

# **Analytical Gravitational Self-force:**

## **Non-circular motion+application to EOB**

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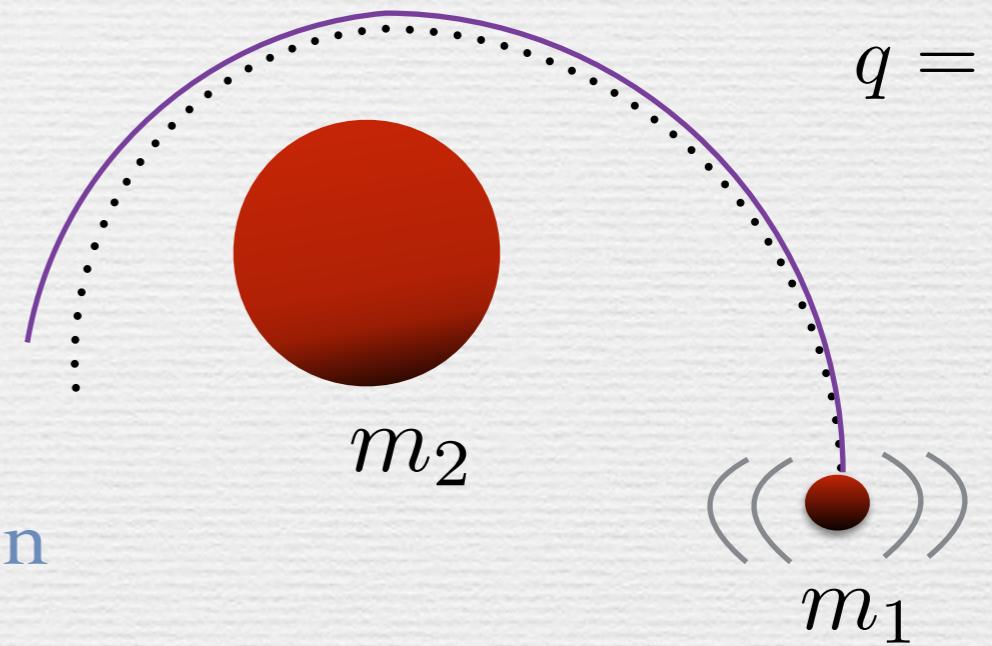
(mostly) based on arXiv:1706.00459 [gr-qc]

What do we mean? See M. Van de Meent tomorrow!

$$m_1 \frac{Du^\alpha}{d\tau} = q^2 F^\alpha[h]$$
$$\delta G_{\mu\nu}[h_{\mu\nu}] = 8\pi T_{\mu\nu}$$

$$g_{\mu\nu} = g_{\mu\nu}^0 + q h_{\mu\nu} + \mathcal{O}(q^2)$$

$$q = \frac{m_1}{m_2}$$



- Osculating assumption--for now, motion is geodesic.

### Aim

- Calculate anything of interest in EMRI modelling without numerics--typically asymptotic expansions (think PN)

analytical SF  $\longleftrightarrow$  post-Newtonian SF

# Dissipative + conservative effects

$$F^\alpha[h] = F_{\text{diss}}^\alpha[h] + F_{\text{cons}}^\alpha[h] \longrightarrow \text{e.g. shift in periastron adv.}$$

e.g.  $\langle \frac{1}{u^t} F_t \rangle = \langle \frac{dE^\infty}{dt} \rangle$

- Fluxes are dominant, but ‘easy’ — exp. convergence

$$\frac{dE^\infty}{dt} = \sum_{\ell m} \frac{1}{4\pi\omega_m^2} |Z_{\ell m \omega}|^2 \quad \psi_{0,\ell m \omega} \rightarrow Z_{\ell m \omega} r^3 e^{i\omega r^*}$$

$$\frac{dE^\infty}{dt} = -\frac{32}{5} \frac{1}{r^5} \left( 1 - \frac{1247}{336r} + \frac{4\pi}{r^{1.5}} + \dots \right) \quad - \text{Teukolsky equation}$$

Finite PN order



Finite  $\ell$  contribution

- Their calculation is well developed analytically
  - Mano, Suzuki & Takasugi (~1996): analytic solns Teukolsky equation
  - eg. 22PN fluxes Schw & 11PN Kerr (Fujita 2015), Generic eccentric inclined Kerr (Sago, Fujita 2015).

