



**inflation**  
**@landscape**

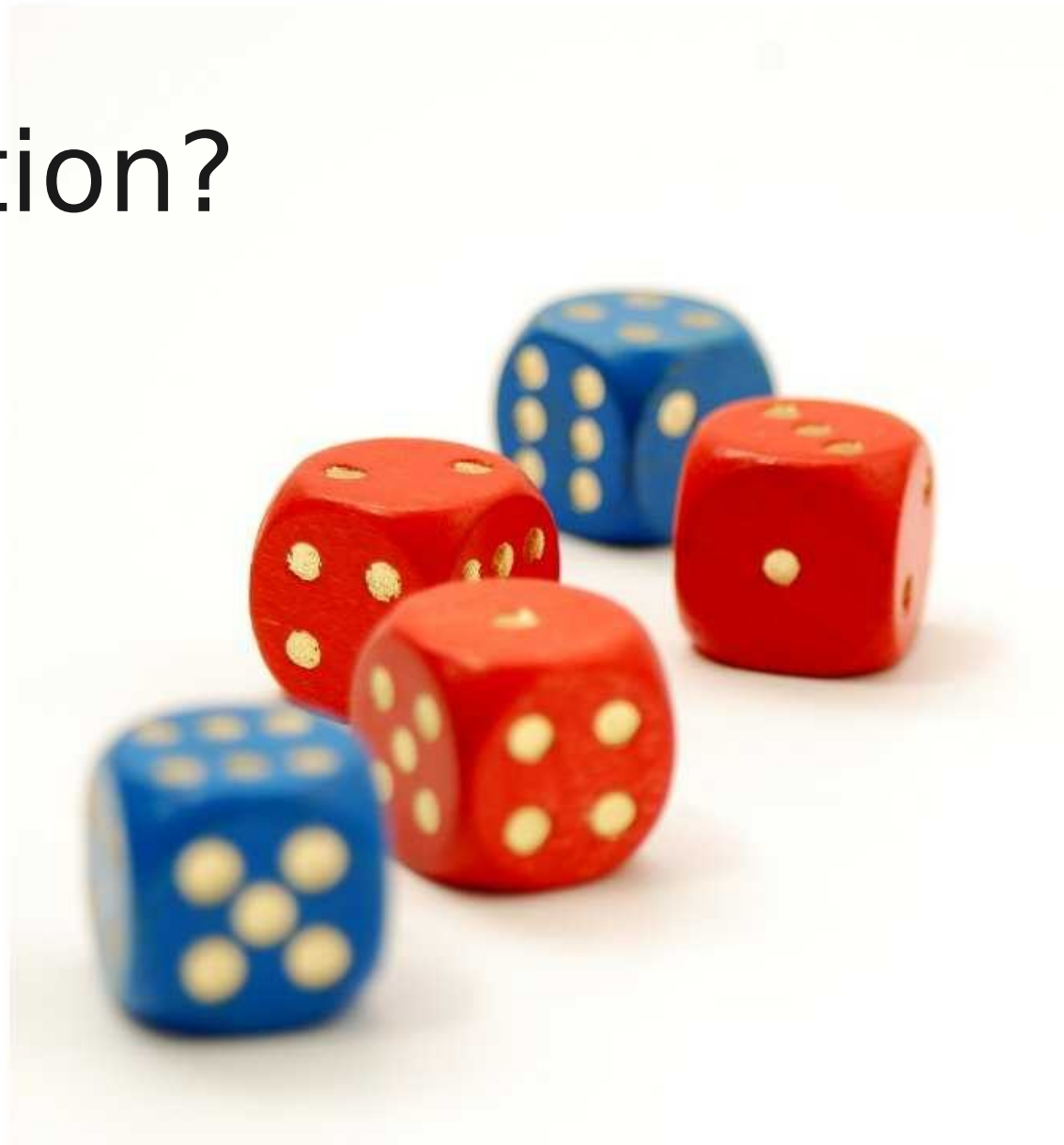
# Inflation in a Landscape

Yi Wang, McGill University, June 2011

Based on:

- X. Chen, YW, 0909.0496, 0911.3380
- M. Li, YW 0903.2123
- N. Afshordi, A. Slosar, YW, 1006.5021
- Y. Cai, S. Pi, YW, 110?.\*
- R. Brandenberger, F. Duplessis, YW, 110?.\*
- J. Cline, G. Moore, YW, 1106.2188

Bet on inflation?

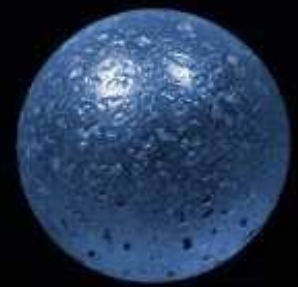




multiple fields

or

single field



The potential is

flat

or

bumpy



Dynamics: Simple,

or

Complicated?

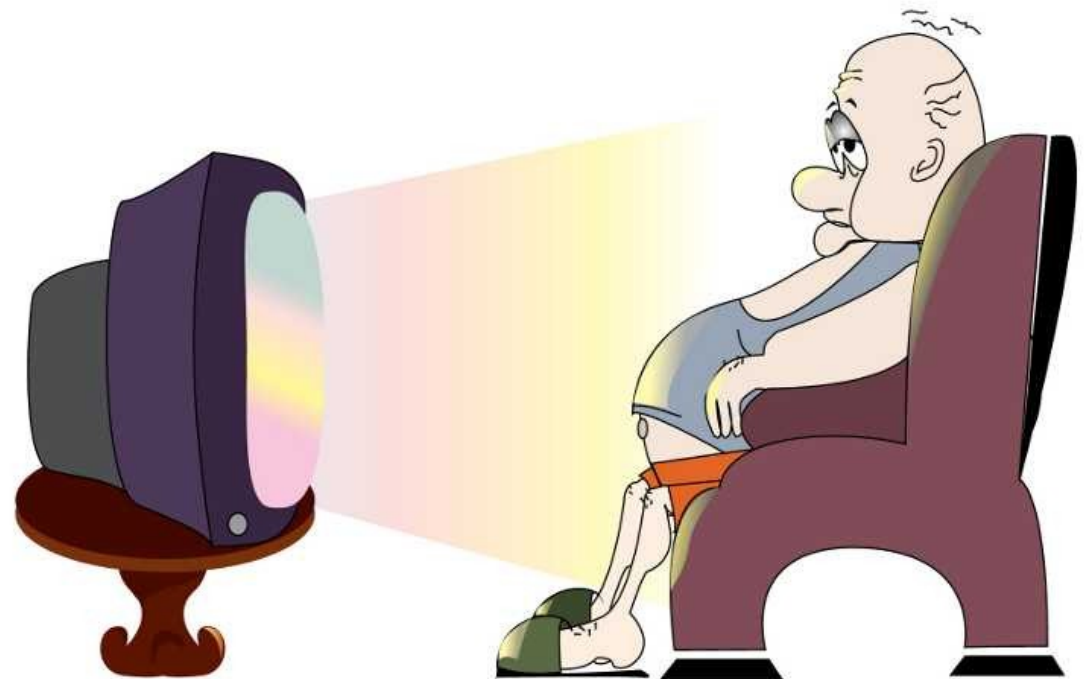


Alternatives ...





work hard or  
go home & watch TV

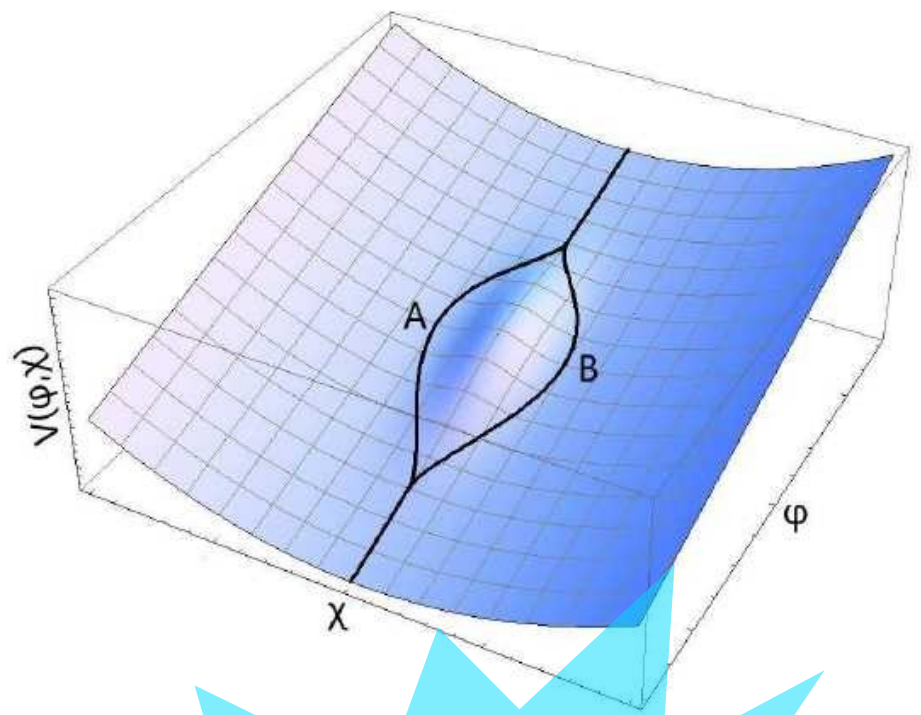
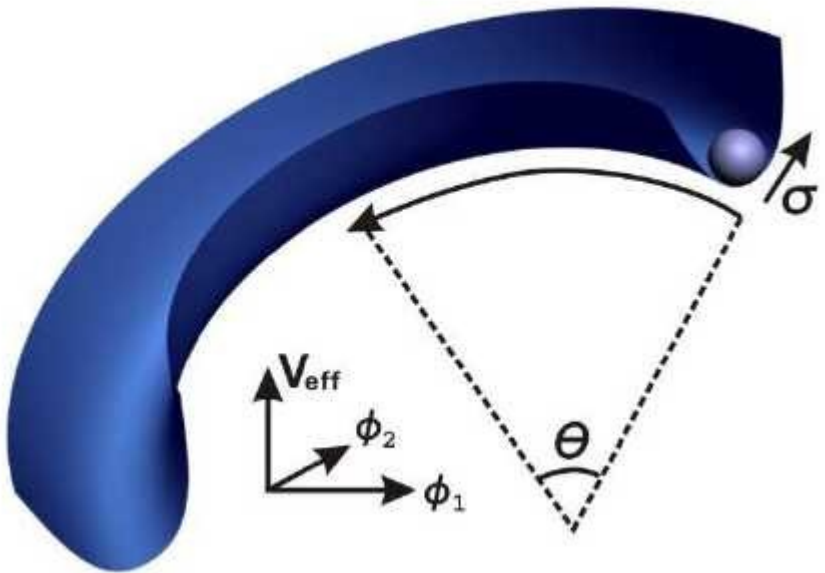






