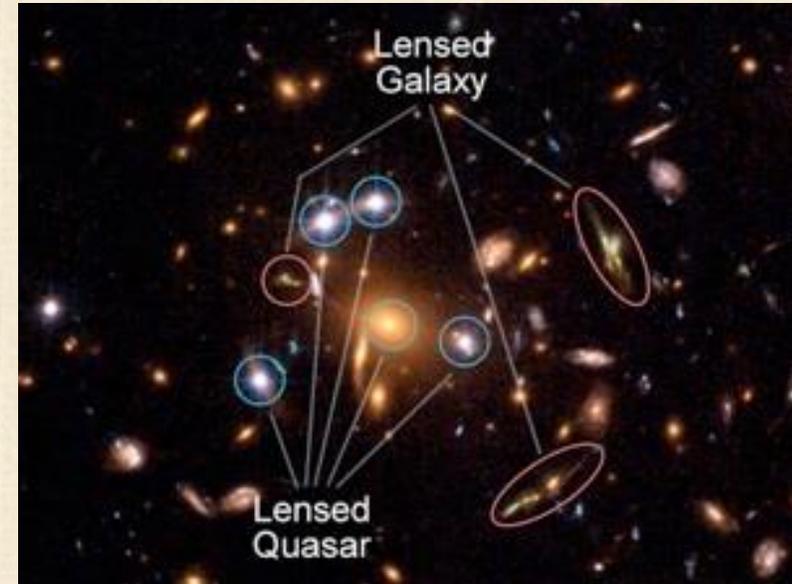
Euclid space mission

Calibration of GL signal using hydrodynamical cosmological simulation

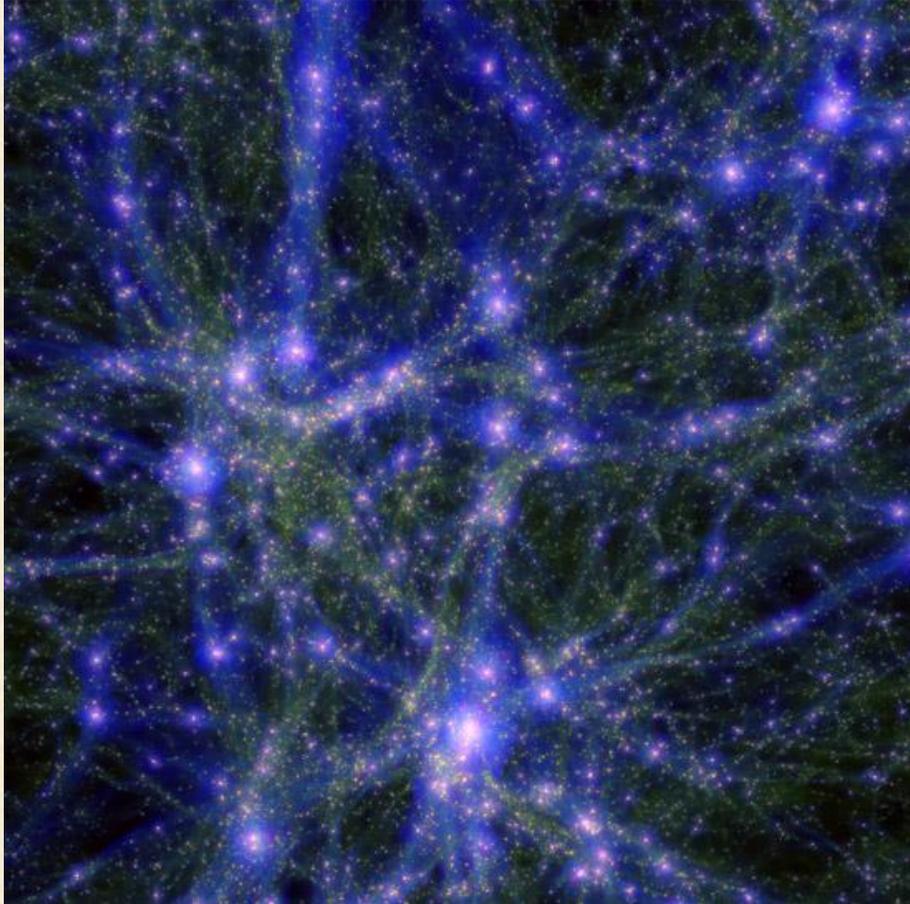
# 1. Prediction of GL signal from hydrodynamical cosmological simulation

## On small scales

- The slope of the central density profile
  - Cusp core problem?
- cross section of strong lensing
  - Is arc abundance changed by baryonic processes ?
- Time delay
  - Constraining the Hubble constant
  - The distribution of mass along the l.o.s ?



# 1. Prediction of GL signal from hydrodynamical cosmological simulation



Horizon AGN simulation

Baryonic processes modeled in Horizon AGN Simulation with RAMSES code :

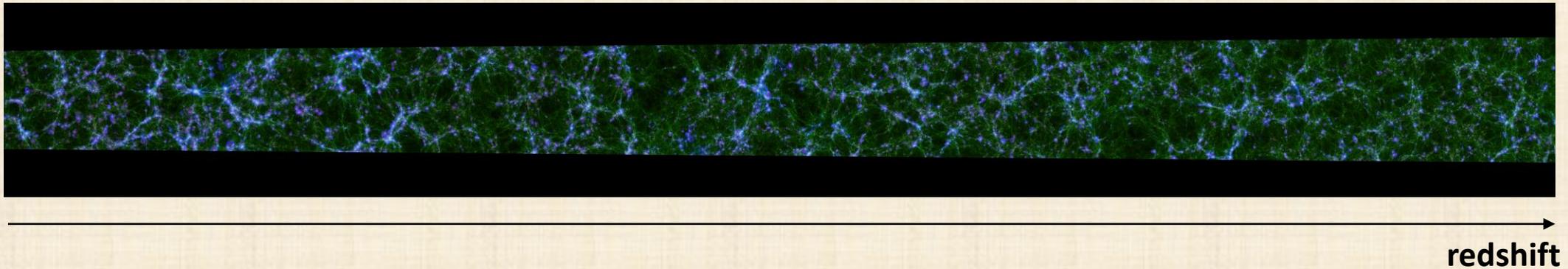
- Gas dynamics, gas cooling/heating
- Star formation
- SN & AGN feedback

Horizon AGN  $\sim$  100 000 galaxies within a box 100 Mpc/h

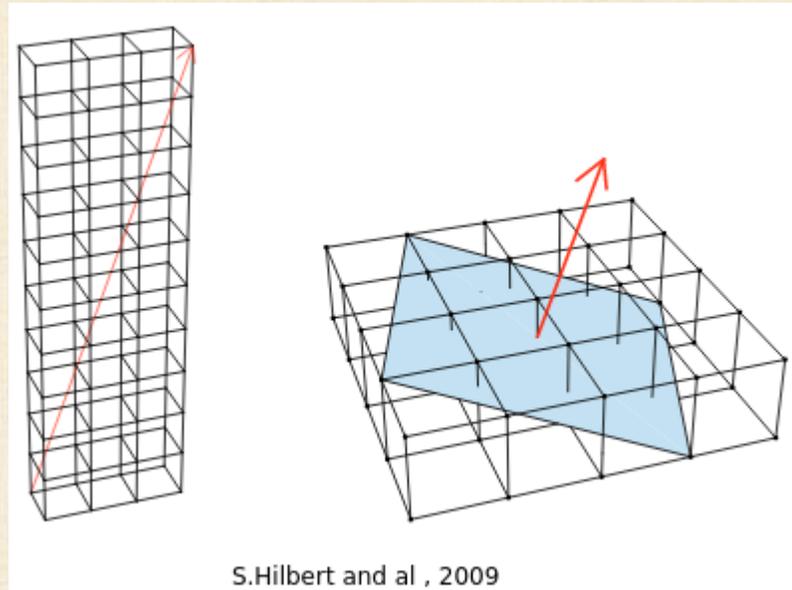
(See Presentation of Sugata Kaviraj)