# Post-Newtonian SF: conservative sector

- Focus in recent years has been on linear in the mass ratio conservative effects, e.g. SF correction to classical redshift

- ♦ Need to deal with regularisation - ~solved (Barack+collaborators, Detweiler-Whiting), e.g. mode-sum approach

$$F^\mu = \sum_{\ell=0}^{\infty} (F^{\mu,\ell} \mp A^\mu(2\ell + 1) - B^\mu + \mathcal{O}(\ell^{-2}))$$

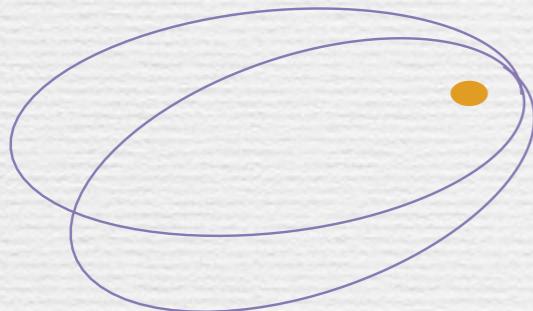


Polynomial convergence

- Not a problem for PN

Bini, Damour (2014)- Soln's Teukolsky equation valid for all  $\ell$ 's

## Non-circular case: Situation to date



$$r_p(\chi) = \frac{pm_2}{1 + e \cos \chi}, \quad 0 < \chi < 2\pi$$

$$\frac{p}{1+e} < r_p/m_2 < \frac{p}{1-e} \quad \Omega_r \neq \Omega_\varphi$$

$\frac{1}{p}, e \ll 1 :$



SF correction to  $u^t$ , encapsulated in  $h_{\mu\nu} u^\mu u^\nu$

Generalized redshift  $\langle U \rangle$ :

- Bini et al (2016): 6.5PN,  $e^2$  Schwarzschild
- Hopper et al (2016): 4PN,  $e^{10}$  Schw.
- Bini et al (2016): 8.5PN,  $e^2$  Kerr.

we want more...

# Non-circular case: next level

- Want self-force
- Schwarzschild, but Radiation Gauge (eye on Kerr)

$$F^\mu = \mathcal{P}^{\mu\nu\lambda\rho}(2h_{\nu\lambda;\rho} - h_{\lambda\rho;\nu})$$

$$\cancel{\delta G_{\mu\nu}[h_{\mu\nu}] = 8\pi T_{\mu\nu}}$$



Teukolsky equation

$${}_s\mathcal{O}\psi = \mathcal{T}[x^\mu(t)]$$



$$h_{\mu\nu}, h_{\mu\nu,\rho}$$

CCK reconstruction\*

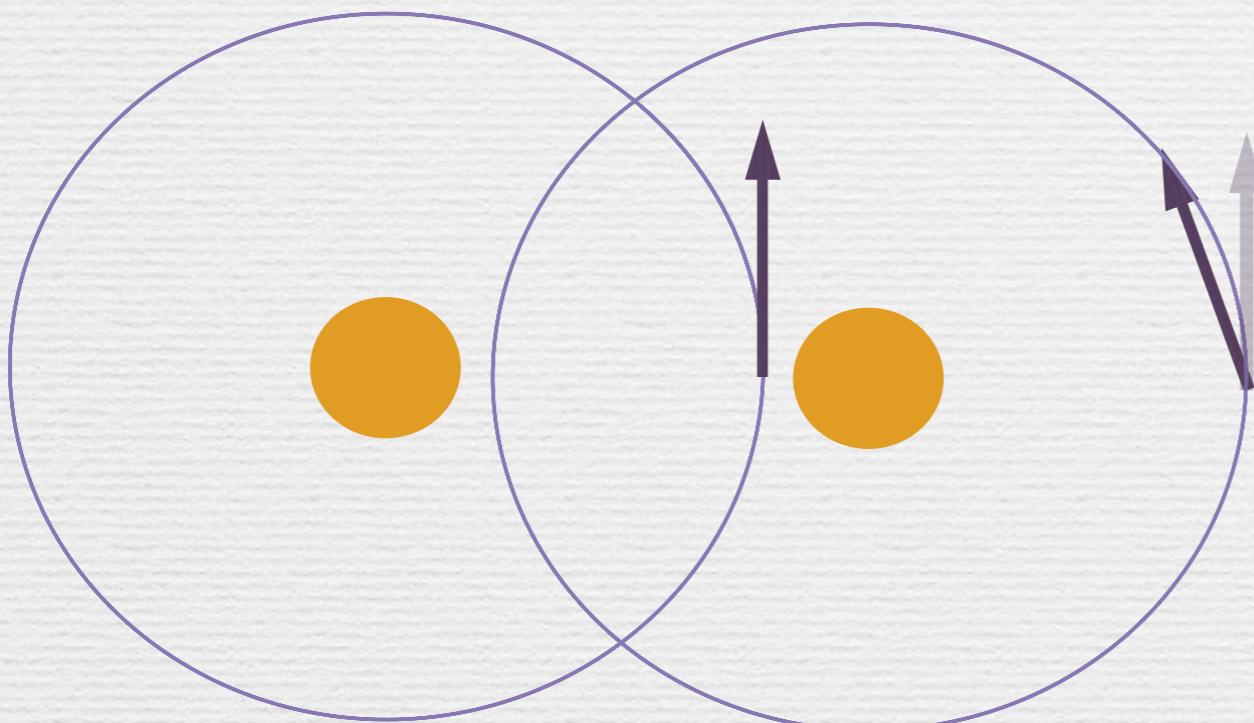
$$h_{tt,t}^{\ell=2} = (4 \sin \chi e + 4 \sin(2\chi) e^2 + \dots) p^{-5/2} + (6 + 20 \cos \chi e + (10 + 14 \cos(2\chi)) e^2 + \dots) p^{-3} + \dots$$

