Gravitational waves from core collapse supernovae



für Astrophysik

Ewald Müller: IAP Colloquium, Paris (France), 06/30/2017

Aim of the talk:

- review of the current status of models/simulations aiming to predict the gravitational wave signal of core collapse supernovae (CCSN)

Outline of the talk:

- why to occupy oneself with gravitational waves from CCSN?
- processes that generate gravitational waves in CCSN
- generic features of the gravitational wave signal from CCSN
- achievements and predictions of the most recent models

Disclaimer:

- only selected references are given, no attempt for completness
- more references, see:
 - http://wwwmpa.mpa-garching.mpg.de/rel_hydro/GWlit_catalog.shtml

Collaborators: Haakon Andresen, Thomas Janka, Bernhard Müller



"Looking" into the "engine" of a core collapse supernova

- through observations of neutrinos up to now only SN1987A
- through observations of gravitational waves not yet occured! Would provide kind of Rosetta stone!
- through simulations

already a 50 year effort ; extremely complex & expensive 6D radiation-hydrodynamics problem requiring ~50 million CPU-hrs / simulation

Core collapse supernovae: neutrino-driven delayed explosion



Explosion mechanism

the computational challenge

6D time-dependent radiation-hydrodynamics problem

Boltzmann equation determines neutrino distribution function in phase space

Integration over momentum space yields source terms for hydrodynamics

Various modelling approaches

- **3D** hydro + "**ray-by-ray-plus**" variable Eddington factor method: current method used at MPA (Janka et al.)
- **3D** hydro + **two moment closure** of Boltzmann equation (next feasible step)
- **3D** hydro + **6D** direct discretization of Boltzmann equation: no serious attempt yet



Required resources 0.1 - 1 PFLOP/s 1 - 10 PFLOP/s

~100 PFLOP/s (sustained!)

Einstein quadrupole formula (adequate for CCSN (Shibata+ '05, Reiswig+ '11))

$$h_{ij}^{TT}(\boldsymbol{X},t) = \frac{2G}{c^4 R} P_{ijkl}(\boldsymbol{N}) \quad \frac{\partial^2}{\partial t^2} Q_{kl}(t - \frac{R}{c})$$

mass-quadrupole tensor

$$Q_{ij}(t) = \int dV \rho(\mathbf{x}, t) (x_i x_j - \frac{1}{3} \delta_{ij} \mathbf{x}^2)$$

numerically more suitable form (Nakamura & Oohara 1989, Blanchet+ 1990)

$$h_{ij}^{TT}(\boldsymbol{X},t) = \frac{2G}{c^4 R} P_{ijkl}(\boldsymbol{N}) \int d\boldsymbol{V} \rho(2v^k v^l - x^k \partial_l \Phi - x^l \partial_k \Phi)$$

Gravitational waves

Einstein quadrupole formula
$$h_{ij}^{TT} = \frac{2G}{c^4 R} \frac{\partial^2}{\partial t^2} Q_{kl} \sim \frac{R_s}{R} \frac{v^2}{c^2}$$

 $R_s = 3 \text{ km}, \text{ v/c} = 0.1, \text{ R} = 10 \text{ kpc} --> h = 10^{-20} \text{ *}$

time-dependent mass-energy quadrupole moment in core collapse supernovae due to

- convection in proto-neutron star
- convection in neutrino heated hot bubble
- anisotropic neutrino emission
- any other non-radial instability (e.g. SASI, AAC)

generically produced by <u>any</u> CCSN

and due to <u>rotation</u> and <u>magnetic fields</u>

* [measuring the distance earth-sun with an accuracy of 1 nm]

state-of-the-art 3D simulation of a 20 solar mass model



a moderate (g=1.26 - 0.2) strangeness-dependent reduction of the axialvector coupling constant turns an unsuccesful model into a succesful one

(Melson+ '15, ApJL 808)



GW emission from postbounce phases Murphy, Ott & Burrows, 2009



- progenitors of 12, 15,
 20 and 40 solar mass
- 2D hydro (N)
- EOS: Shen+ '98
- approximate
 v heating/cooling
- parametrized explosions

distinct phases of GW emission (non-rotating cores!):

- prompt convection (negative entropy gradient left by stalling shock; generic?)
- PNS & postshock convection
- SASI (spikes from narrow downflows striking the PNS "surface")
- explosion

GW emission from postbounce phases Murphy, Ott & Burrows, 2009



GW signal from explosion phase contains information of the effects of global asymmetries in explosion

Towards realistic theoretical GW signals from CSSN

further "dimensions": rotation, magnetic fields, initial models



Some achievements/discoveries from ~40 yrs of studies of the GW signature of CCSN

- spheroidal / ellipsoidal, <u>one-zone</u> models Saenz & Shapiro '78, '79, '81

 $\Delta E/Mc^2 \sim 10^{-6} \dots 10^{-4}$

- 2D Newtonian hydro, prompt explosion models, evolutionary progenitor models, microscopic EOS Müller '82; Mönchmeyer+I '89, '91 $\Delta E/Mc^2 \sim 10^{-10} \dots 10^{-7}$
- 2D Newtonian rotating polytropes, parameter study
 Finn & Evans '90; Yamada & Sato '95; Zwerger & Müller '97;
 Rampp, Müller & Ruffert '98 (3D)
 centrifugal bounce can occur

<u>GR</u> parameter study, 2D simplified EOS, no ν-transport
 Dimmelmeier, Font & Müller '02 (CFC); Shibata '03; Shibata & Sekiguchi '05 (3D)