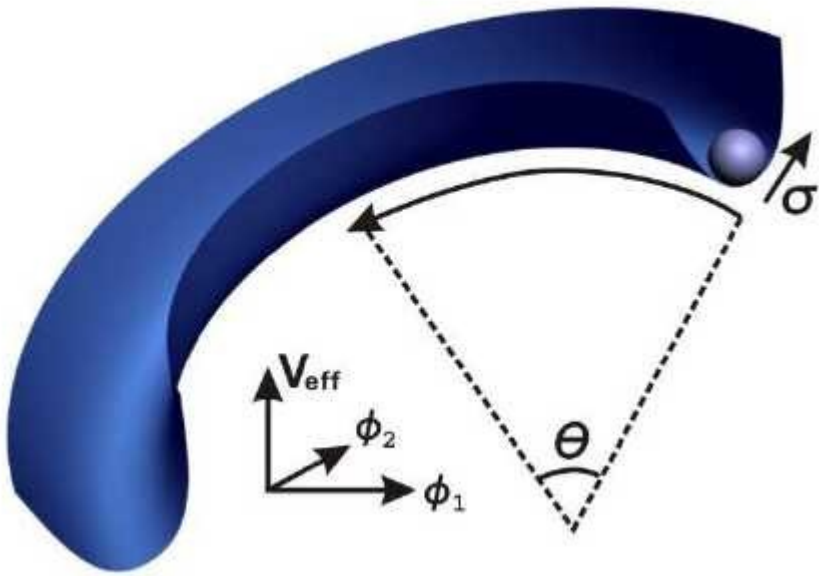
work hard





0	0	0	0	1	1	2	2	3	3	3	2	2	1	1	0	0
0	0	0	0	0	1	1	2	2	3	2	2	1	1	0	0	0
0	0	0	0	0	0	1	1	2	1	2	1	1	0	0	0	0
0	0	0	0	0	0	0	1	0	1	1	1	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	1	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0

**BET**



Part I: quasi-single field inflation

Definition of  
single field?





Background:

$$m^2 < \eta H^2$$

Perturbations:

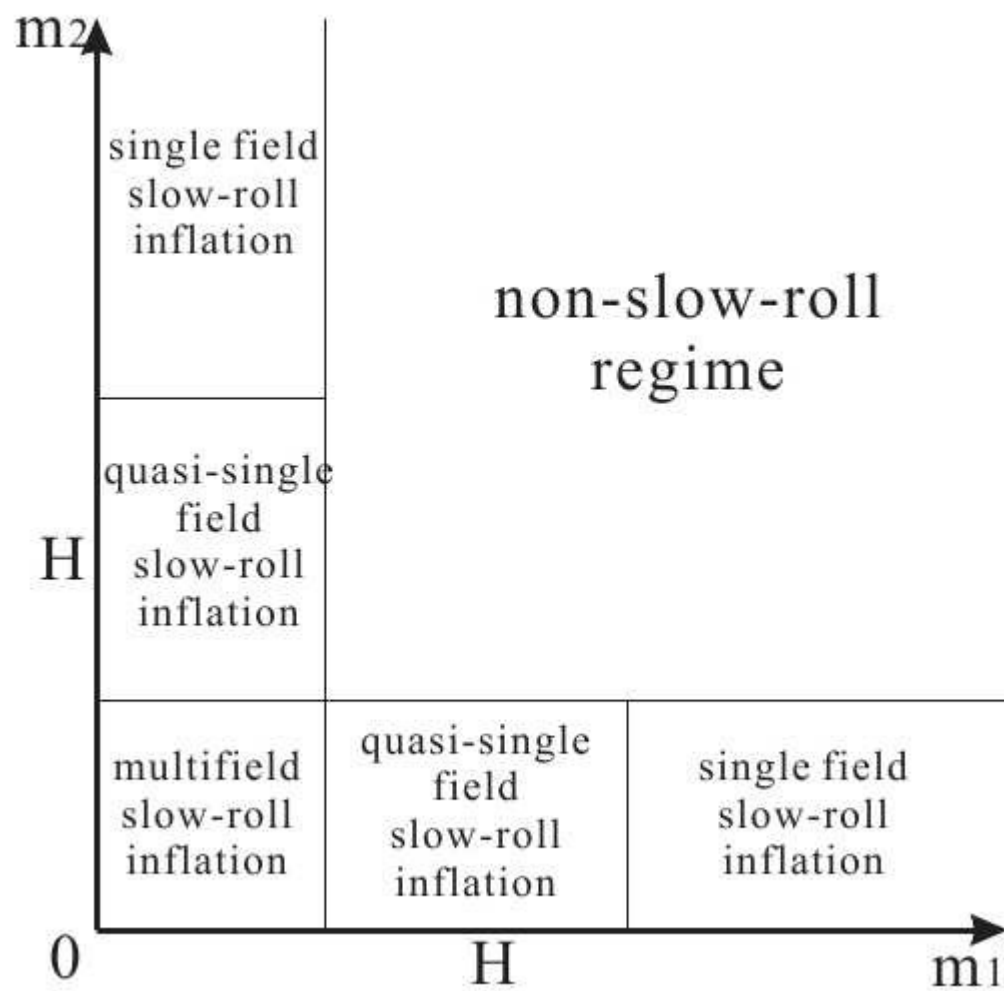
$$m < H$$



Quasi-single field:

$$\eta H^2 < m^2 < H^2$$





Why to care  $\eta H^2 < m^2 < H^2$  ?

*a priori* large parameter space!



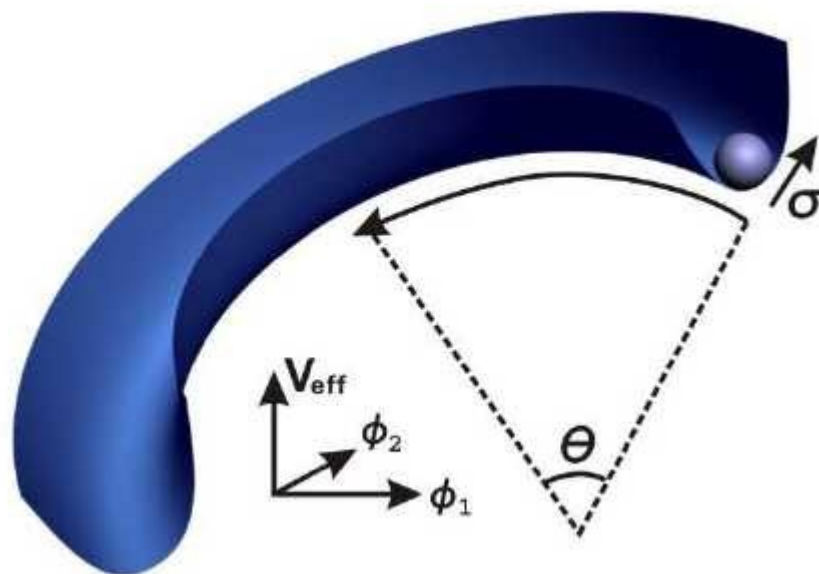
Why to care  $\eta H^2 < m^2 < H^2$  ?

*a priori* large parameter space!