7 orders in  $1/p$   
including to  $e^4$   
all  $\ell$

Can we do an inspiral? maybe, but first...

# Spin-precession: Circular overview

The orientation of a test spin will precess on a curved background



$$u^a \nabla_a s^b = 0$$

$$\frac{d\mathbf{s}}{d\tau} = \omega_s \times \mathbf{s}$$

$$\psi = 1 - \frac{|\omega_s|}{u^\varphi}$$

- Dolan et al (2014)-- 1sf correction to this is gauge invariant - schw background

$$\psi = \psi_0 + q\Delta\psi + \dots$$

# Spin-precession: Eccentric Schwarzschild

Akcay, Dolan & Dempsey (2016)

$$\psi = 1 - \frac{\Psi}{\Phi}$$

Accumulated precession/ radial period

Acc. azimuthal phase/radial period

+prescription for calculating the 1-SF correction, explicitly requiring knowledge of the SF, and derivatives of the metric perturbation.

ADD give, strong field numerical data in Lorenz gauge, and a PN expansion from the full PN Hamiltonian:

# Spin-precession: Eccentric Schwarzschild

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Akcay, Dolan & Dempsey (2016)

ADD give, strong field numerical data in Lorenz gauge, and a 3PN expansion from the full PN hamiltonian:



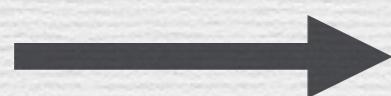
- Add terms to 6PN
- Independent check of the gauge invariance.
- Complete the parameter space.
- Validate our code.

# Spin-precession: Results

$$\Delta\psi = \Delta\psi_0 + \Delta\psi_2 e^2 + \mathcal{O}(e^4)$$

$$\begin{aligned} \Delta\psi_2 = & \frac{1}{p^2} + \frac{\frac{341}{16} - \frac{123\pi^2}{256}}{p^3} + \frac{-\frac{164123}{480} + \frac{536\gamma}{5} - \frac{268\ln p}{5} - \frac{23729\pi^2}{4096} + \frac{11720\ln 2}{3} - \frac{10206\ln 3}{5}}{p^4} \\ & + \frac{-\frac{89576921}{57600} - \frac{22682\gamma}{15} + \frac{11341\ln p}{15} + \frac{21333485\pi^2}{49152} - \frac{4836254\ln 2}{105} + \frac{4430133\ln 3}{320} + \frac{9765625\ln 5}{1344} + \frac{319609\pi}{630}}{p^5} + \frac{319609\pi}{630} \\ & - \frac{\frac{464068669129}{5080320} + \frac{2508913\gamma}{945} - \frac{2508913\ln p}{1890} - \frac{32088966503\pi^2}{2359296} + \frac{146026515\pi^4}{1048576} - \frac{273329813\ln 2}{945} + \frac{159335343\ln 3}{8960} + \frac{17193359375\ln 5}{145152}}{p^6} \\ & + \mathcal{O}(p^{-13/2}) \end{aligned}$$

- Leading orders agree
- Agreement within numerical error



Passed strenuous test of code.