## 2. Ray tracing through Horizon AGN light cone

Horizon AGN light cone



Goal : Tracing light ray through the light cone



Example of light cone creation in a box simulation

## 2. Ray tracing through Horizon AGN light cone

### Thin lens theory

Lensing equation :

$$\vec{\beta} = \vec{\theta} - \frac{D_{LS}}{D_s} \vec{\alpha}$$

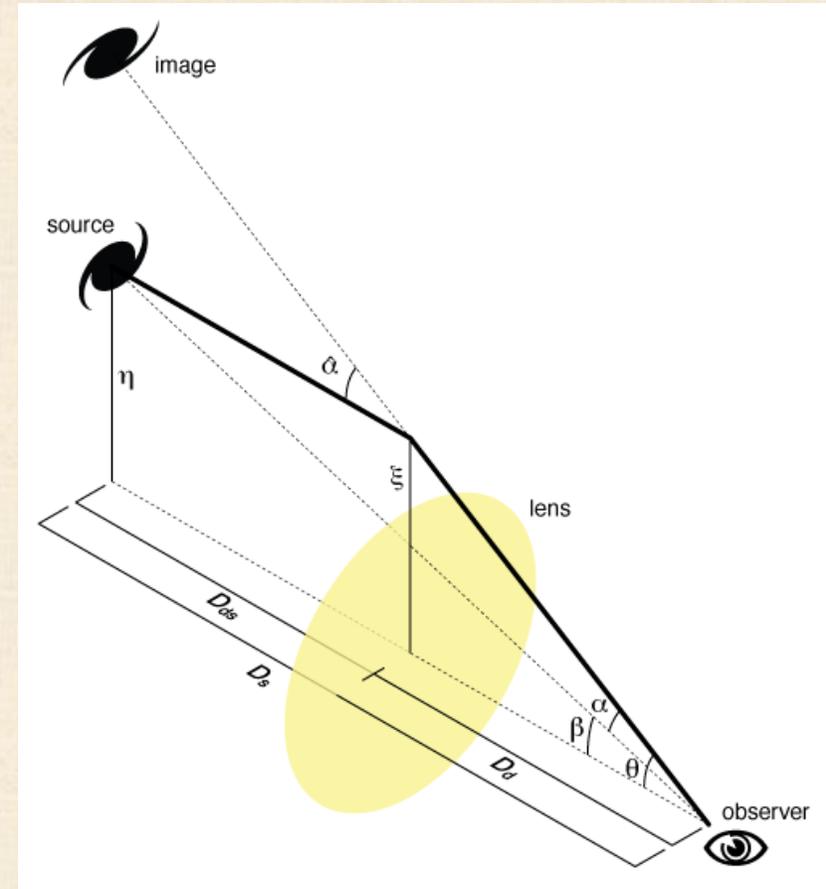
Lensing potential (2D projected) :

$$\varphi(\vec{\theta}) = \frac{2 D_{LS}}{c^2 D_s D_L} \int d\chi \phi_{3D}$$

$$\vec{\alpha} = \vec{\nabla}_{\theta} \varphi$$

Matrix amplification :

$$A^{-1} = \frac{\partial \vec{\beta}}{\partial \vec{\theta}} = \begin{pmatrix} 1 - \kappa & 0 \\ 0 & 1 - \kappa \end{pmatrix} + \begin{pmatrix} -\gamma_1 & -\gamma_2 \\ -\gamma_2 & \gamma_1 \end{pmatrix}$$



## 2. Ray tracing through Horizon AGN light cone

thin lens theory

Lensing quantities :

$$\text{Convergence : } \kappa = \frac{\Sigma}{\Sigma_{crit}} = \frac{1}{2} \Delta\varphi$$

$$\text{Shear real : } \gamma_1 = \frac{1}{2} \left( \frac{\partial^2 \varphi}{\partial \theta_1^2} - \frac{\partial^2 \varphi}{\partial \theta_2^2} \right)$$

$$\text{Shear imaginary : } \gamma_2 = \frac{\partial^2 \varphi}{\partial \theta_1 \partial \theta_2}$$

	< 0	> 0
$\kappa$		
Re[ $\gamma$ ]		
Im[ $\gamma$ ]		

# Resolution of differential lens equation in Fourier space

*Differential equation*

$$\kappa = \frac{1}{2} \Delta \varphi$$

$$\gamma_1 = \frac{1}{2} \left( \frac{\partial^2 \varphi}{\partial \theta_1^2} - \frac{\partial^2 \varphi}{\partial \theta_2^2} \right)$$

$$\gamma_2 = \frac{1}{2} \left( \frac{\partial^2 \varphi}{\partial \theta_1^2} - \frac{\partial^2 \varphi}{\partial \theta_2^2} \right)$$

$$\vec{\alpha} = \vec{\nabla}_\theta \varphi$$

*Fourier space*

$$\hat{\psi} = -2 \frac{\hat{\kappa}}{(k_1 + k_2)^2}$$

$$\hat{\gamma}_1 = -\frac{1}{2} (k_1^2 - k_2^2) \hat{\psi}$$

$$\hat{\gamma}_2 = -k_1 k_2 \hat{\psi}$$

$$\hat{\alpha}_1 = -i k_1 \hat{\psi}$$

$$\hat{\alpha}_2 = -i k_2 \hat{\psi}$$

## 2. Ray tracing through Horizon AGN light cone

Multiple lens planes

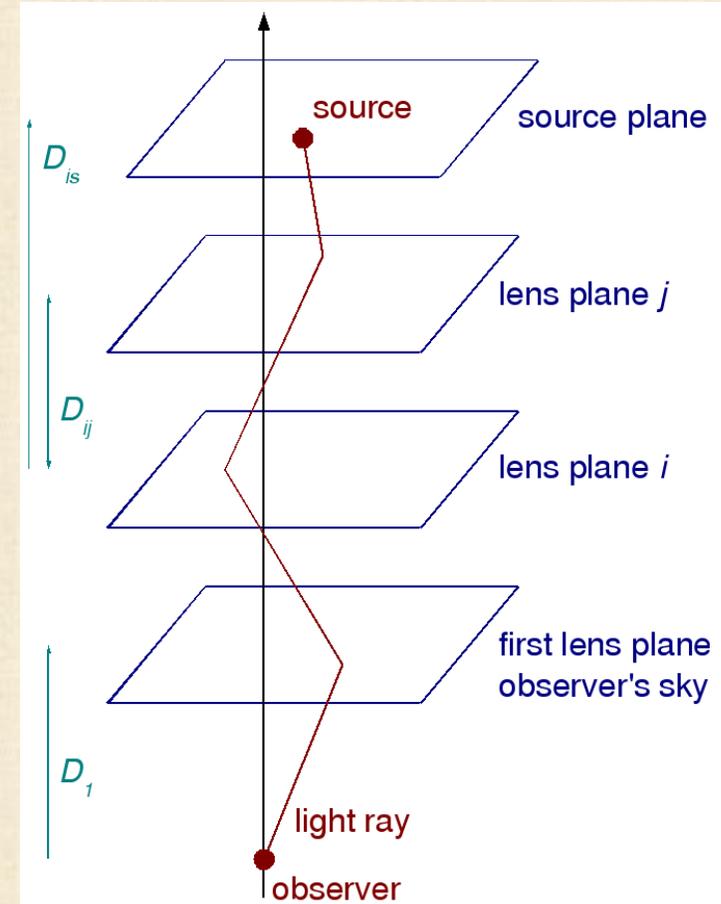
Lens equation applies on each of the multiple lens planes :