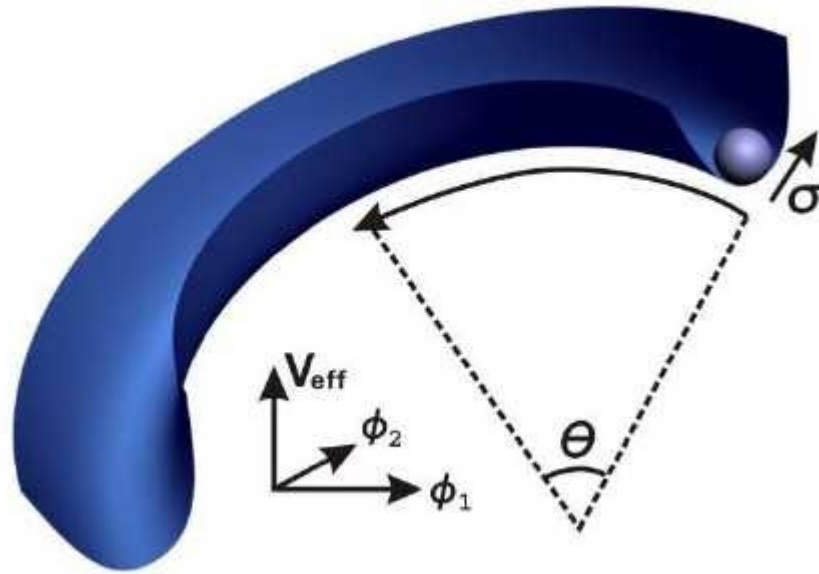
Fine tuning problem of inflation:

To tune once is already a lot

A simple model:



# A simple model:



$$S_m = \int d^4x \sqrt{-g} \left[ -\frac{1}{2}(\tilde{R} + \sigma)^2 g^{\mu\nu} \partial_\mu \theta \partial_\nu \theta - \frac{1}{2} g^{\mu\nu} \partial_\mu \sigma \partial_\nu \sigma - V_{\text{sr}}(\theta) - V(\sigma) \right]$$

# Method of calculation:

## in-in formalism

$$\langle Q(t) \rangle \equiv \langle 0 | \left[ \bar{T} \exp \left( i \int_{t_0}^t dt' H_I(t') \right) \right] Q_I(t) \left[ T \exp \left( -i \int_{t_0}^t dt' H_I(t') \right) \right] | 0 \rangle$$

$$i^n \int_{t_0}^t dt_1 \int_{t_1}^t dt_2 \cdots \int_{t_{n-1}}^t dt_n \langle [H_I(t_1), [H_I(t_2), \cdots, [H_I(t_n), Q_I(t)] \cdots]] \rangle$$



# Transfer vertex

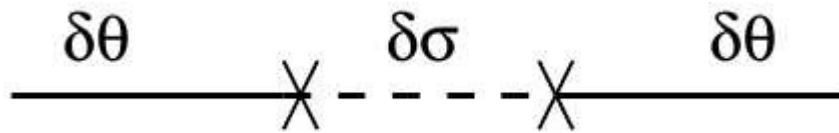
$$S_m = \int d^4x \sqrt{-g} \left[ -\frac{1}{2} (\tilde{R} + \sigma)^2 g^{\mu\nu} \partial_\mu \theta \partial_\nu \theta - \frac{1}{2} g^{\mu\nu} \partial_\mu \sigma \partial_\nu \sigma - V_{\text{sr}}(\theta) - V(\sigma) \right]$$



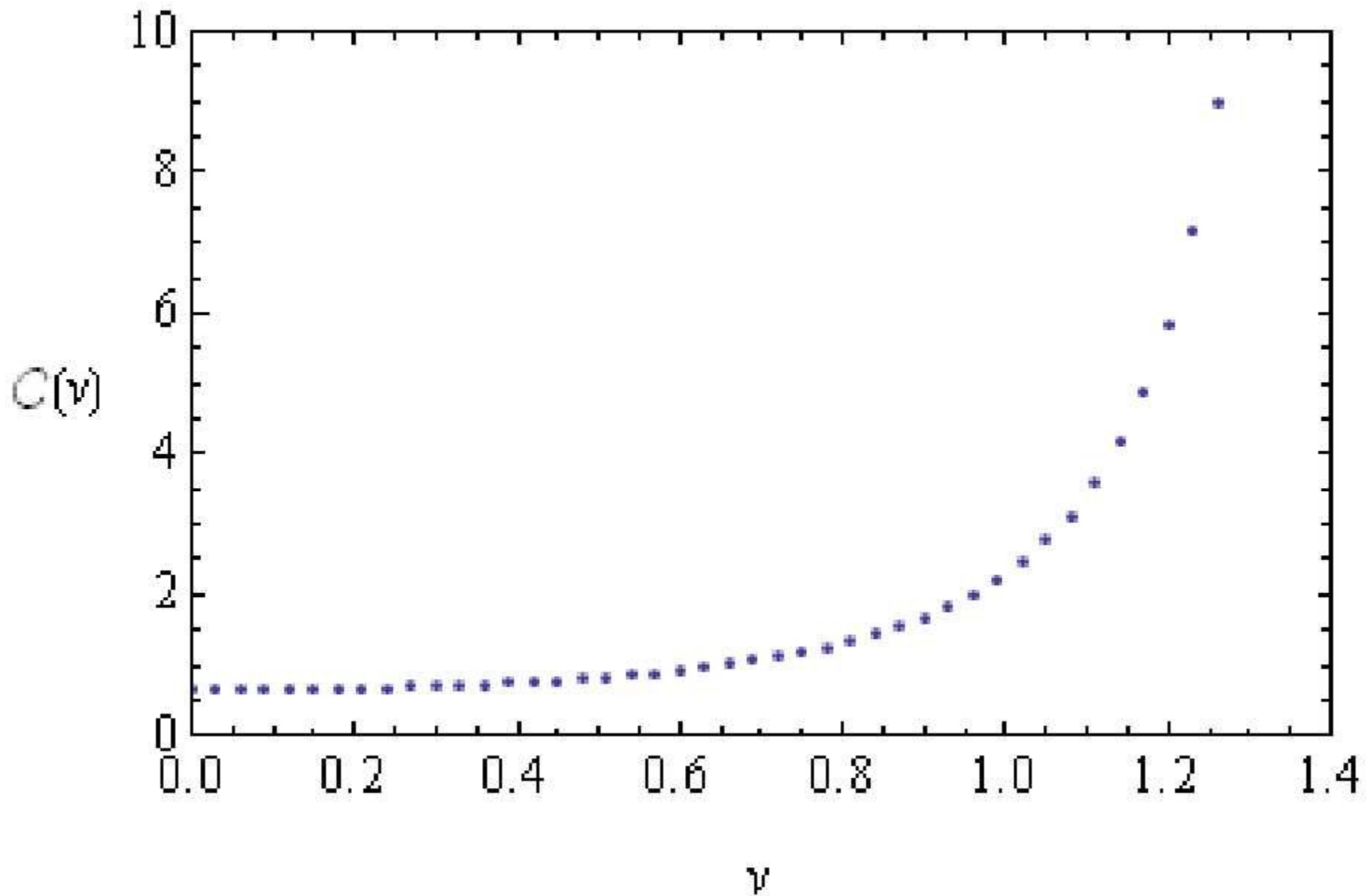
2pt coupling from turning trajectory

$$\delta \mathcal{L}_2 = 2a^3 R \dot{\theta}_0 \delta \sigma \delta \dot{\theta} \quad \begin{array}{c} \delta \theta \\ \text{---} \times \text{---} \delta \sigma \\ \dot{\theta} / H \end{array}$$

# Power spectrum



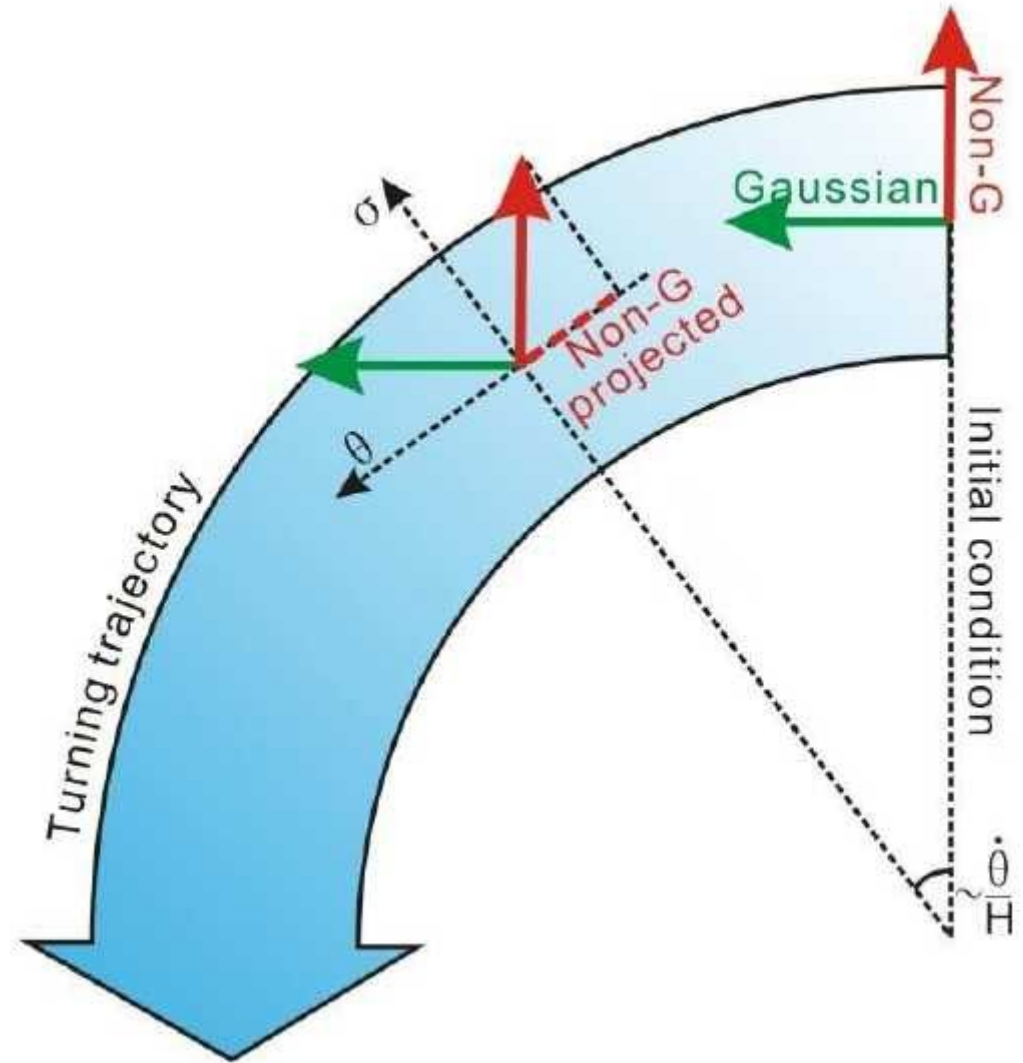
$$\delta P_\zeta \sim \left(\dot{\theta}/H\right)^2 P_\zeta$$