# Effective-one-body approach: See T. Damour tomorrow!

Proposal: Knowledge of the SF spin precession can determine terms in the EOB Hamiltonian (circ. case Bini,Damour...)

$$\mathcal{H}(R, P, S_1, S_2) = Mc^2 \sqrt{1 + 2\nu \left( \frac{\mathcal{H}_{\text{eff}}}{\mu c^2} - 1 \right)}$$

$\nu$ -symmetric mass-ratio

$$\mathcal{H}_{\text{eff}} = \mathcal{H}_{\text{eff}}^{\text{O}} + \mathcal{H}_{\text{eff}}^{\text{SO}}$$
$$\mathcal{H}_{\text{eff}}^{\text{SO}} = \frac{G}{c^2 R^3} (g_{\mathbf{S}} \mathbf{L} \cdot \mathbf{S} + g_{\mathbf{S}^*} \mathbf{L} \cdot \mathbf{S}^*)$$
$$\mathbf{S} = \mathbf{S}_1 + \mathbf{S}_2$$
$$\mathbf{S}^* = \frac{m_2}{m_1} \mathbf{S}_1 + \frac{m_1}{m_2} \mathbf{S}_2$$

The spin vectors satisfy a familiar equation..

$$\begin{aligned} \frac{d\mathbf{S}_a}{dt} &= \{\mathcal{H}, \mathbf{S}_a\} \\ &= \Omega_{\mathbf{S}_a} \times \mathbf{S}_a \end{aligned}$$

$$\Omega_{\mathbf{S}_a} = \frac{\partial \mathcal{H}}{\partial \mathbf{S}_a}$$

# EOB spin-precession

$$\begin{aligned}\frac{d\mathbf{S}_a}{dt} &= \{\mathcal{H}, \mathbf{S}_a\} \\ &= \boldsymbol{\Omega}_{\mathbf{S}_a} \times \mathbf{S}_a\end{aligned}$$

$$\boldsymbol{\Omega}_{\mathbf{S}_a} = \frac{\partial \mathcal{H}}{\partial \mathbf{S}_a}$$

Strategy: 1) Set up same type situation..

$$m_1 \ll m_2, \quad S_1 \ll 1, \quad S_2 = 0, \quad P_R \neq 0, \quad \mathbf{L} \cdot \mathbf{S} = P_\varphi S$$


$$\psi \equiv \frac{\langle \boldsymbol{\Omega}_{\mathbf{S}_1} \rangle}{\Omega_\varphi}$$

is the same function!

2) Extract SF info by equating

$$\mathcal{O}(\nu) \text{ piece of } \psi^{\text{EOB}} \quad \longleftrightarrow \quad \Delta\psi$$

via gauge inv. parameterisation.

(straightforward, but delicate)

# EOB spin-precession: what does SF give?

e.g.

$$M\Omega_{S_1} = p_\phi(1 - \nu \hat{\mathcal{E}}_{\text{eff}})u^3[\nu g_S + (1 - \nu)g_{S*}] + \mathcal{O}(\nu^2)$$



$$\begin{aligned} g_S &= 2 + \nu g_S^1(u, p_r) + \mathcal{O}(\nu^2) \\ g_{S*} &= g_{S*}^0(u, p_r) + \nu g_{S*}^1(u, p_r) + \mathcal{O}(\nu^2) \end{aligned}$$

$\Delta\psi$  picks out  $g_{S*}^1(u, p_r)$

$$g_{S*}^1(u, p_r) = g_{S*}^{1,0}(u) + g_{S*}^{1,2}(u)p_r^2 + \mathcal{O}(p_r^4)$$

$$p_r^2 \sim \mathcal{O}(e^2)$$

# EOB spin-precession: Results

$$g_{S^*}^{1,2}(u) = -\frac{9}{4} - \frac{9}{4}u - \frac{717}{32}u^2 + \left( \frac{1447441}{960} - \frac{4829}{256}\pi^2 - \frac{16038}{5}\ln(3) + \frac{46976}{15}\ln(2) - \frac{512}{5} + \gamma - \frac{256}{5}\ln(u) \right)u^3 \\ - \left( \frac{185195453}{38400} + \frac{19162}{35}\gamma + \frac{2097479}{8192}\pi^2 + \frac{454167}{20}\ln(3) - \frac{1081966}{35}\ln(2) + \frac{9581}{35}\ln(u) \right)u^4 + \dots$$

..The coefficient of  $p_r^2$  in the linear in  $\nu$  piece of

$$\mathcal{H}_{\text{eff}}^{\text{SO}} = \frac{G}{c^2 R^3} (g_S \mathbf{L} \cdot \mathbf{S} + g_{S^*} \mathbf{L} \cdot \mathbf{S}^*)$$

i.e. EOB can handle spin-orbit interactions  
with greater accuracy!

# Conclusions and outlook

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## Results

1.  $h_{\mu\nu}, h_{\mu\nu,\rho}, F^\mu$  – high order PN + eccentricity expansions
2. Extended knowledge of  $\Delta\psi$ , verified results of Akcay Dolan & Dempsey
3. Improved knowledge of gyrogravitomagnetic ratio  $g_S^*$

## What's next?

1. Investigate the SF in radiation gauge?
2. Application to osculating inspiral?
3. PN always needs resummation!
4. SF in Kerr spacetime