Source position :

$$\vec{\beta}_k = \vec{\theta} - \sum_{i=1}^{k-1} \frac{D_{ik} D_s}{D_k D_{is}} \vec{\alpha}_i(\vec{\beta}_i)$$

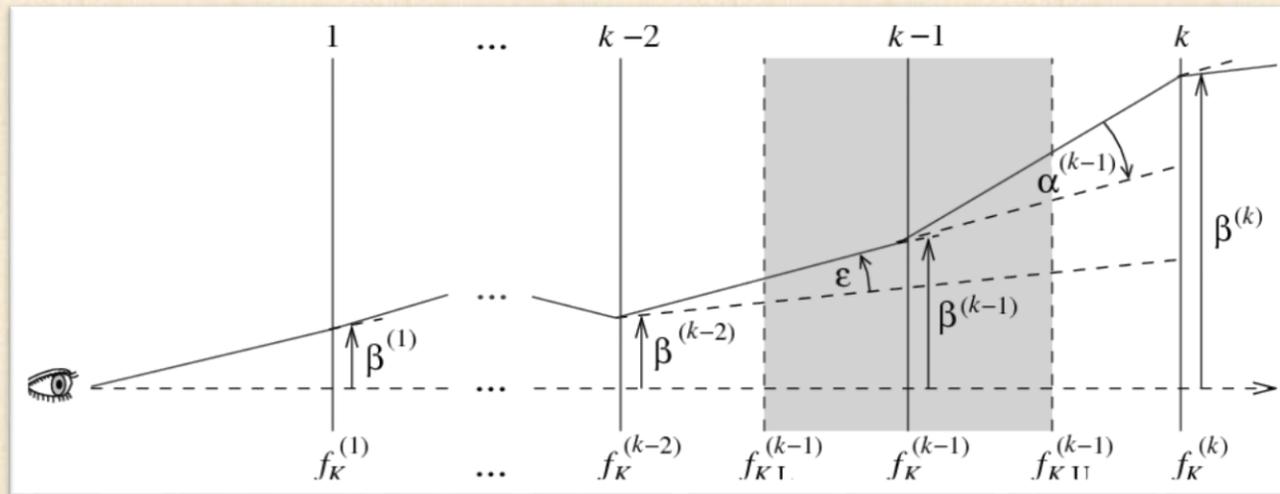
Equation for the jacobian of the lens mapping for the k-th lens plane:

$$A^k = \frac{\partial \vec{\beta}_k}{\partial \vec{\theta}} = I - \sum_{i=1}^{k-1} \frac{D_{ik} D_s}{D_k D_{is}} \frac{\partial \vec{\beta}_i}{\partial \vec{\theta}} \frac{\partial \vec{\alpha}_i}{\partial \vec{\beta}_i}$$



## 2. Ray tracing through Horizon AGN light cone

### Multiple lens planes



Credit to S. Hilbert and al, 2009

iterative method : Much less memory-demanding  
(store 3 planes, instead of several tens or hundreds)

Recurrence relation for  
source position :

$$\beta^k \propto \beta^{k-1}, \beta^{k-2}, \alpha^{k-1}$$

Recurrence relation for  
amplification matrix :

$$A^k \propto A^{k-1}, A^{k-2}, \frac{\partial \alpha^{k-1}}{\partial \beta^{k-1}}$$

## 2. Lensing signal from Horizon AGN light cone

Linking RAMSES accelerations with deflections

Acceleration field are sample by particles :

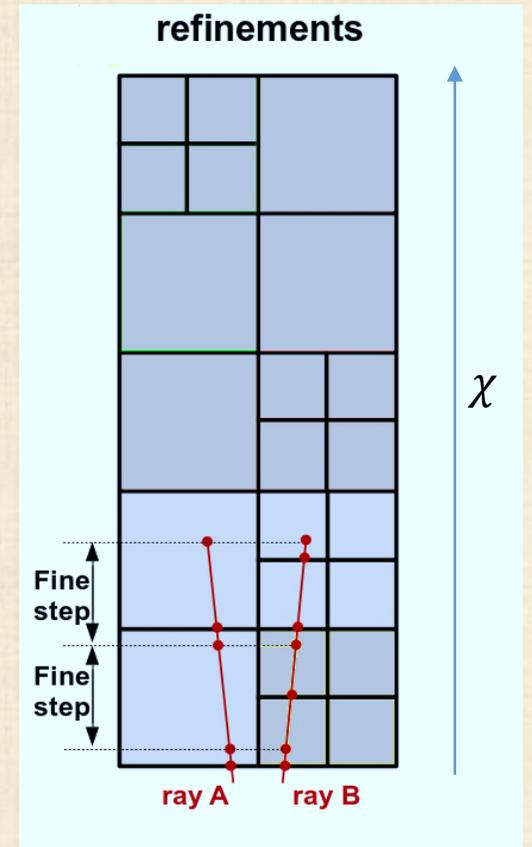
$$\vec{a} \rightarrow \vec{F}_g \rightarrow \vec{\nabla} \phi$$

$$\text{And } \vec{\alpha} = \frac{2 D_{LS}}{c^2 D_L} \int \vec{\nabla}_{\perp} \phi \, d\chi \text{ (integrate along l.o.s)}$$

Using the gas component :

- Gas particles follow the RAMSES grid

- Given the size of RAMSES cells :  $\sigma = \left( \frac{m_{gas}}{\rho_{gas}} \right)^{1/3}$



## 2. Linking RAMSES accelerations with deflections

### Adaptative Gaussian kernel 2D ( x y direction)

Each gas particle is defined by  $\sigma_p, \vec{x}_p$ , and is treated as truncated gaussian kernels

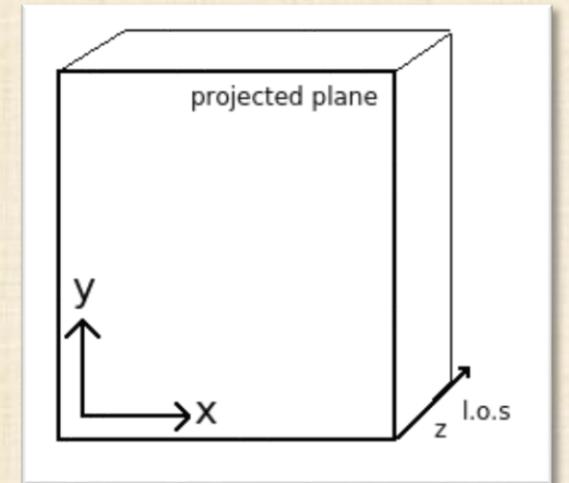
$$f(x) = e^{-\frac{(x-x_p)^2}{2\sigma_p^2}} \text{ for } \frac{x_{pixel} - x_{part}}{\sigma_p} \leq 4$$

With pixel weight

$$F_x = \int_{x_{pixel\_min}}^{x_{pixel\_max}} e^{-\frac{(x-x_p)^2}{2\sigma_p^2}} dx$$

Along the l.o.s (z direction) :

$$F_z \propto \sigma_p$$



Tabulation of the erf function (time consuming)