$$\delta P_\zeta = C(\nu) \left( \dot{\theta}/H \right)^2 P_\zeta$$

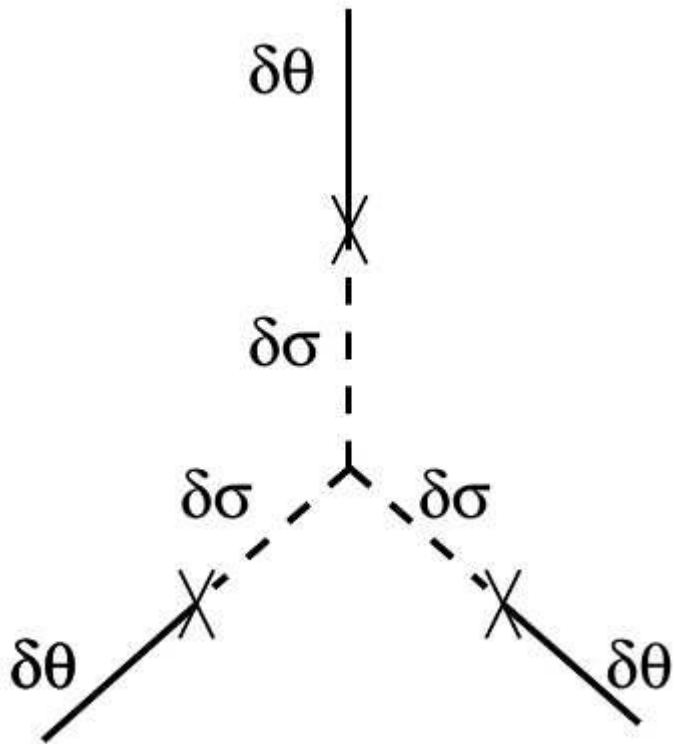
$$\nu = \sqrt{9/4 - m^2/H^2}$$

# Non-Gaussianity





# Bispectrum: amplitude



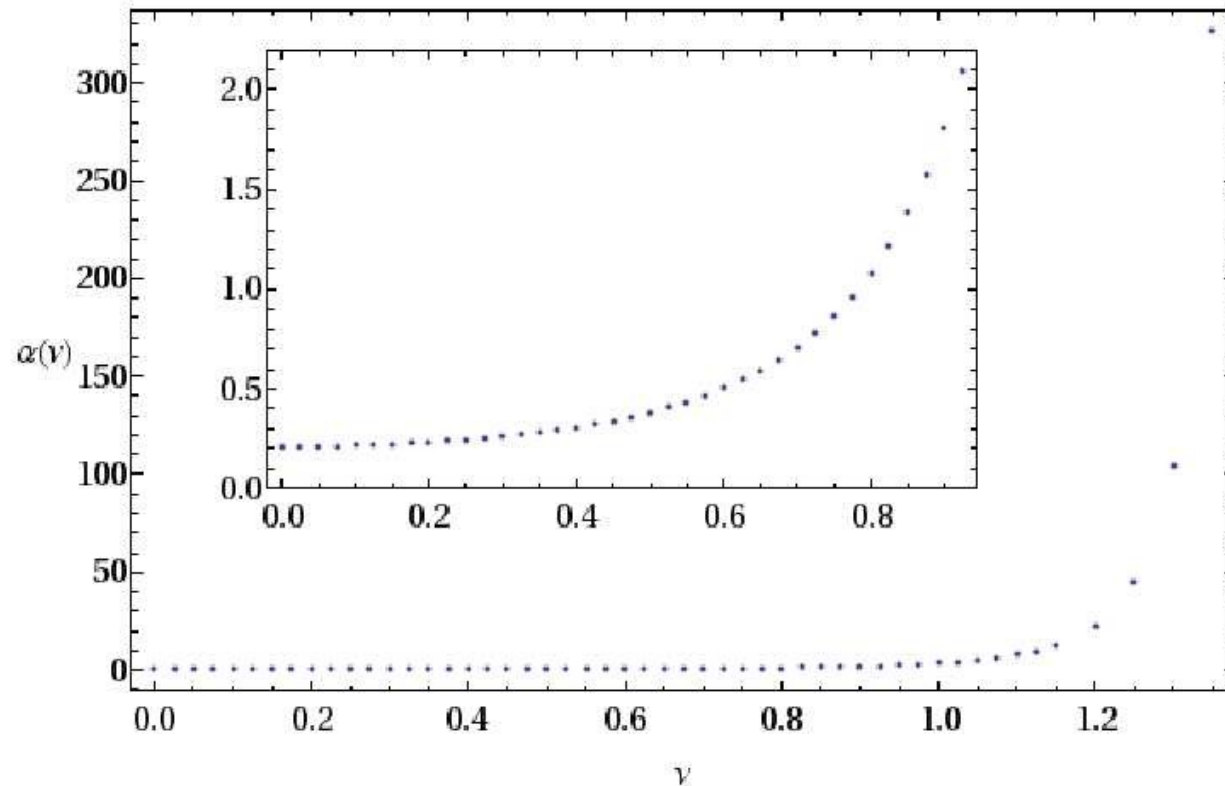
$$f_{NL} \sim P_{\zeta}^{-1/2} \left( \dot{\theta} / H \right)^3 (V''' / H)$$



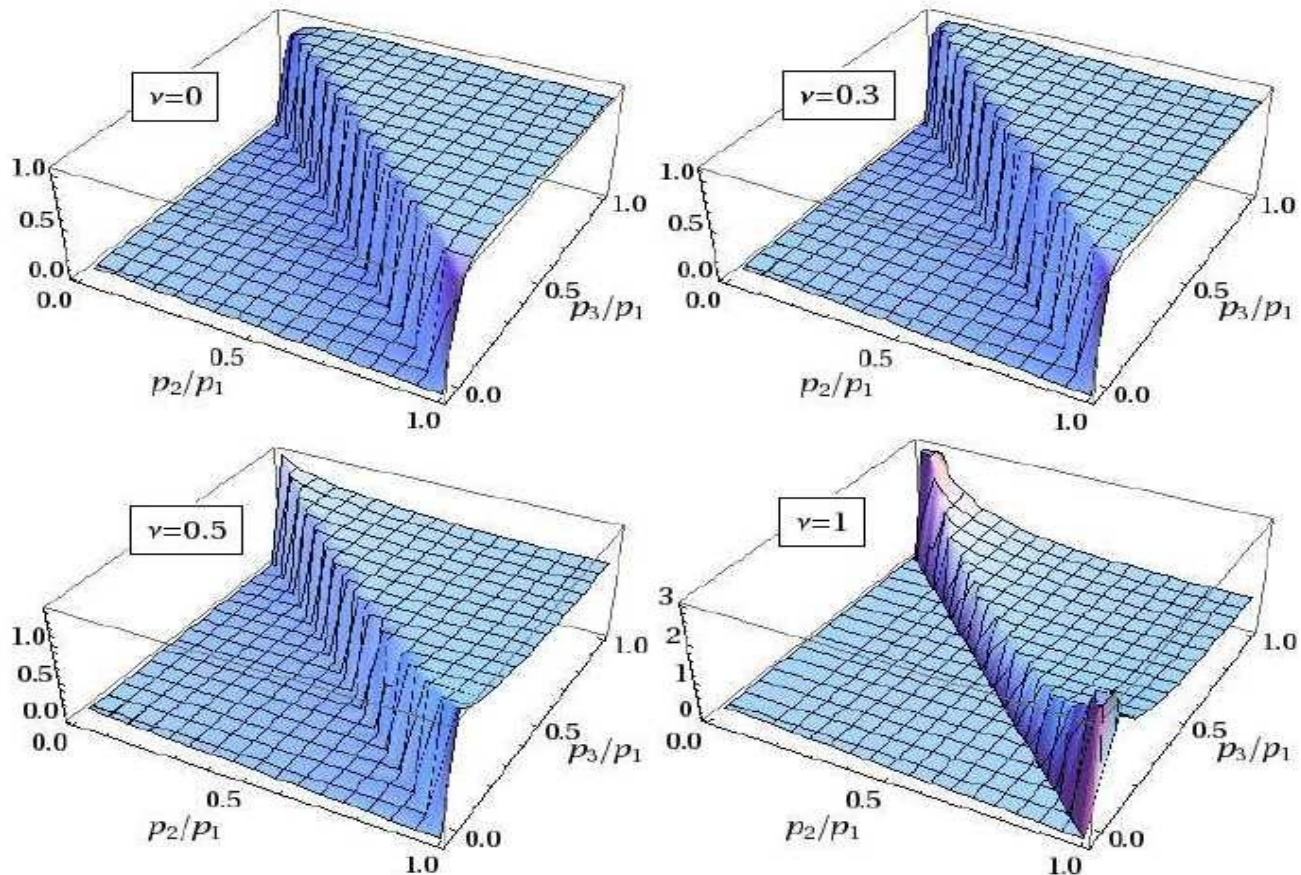
$$\begin{aligned} \langle \zeta^3 \rangle &\sim \text{coupling} \times P_{\zeta}^{3/2} \\ &\sim f_{NL} \times P_{\zeta}^2 \end{aligned}$$