GR potential deeper --> larger bounce densities, more compact PNS --> centrifugal bounce occurs only rarely



- MHD models Kotake+ '04 (N); Yamada & Sawai '04 (N); Obergaulinger+ '06 (N & GR-pot); Scheidegger+ '08, '10 (3D, EOS, GR-pot),

only strong initial fields ($B > 10^{11} G$) influence overall dynamics

---> GW amplitude affected by

* magnetic fields that contribute significantly to total energy density

* bipolor MHD driven jet outflow giving rise to a new signal type

* initially strong toroidal field: jet suppressed by fast growing spiral SASI





- <u>GW from neutrinos</u> (time-dependent asymmetric energy/mass flux) Müller & Janka '97, Müller+ '04 (2D, GR-pot)
 - * frequencies lower (\leq 30 Hz) than those of matter signal (100 Hz ... 1000 Hz) \rightarrow harder to observe because of seismic detector noise

* amplitudes overall larger





non-rotating 11 solar mass 2D model

rotating 15 solar mass 2D model

Parametrized 3D models of neutrino-driven core collapse supernovae Müller, Janka, Wongwathanarat (2012)



GW amplitudes due to aspherical flow & corresponding spectograms $dE_{_M}/d\nu$

total GW amplitudes (including ν)

GW signature of parametrized 3D models



Rh^{TT}_{*} - 32.5 cm 35.9 cm



B.Müller, Janka & Marek 2013: 2D, CFC, multi-group three-flavor ray-by-ray-plus v-transport

relativistic effects in combination with detailed v-transport essential for quantitative predictions, determine structure of PNS surface layer and its characteristic g-mode frequency



solid lines: matter GW signal; dashed lines: neutrino GW signal (different scales!)

Cerda-Duran+ 2013: 2D, XCFC, 35 solar mass, rapidly rotating (2rad/s), low-metallicity progenitor, gray v-leakage scheme, LS220-EOS,

two models simulated until BH formation (M_{PNS} > 2.04 M_{SUP})



lines show frequency evolution of g-modes at PNS surface (green), g-modes in cold inner core (solid red), quasi-radial mode (dashed red), and f-mode (dotted blue)

Kuroda, Takiwaki & Kotake 2014:

3D, **GR**, 15 solar mass progenitor, 5 rotating models (Ω [rad/sec) = $\pi/6$, $\pi/3$, and π), Shen EOS, approximate v-transport (M1 closure, variable Eddington factor of Levermore 1984)

- results consistent with previous work in 3D (Ott+ '07, '12, '13; simpler transport, softer EOS)
- signals from prompt convection qualitatively similar (except for most rapidly rotating model)
- nonaxisymmetric instabilities (spiral SASI) essential for GW signature of rotating models (see also 3D Newtonian study with ZEUS-MP of Kotake+ '11)



Hayama+ 2016: circular polarization of GW from CCSN provide a clear indication of rapid rotation rotation period < few seconds \rightarrow galactic events detectable with network of 2nd generation detectors

Andresen, B.Müller, E.Müller & Janka 2016:

3D, **GR-pot**, **11**.2, 20, and 27 solar mass progenitors, LS-220 EOS, ray-by-ray-plus (Buras+ '06) energy-dependent two-moment multi-flavor v-transport (VERTEX, Rampp & Janka '02)

See also: Kuroda, Kotake & Takiwaki '16, 15 solar mass model, different nuclear EOS, grey ν -transport

- GW emission (in pre-explosion phase) strongly depends on whether post-shock flow is dominated by SASI or convection (driven by neutrino heating)

- SASI-dominated models:

strong emission at low frequencies around 100 ... 200 Hz; not present in 2D models!

but low-frequency emission is no unambiguous signature of SASI, can also also occur after onset of explosion

- GW emission of 3D models differs considerably from that of 2D models amplitudes: $|A_{+,x}| \leq 4$ cm in 3D models, and ~ few 10 cm in 2D models
- shock revival in exploding 20 M_{\odot} model results in enhanced low-frequency emission (preferred scale of convective eddies in PNS convection zone changes)

low-frequency SASI emission GW from prompt convection

Andresen, B.Müller, E.Müller & Janka 2016



Andresen, B.Müller, E.Müller & Janka 2016:

to determine origin of high-frequency and low-frequency components of GW signal the computational volume is divided into three layers

layer A: PNS convection zone

layer B: PNS convectively stable surface layer

layer C: region beyond the PNS"surface"



low-frequency emission: from all 3 layers (even for model s20s!) high-frequency emission: mostly from aspherical motion in layer A



Andresen, B.Müller, E.Müller & Janka 2016:

- SASI produces strong signal component in the frequency range 100 ... 200 Hz

amplitude spectrograms of low-frequency GW signal arising from the three different layers



apparent temporal correlation of low-frequency emission with the SASI, because of a global modulation $(3D \neq 2D!)$ of the accretion flow by the SASI

GW signal traces frequency of I=1,2 SASI modes, but with a frequency doubling for the I=1 spiral mode (integral in d^2Q/dt^2 invariant to a rotation by π in any direction) \rightarrow double-peak structure of low-frequency GW signal

OUTLOOK:

The detection of the gravitational wave signal from a CCSN will provide new & independent insights into the explosion dynamics

However, detection prospects appear rather bleak: $D \leq 3$ kpc (2nd-generation), $D \leq 30$ kpc (3rd-generation) & closer events with high S/N

Predicting the gravitational wave signal of CCSN: 3 decades of fun & ordeal

Müller 1982



Müller+ 2012, Andresen+ 2016