# Linking RAMSES accelerations with deflections

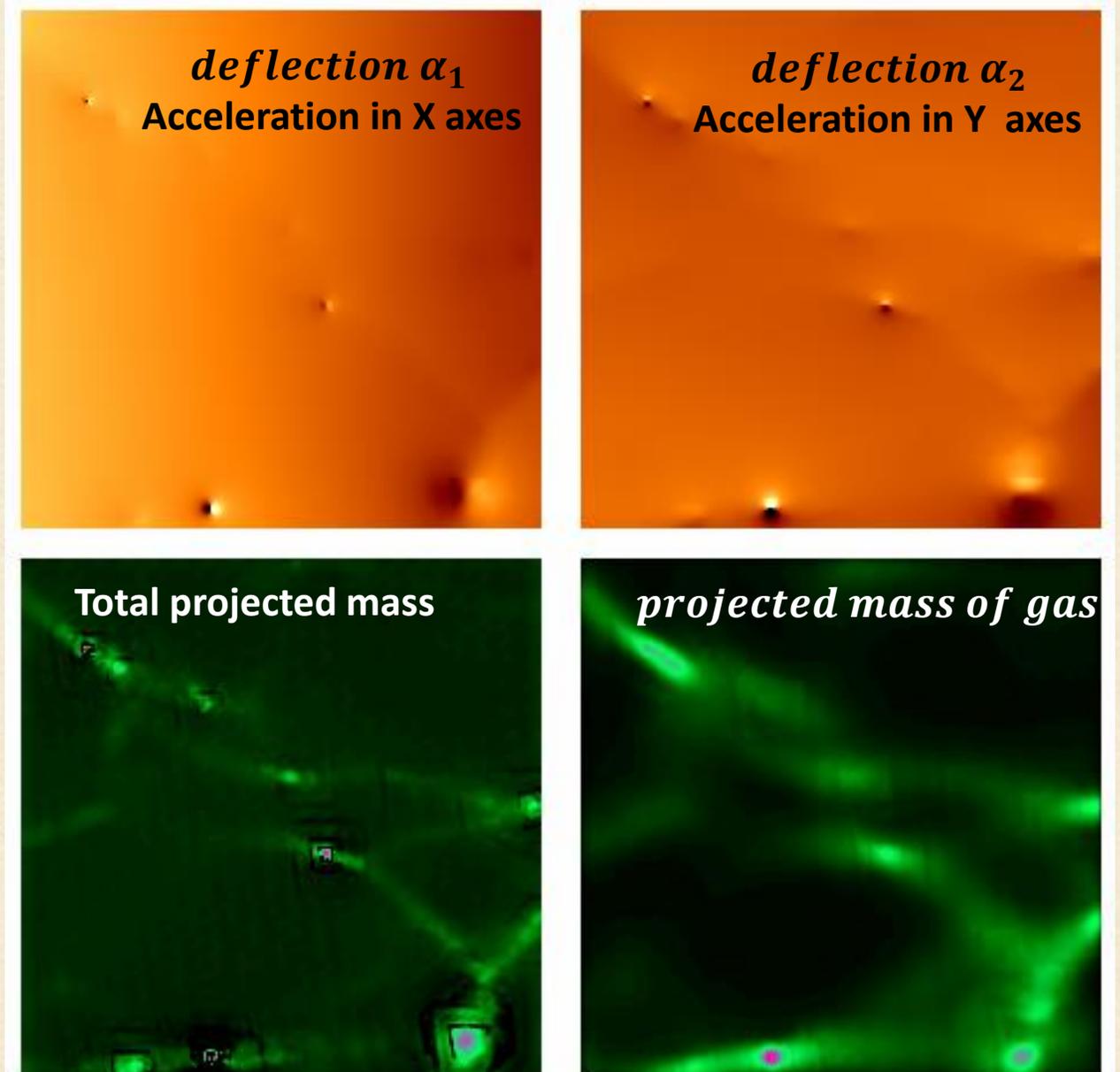
Reconstructed convergence from  
the deflection ( $\propto$  *acceleration*)

$$\vec{\nabla} \cdot \vec{\alpha} = 2 \kappa$$

artefact problems :

- Integration of acceleration along  
the RAMSES grid seems wrong ....

*work in process...*

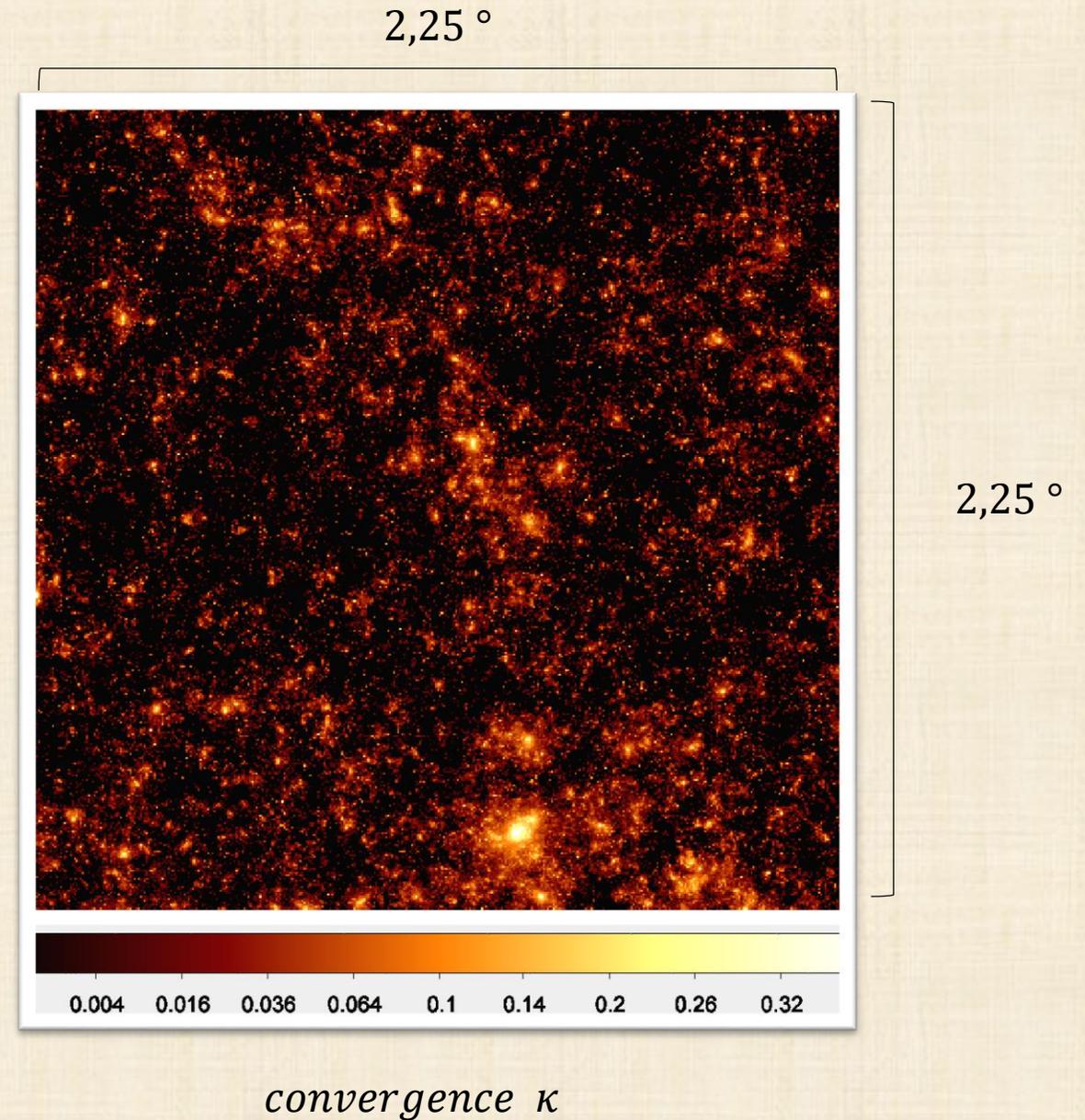
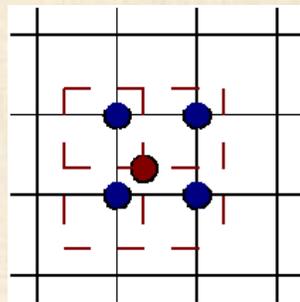