# Bispectrum: amplitude

$$f_{NL} = \alpha(\nu) P_{\zeta}^{-1/2} \left( \dot{\theta}/H \right)^3 (V'''/H)$$



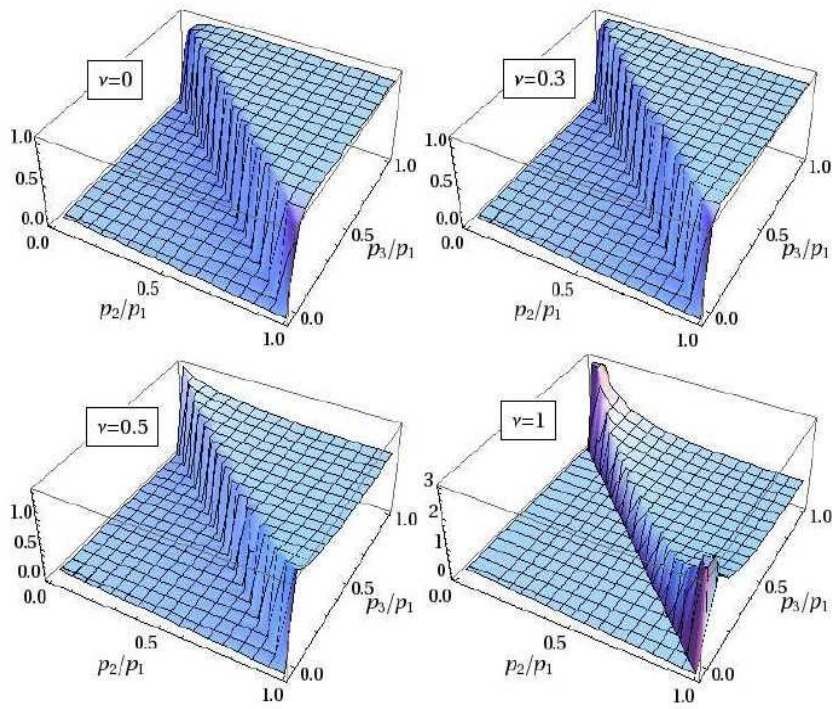
# Bispectrum: shape



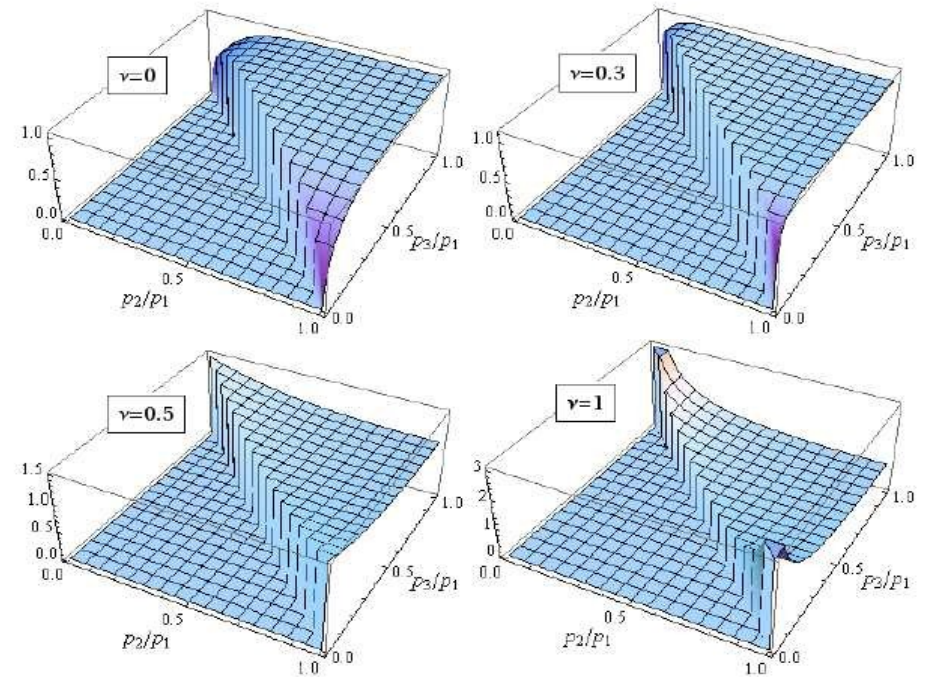
$$\nu = \sqrt{9/4 - m^2/H^2}$$

# Bispectrum: shape ansatz

$$F = \frac{3^{\frac{9}{2}-3\nu}}{10} \frac{f_{NL}^{\text{int}}(p_1^2 + p_2^2 + p_3^2)}{(p_1 p_2 p_3)^{\frac{3}{2}+\nu} (p_1 + p_2 + p_3)^{\frac{7}{2}-3\nu}}$$



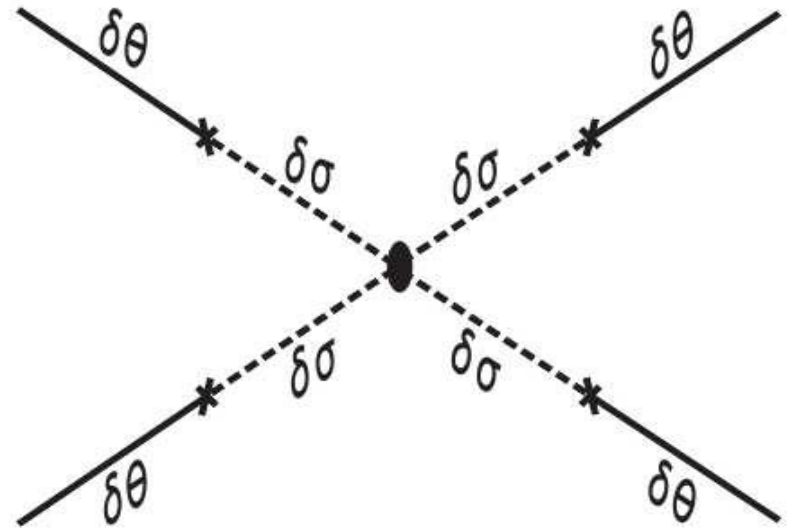
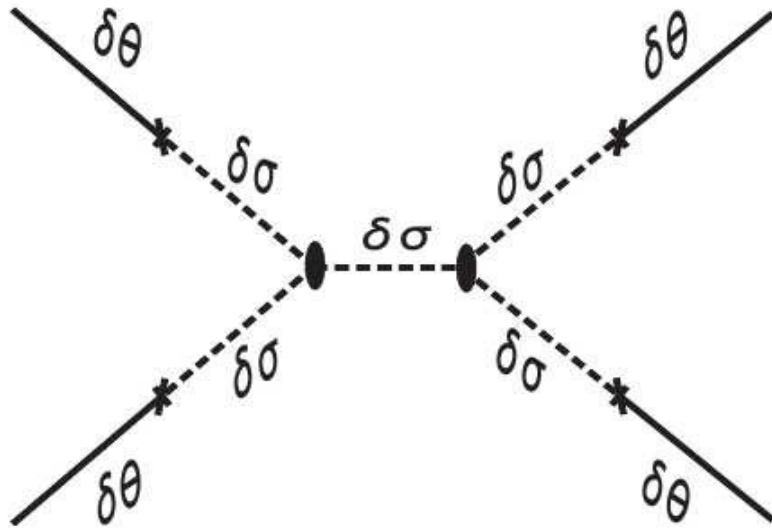
True numeric



Shape ansatz



# Trispectrum



$$t_{NL} \sim \max \left\{ P_{\zeta}^{-1} \left( \dot{\theta}/H \right)^4 (V''''/H)^2, P_{\zeta}^{-1} \left( \dot{\theta}/H \right)^4 V'''' \right\}$$

$$t_{NL} \gg f_{NL}^2 \text{ for } \dot{\theta}/H \ll 1$$



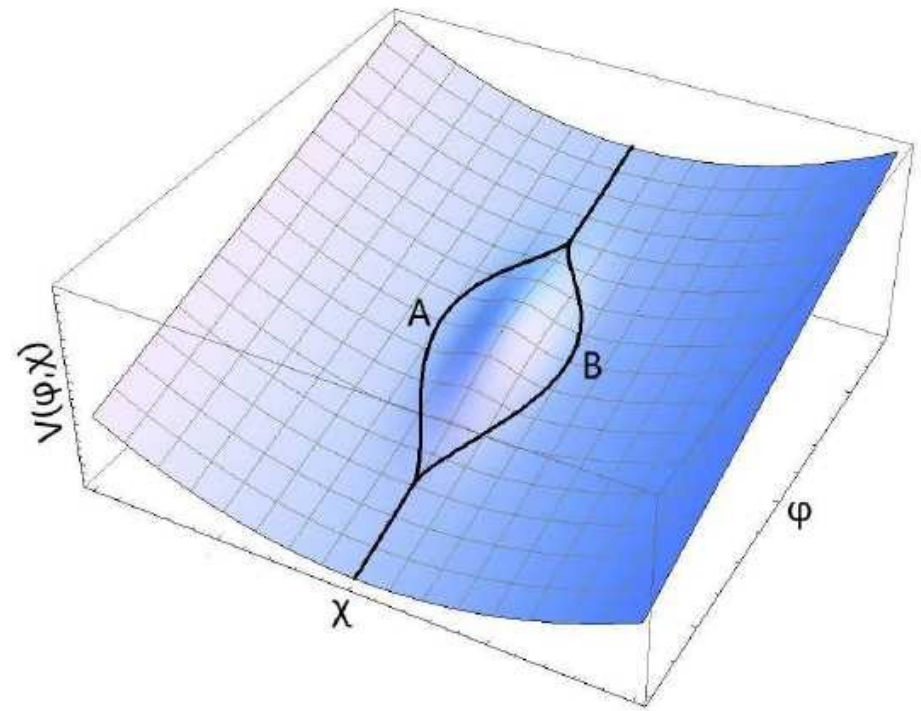
# Higher point correlations

$$h_{NL} \sim P_{\zeta}^{-3/2} \left( \dot{\theta}/H \right)^5 (V'''/H)^3 \sim (\dot{\theta}/H)^{-4} f_{NL}^3 ,$$

$$i_{NL} \sim P_{\zeta}^{-2} \left( \dot{\theta}/H \right)^6 (V'''/H)^4 \sim (\dot{\theta}/H)^{-6} f_{NL}^4 .$$

Might be difficult to probe.