# Using the projected mass Identify the contribution of Different components

## Convergence map $\kappa$ by projected DM particles

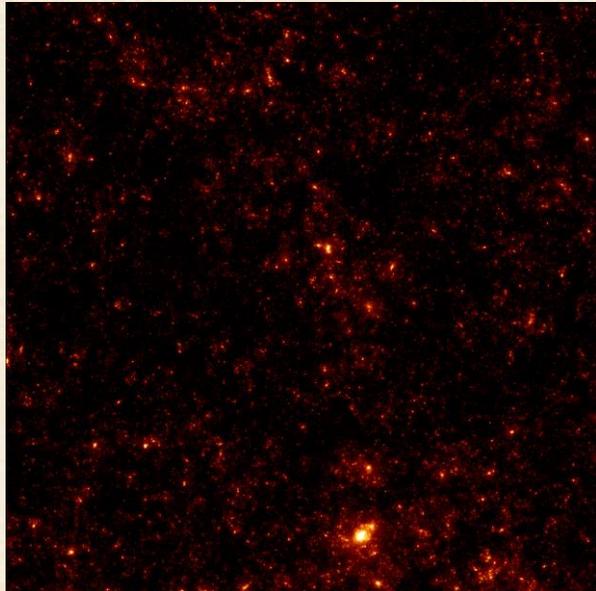
- Redshift of source plane  $z_s \sim 3$
- Propagation light ray
- Angular resolution :  $\Delta\theta = 1 \text{ arcsec}$
- 16 lightcone slices = 1 lens plane  
From  $z \sim 0$  to  $z \sim 3$  : 1064 lens planes

- DM particles only  
with Cloud In Cell



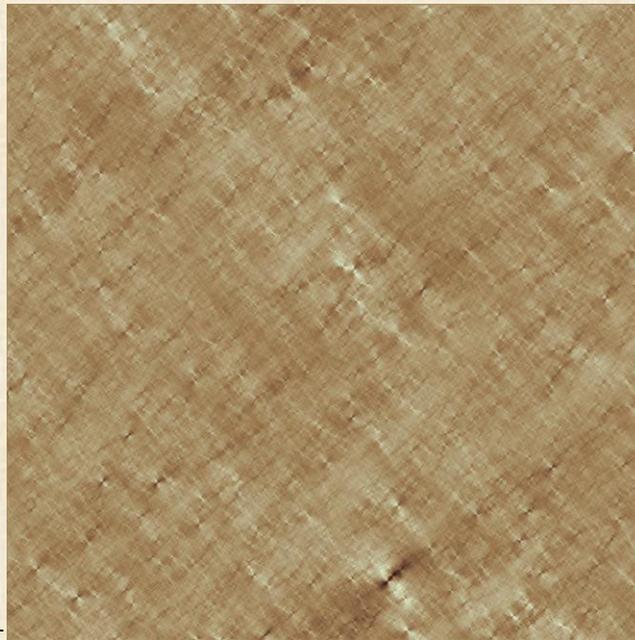
Using the projected mass  
Classical approach

Lensing quantities



*convergence  $\kappa$*

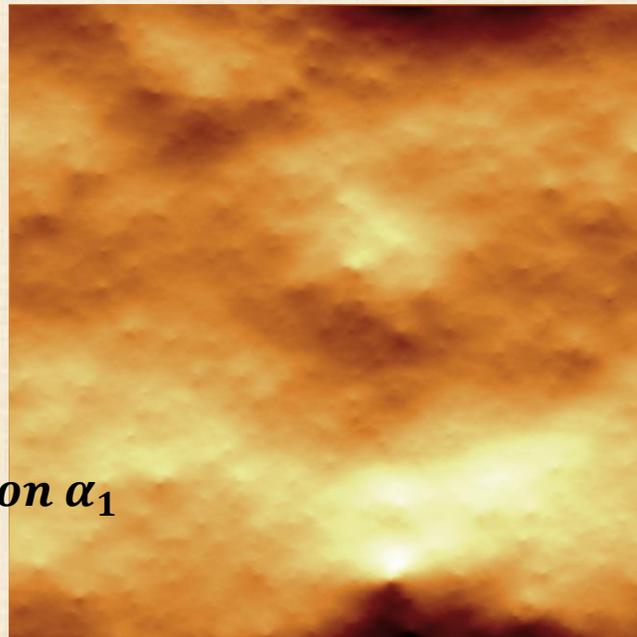
*shear  $\gamma_1$*



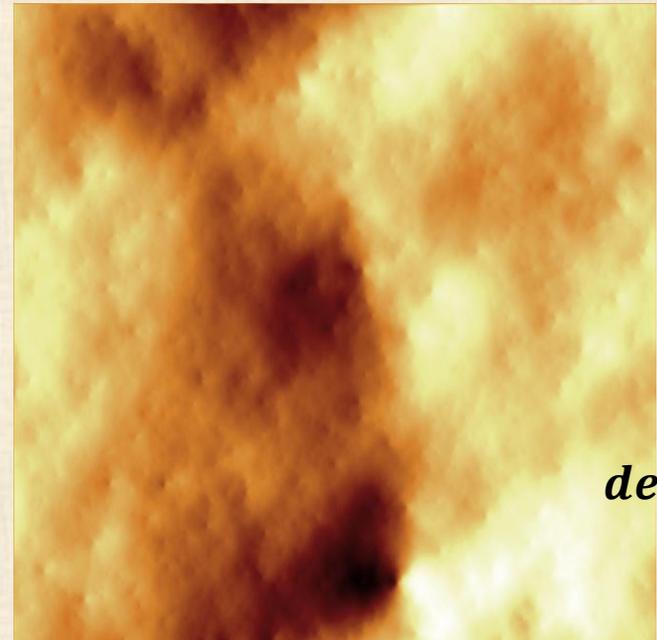
*shear  $\gamma_2$*



*deflection  $\alpha_1$*



*deflection  $\alpha_2$*

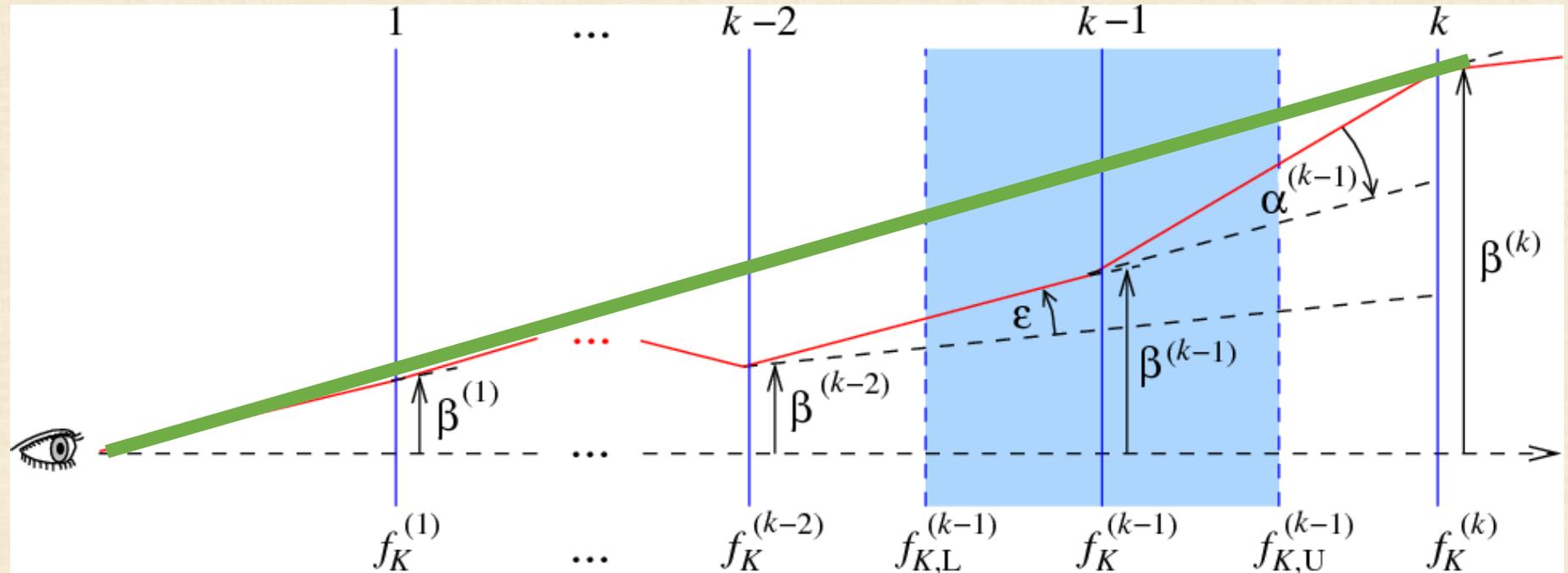


# Test : comparaisn of the propagation of light ray with Born approximation

## Born approximation

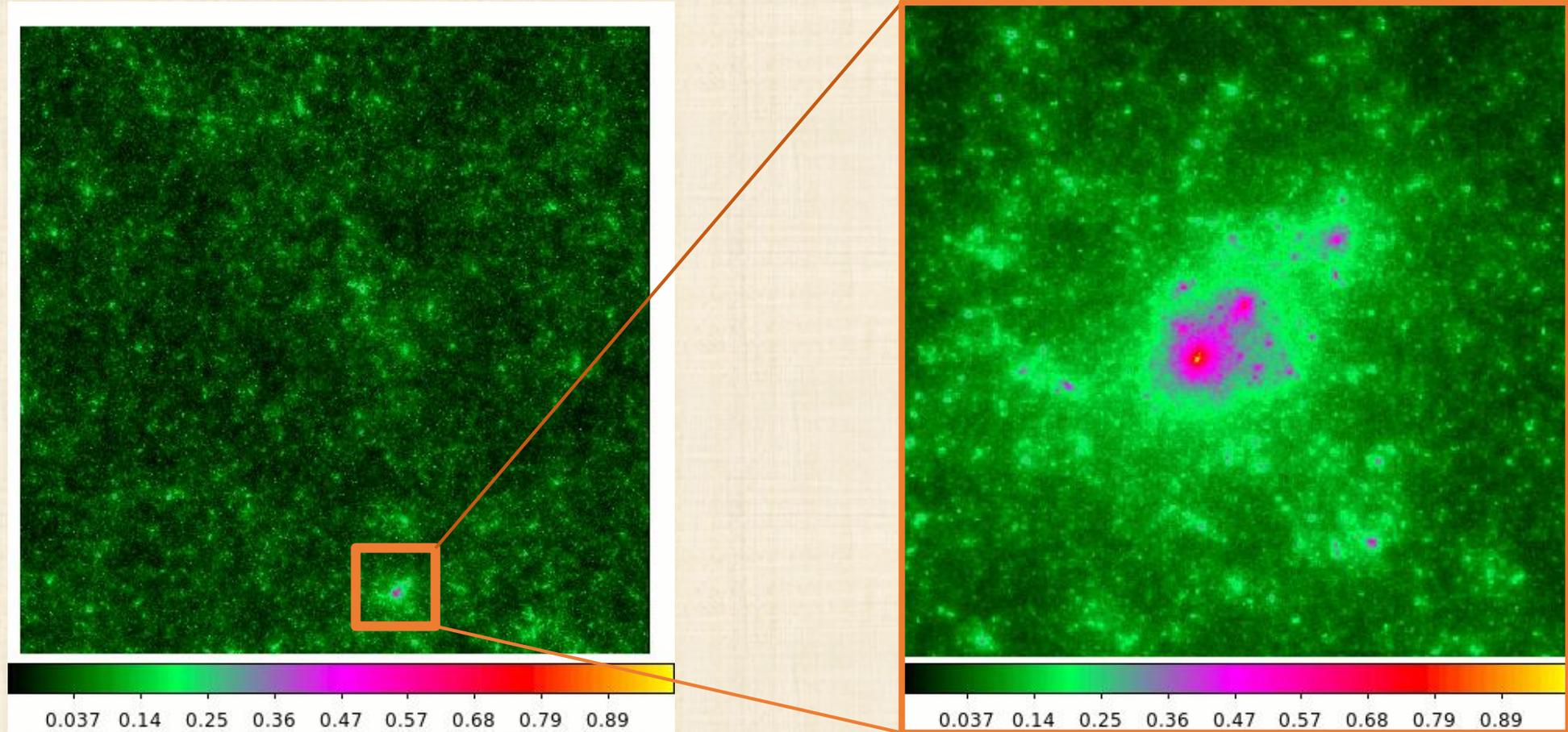
- Summation of mass contribution along a undistrubed path

(ignore deviation of light path at each lens plane )



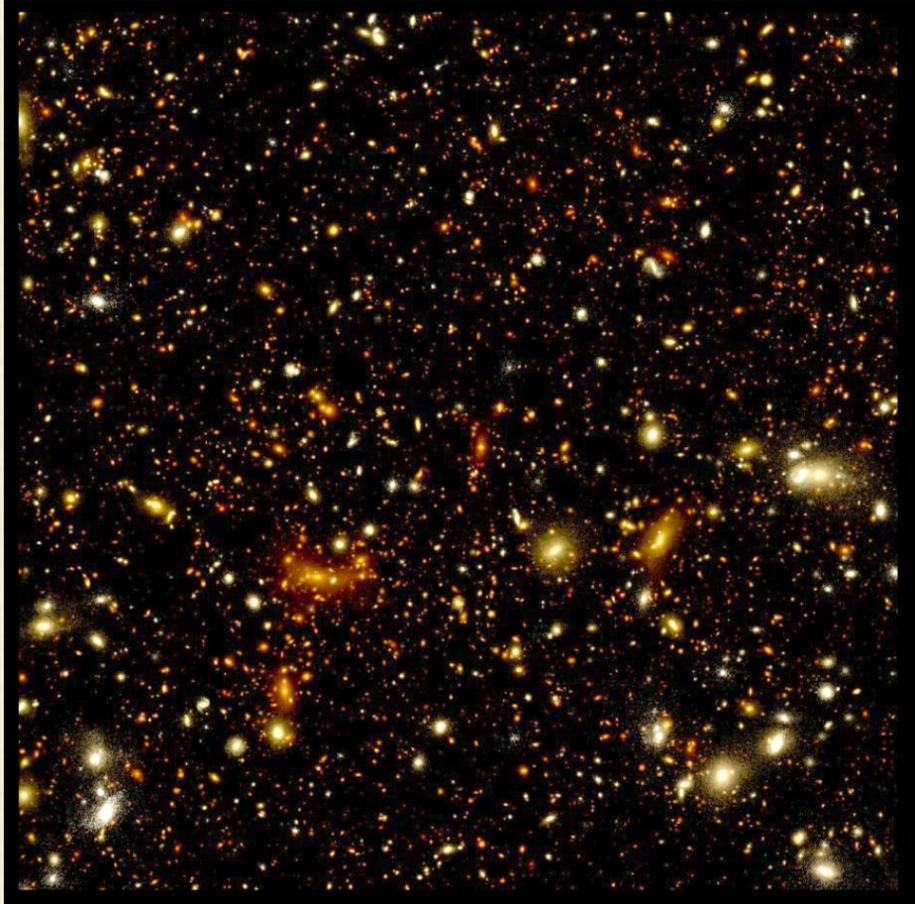
# Comparison of the propagation of light ray / Born approximation

Born approximation is appropriate at large scale but not adapted to SL at small scale