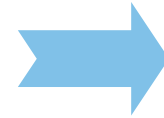
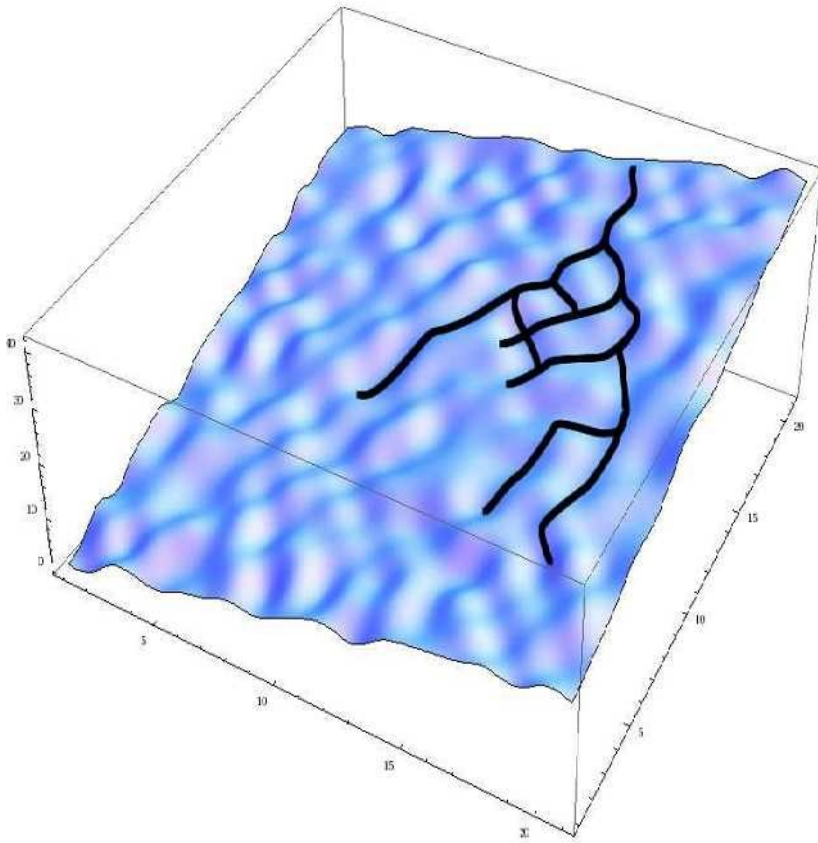
But ... who knows?



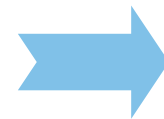
## Part II: multi-stream inflation



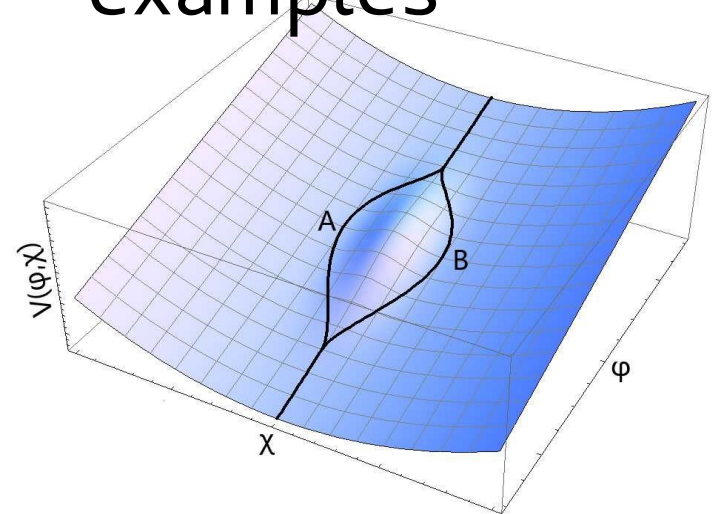
# Multi-stream trajectories for inflation



statistics

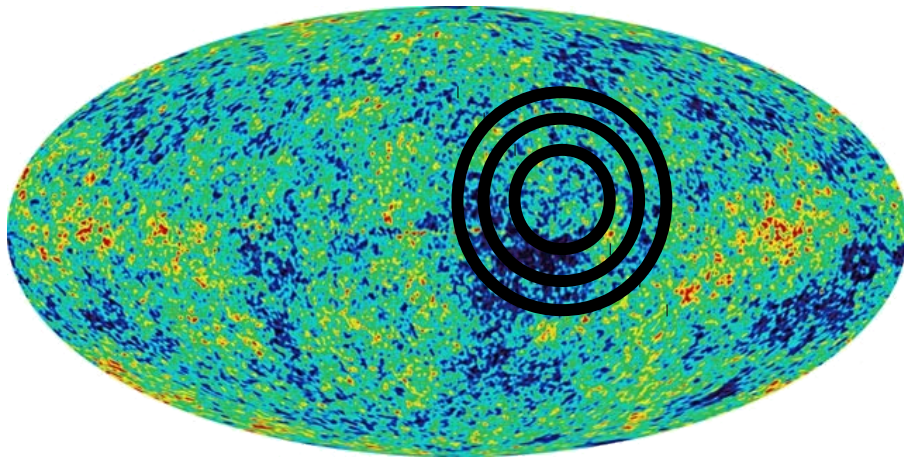
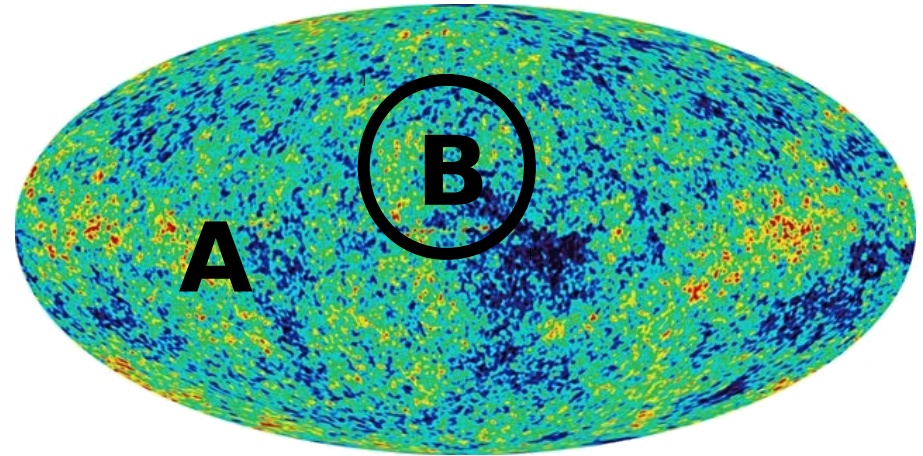
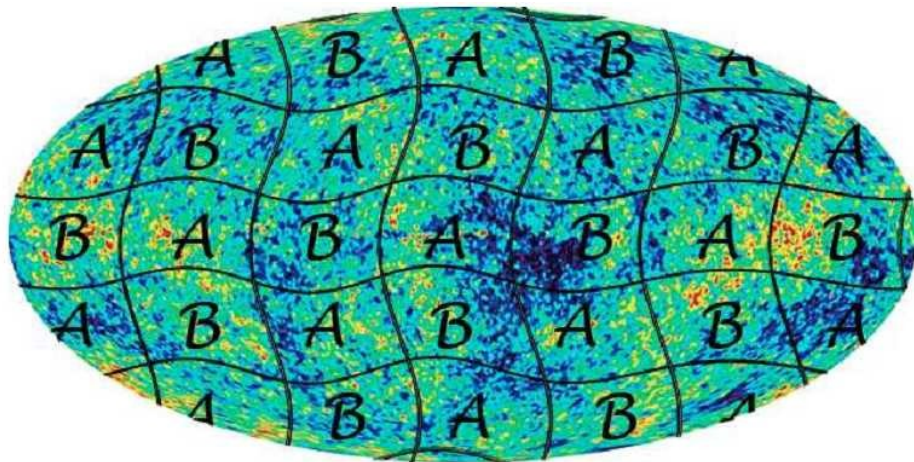