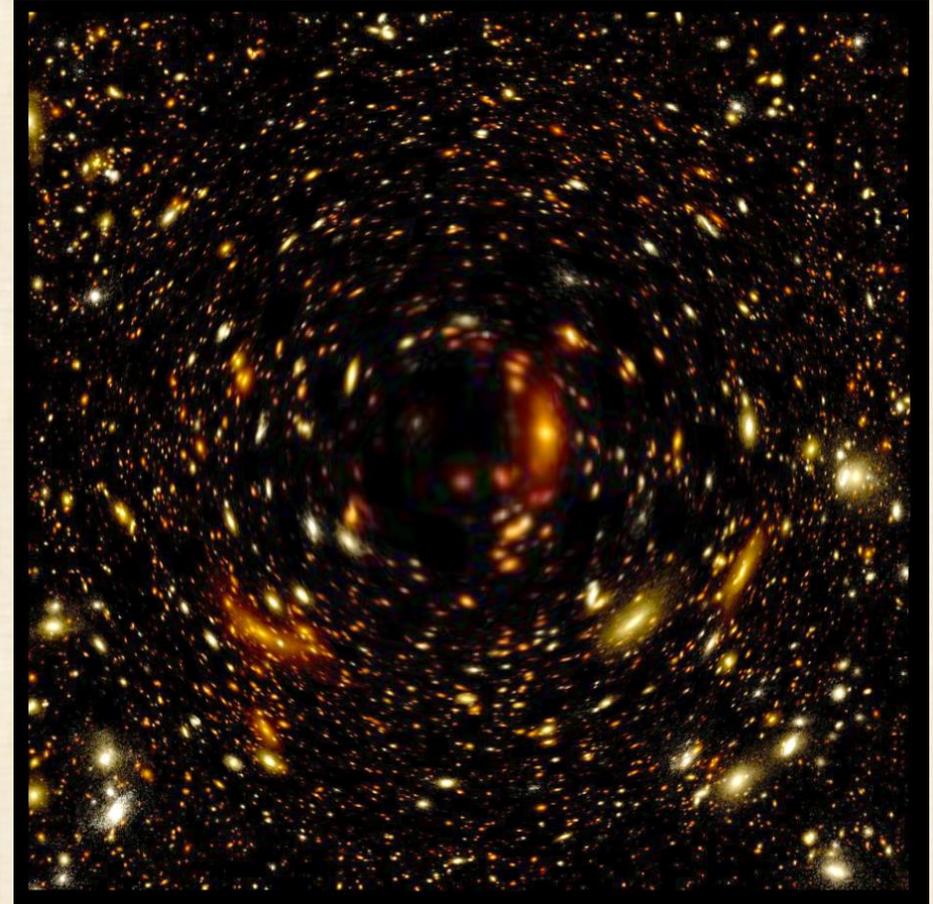


# Short term goal

Adding gravitational lensing signal on mock observation of galaxies (S. Kaviraj and al, 2016)



Mock image ©LaigleC

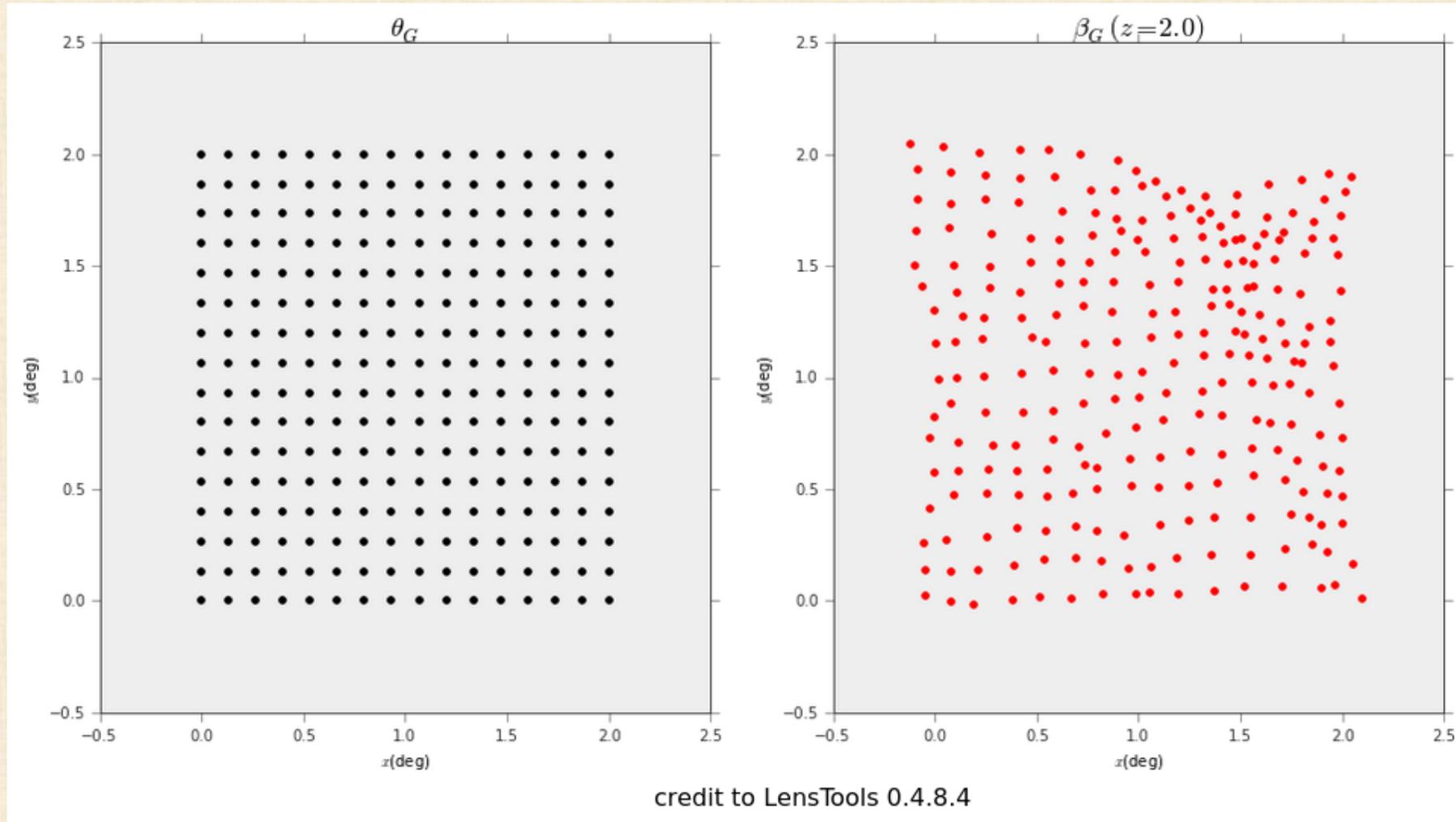


Lensed Mock image by an arbitrary lens

Thank you for your attention

## 2. Ray tracing through Horizon AGN light cone

From a regular grid image to the source position :  $\vec{\beta} = \vec{\theta} - \vec{\alpha}$



# Details on recursive relation

Lens equation apply on the multiple lens planes approximation :

$$\vec{\theta}_j = \vec{\theta}_1 - \sum \frac{D_{ij} D_s}{D_j D_{is}} \vec{\alpha}_i$$

Recursion relation :

$$A^k(\vec{\theta}) = I - \sum_{i=1}^{k-1} \frac{D_{ik} D_s}{D_k D_{is}} A^i U^i$$

$$A_i = \frac{\partial \vec{\beta}_i}{\partial \vec{\theta}_1} \text{ and } U^i = \frac{\partial \vec{\alpha}_i}{\partial \vec{\beta}_i}$$

The number of sub-haloes is it reduced by the presence of baryons? (Guilia Despali & Simona Vegetti, 2016)

- Investigate EAGLE and ILLUSTRIS simulation / concentrate on ETGs
- Especially at low mass ( $\leq 10^{10} M_{\odot}/h$ ), by different amounts depending on the model
- They attempt to predict the DM fraction in subhaloes and the slope of mass function  $\alpha$

# Strong lensing details

Study of substructure via :