examples

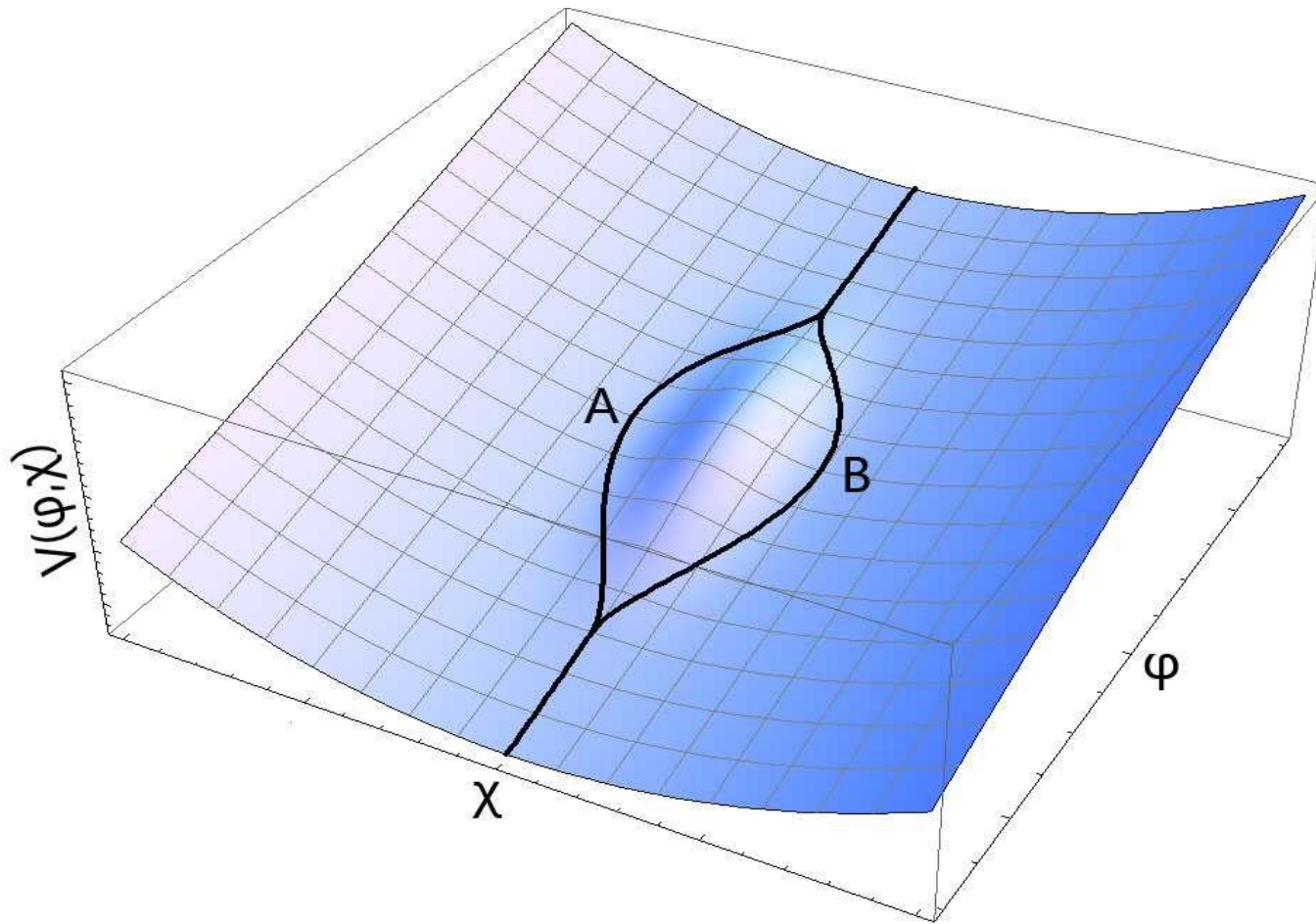




examples (case study):



# Nearly symmetric bifurcations

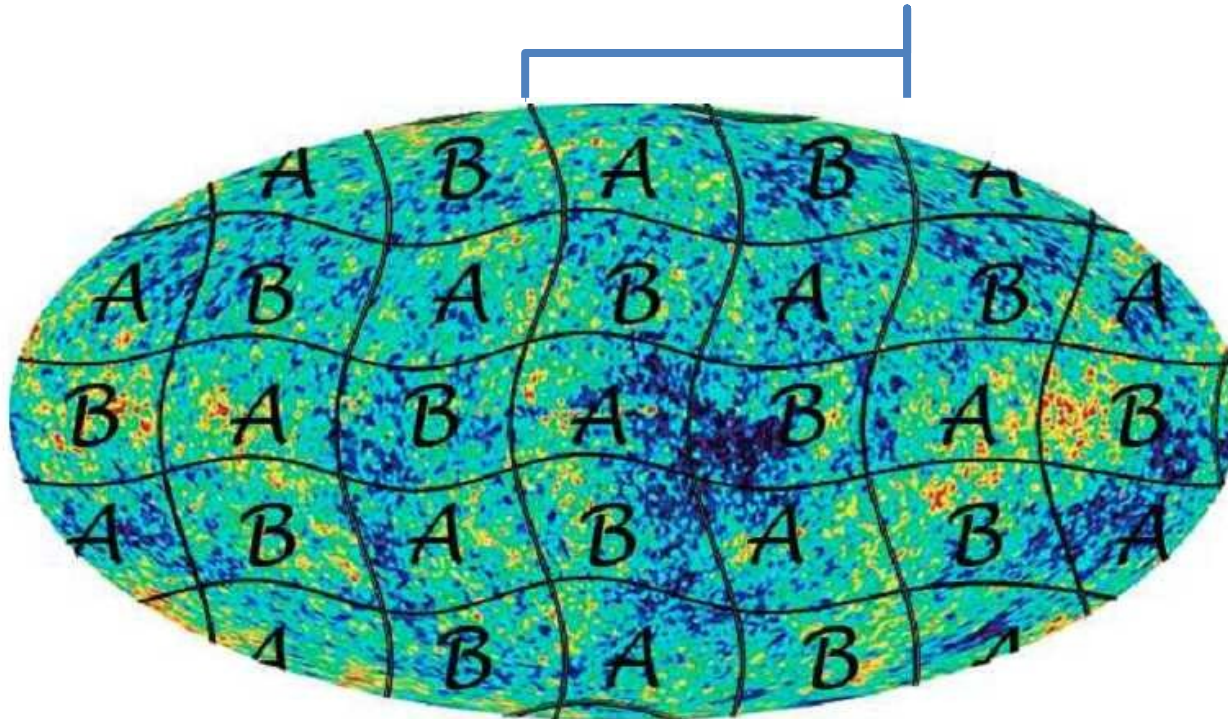




# Nearly symmetric bifurcations

----- Pert. feature

bifurcation scale

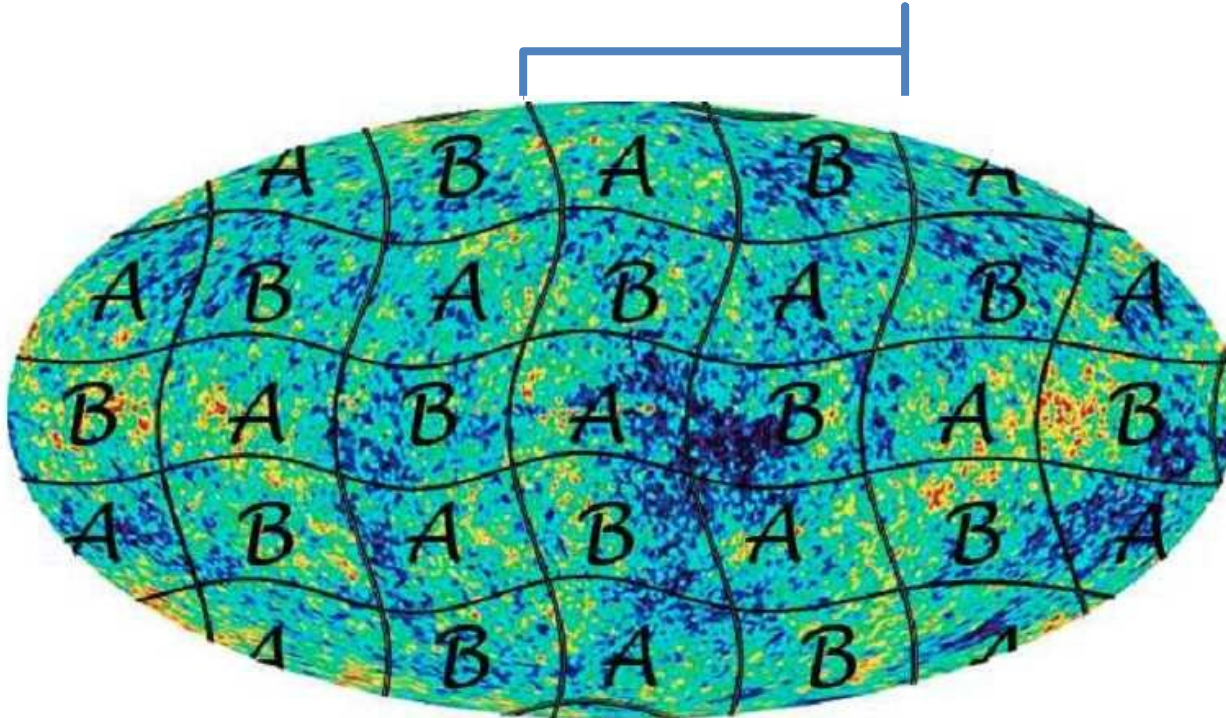


$$\zeta = \delta N$$

# Nearly symmetric bifurcations

----- Pert. asym.