The relative flux of multiply imaged quasar

Their effect on surface brightness of einstein rings & lensed arc

# Euclid



# Time delays

SL : measured time delay between the multiple images & models of mass distribution  
→ determination of time delay distance → cosmological parameters

Excess time :  $t(\theta, \beta) = \frac{D_{\Delta t}}{c} \left( \frac{(\theta - \beta)^2}{2} - \psi(\theta) \right)$  with  $D_{\Delta t}$  time delay distance

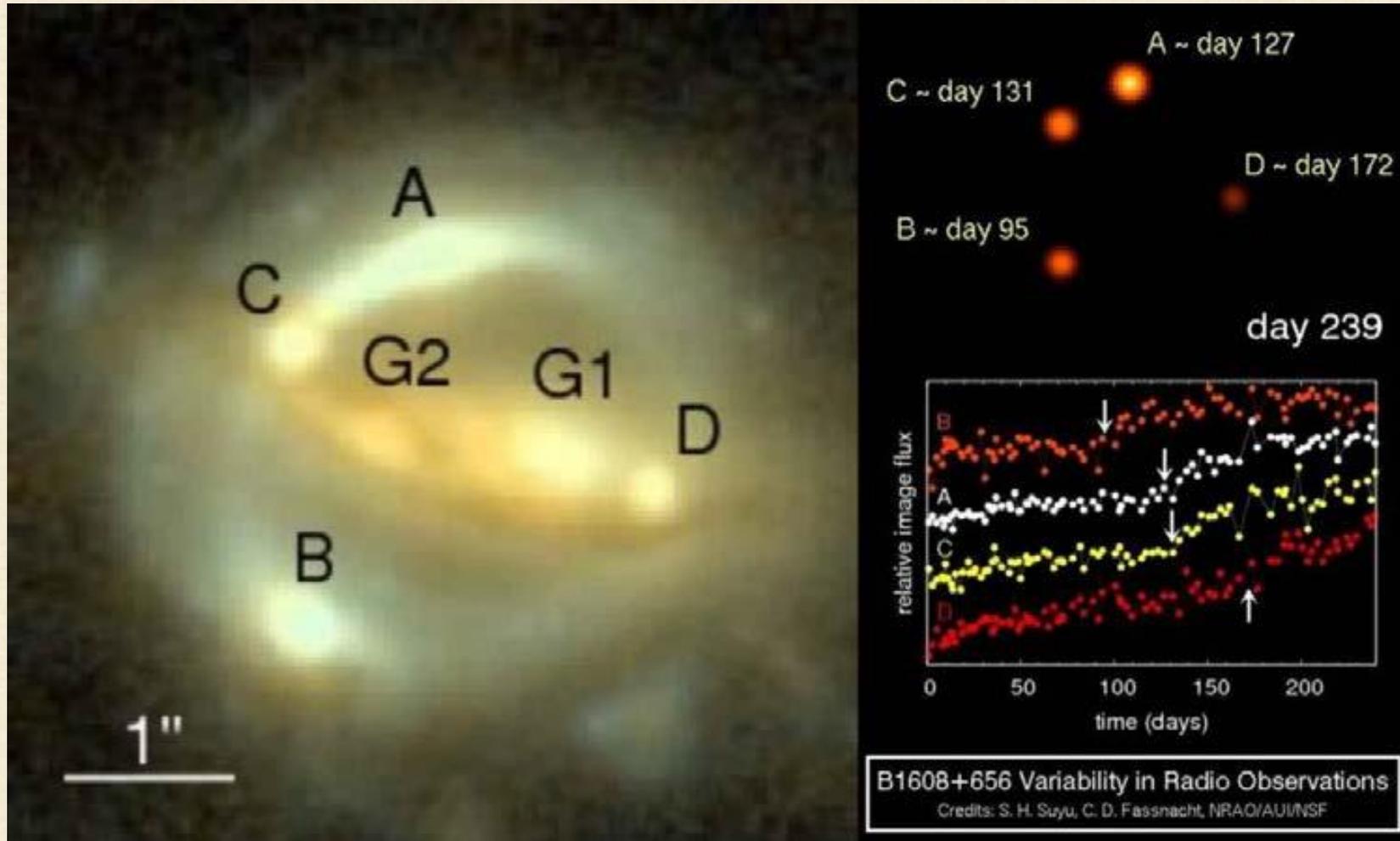
Time delay distance :  $D_{\Delta t} = (1 + z_d) \frac{D_d D_s}{D_{ds}}$

Time delay  $\Delta t_{ij} = t(\theta_i, \beta) - t(\theta_j, \beta) = \frac{D_{\Delta t}}{c} \left( \frac{(\theta_i - \beta)^2}{2} - \psi_i(\theta_i) - \frac{(\theta_j - \beta)^2}{2} - \psi_j(\theta_j) \right)$

Image configuration+morphology →  $\Sigma, \psi$  + source vary in time  $\Delta t$  → obtaint  $D_{\Delta t}$

Contraint cosmological parameter by distance redshift test  $D_{\Delta t} \propto \frac{1}{H_0}$

# Time delay



# « Lens galaxies in Illustris simulation : power law & the bias of Hubble constant from time delay » Xu and al 2015

The total density profile in central region is usually described by Power law  $\rho \propto r^{-\gamma'}$

Radial scale : transition from dominance of BM to DM

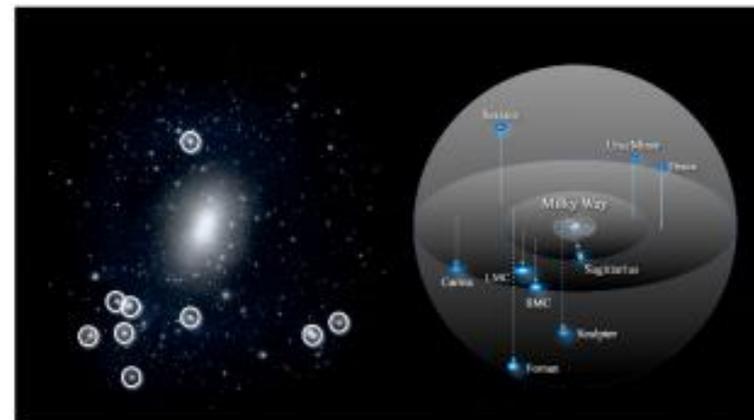
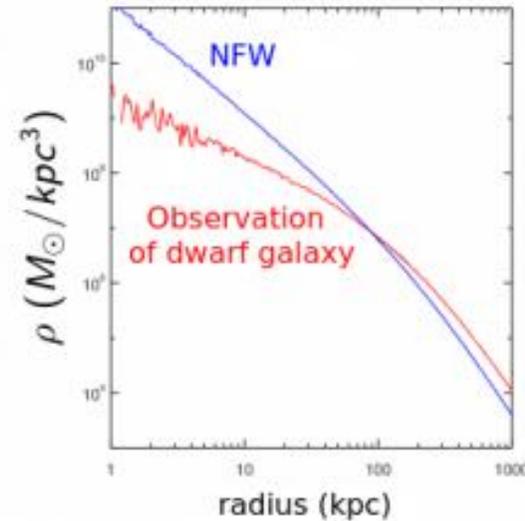
PL affect the product of time delay  $\rightarrow$  bias the determination of  $H_0$

They study dynamic and SL on simulated galaxies (Illustris)

Find : the bias on  $H_0$  introduced by PL assumption can reach 20-50 %

## Testing the CDM Paradigm on Small Non Linear Scales

- density profile of DM halo fitting by universal NFW profile
  - *Cusp-core problem*
- Missing satellites
  - $\rho (M_{\odot}/kpc^3)$
  - CDM model clearly over-predicts the number of substructure in the Milky Way



Simulated Cold Dark Matter Halo

(Weinberg and al 2013)