bifurcation scale



$$P_{\zeta}(A) \neq P_{\zeta}(B)$$

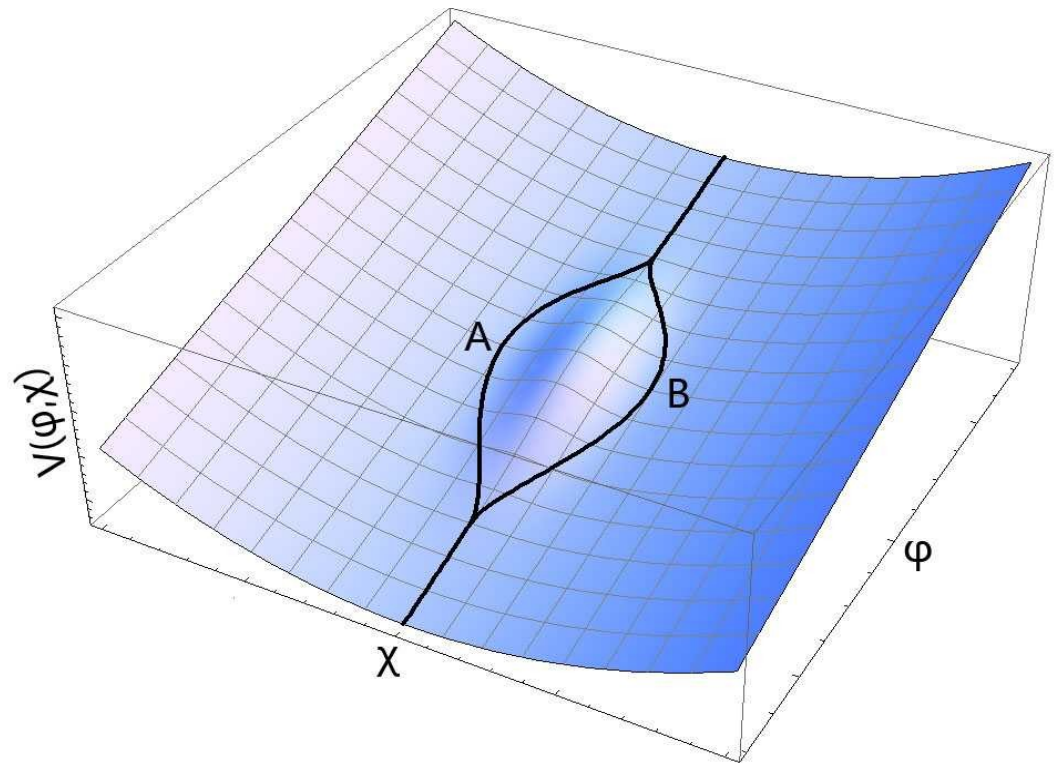
# Nearly symmetric bifurcations

----- Non-G

Correlation between

feature

asymmetry



# Nearly symmetric bifurcations

----- Non-G

$$P(\delta\zeta_{k_1}^{S,N}, \zeta_k) = P(\delta\zeta_{k_1}^{S,N}) \left[ \frac{e^{-\frac{\zeta_k^2}{2\sigma_A^2}}}{\sqrt{2\pi}\sigma_A} \theta(\delta\zeta_{k_1}^S) + \frac{e^{-\frac{\zeta_k^2}{2\sigma_B^2}}}{\sqrt{2\pi}\sigma_B} \theta(-\delta\zeta_{k_1}^S) \right]$$

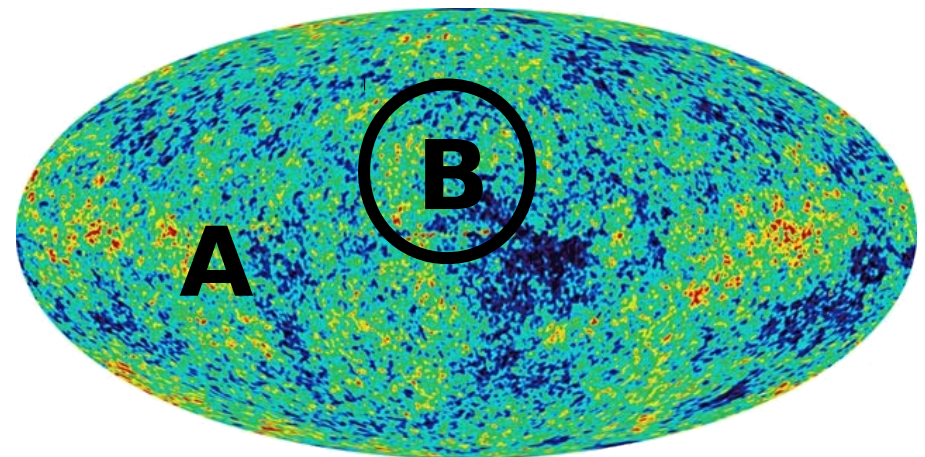
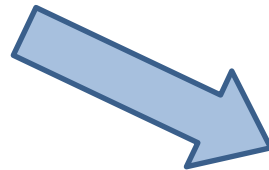
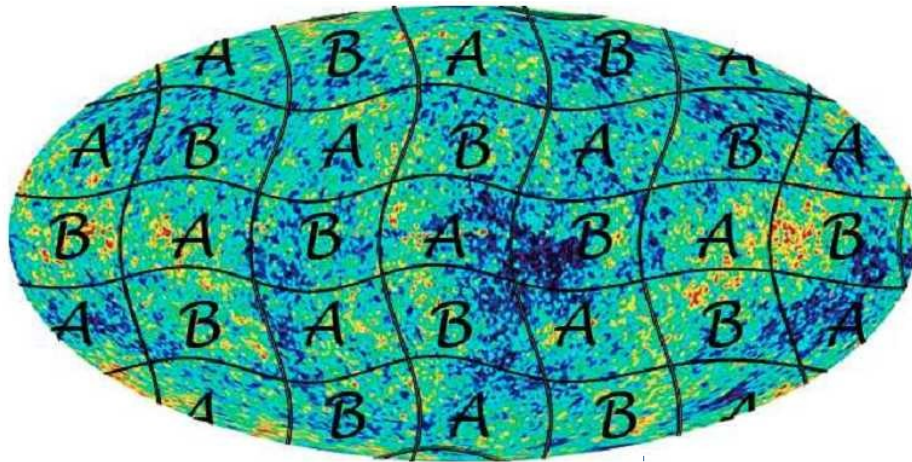
$$f_{NL} \simeq x P_\zeta^{-1/2} \left( \frac{P_\zeta^A - P_\zeta^B}{P_\zeta} \right)$$

$x \equiv \delta\zeta_{k_1}/\zeta_{k_1}$  denotes the fraction of extra fluctuation from the multi-stream effect.



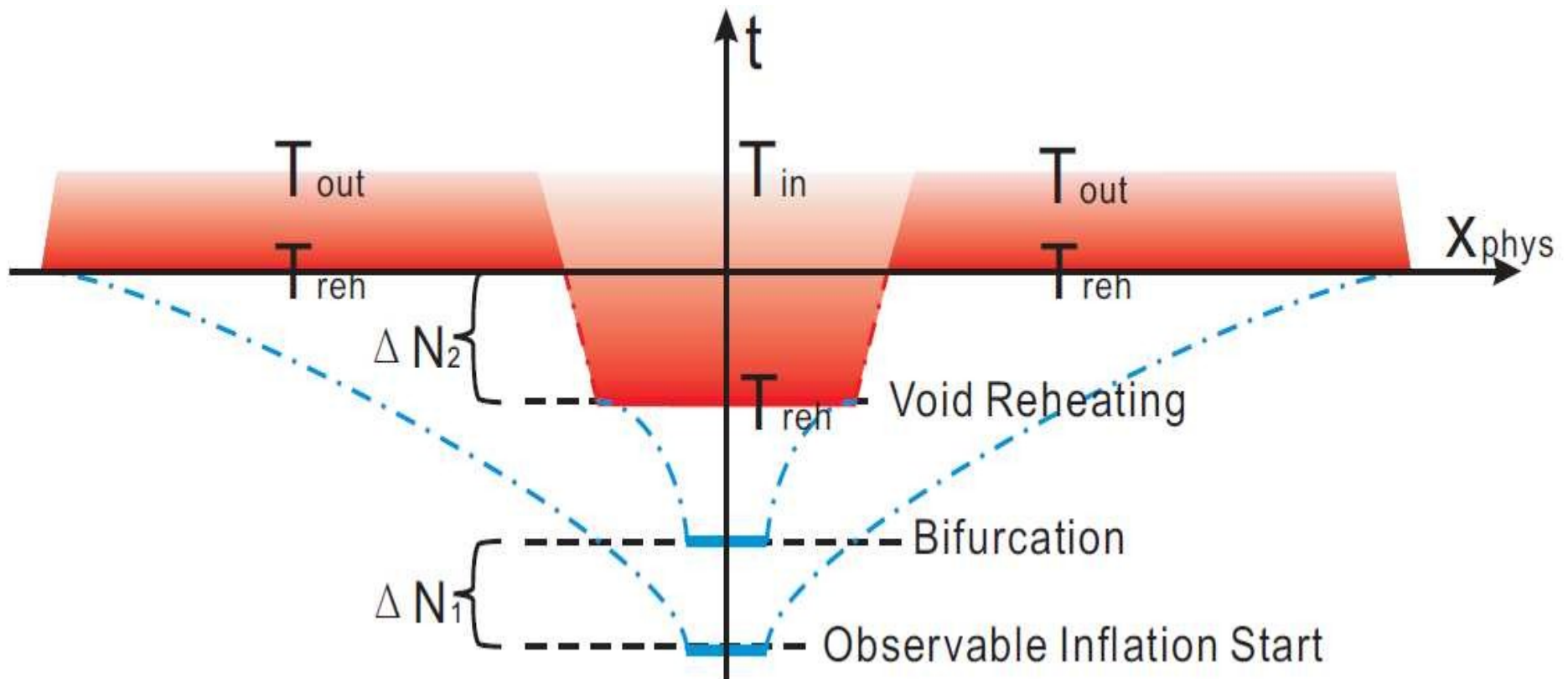
# Very asymmetric bifurcations

---- Overview

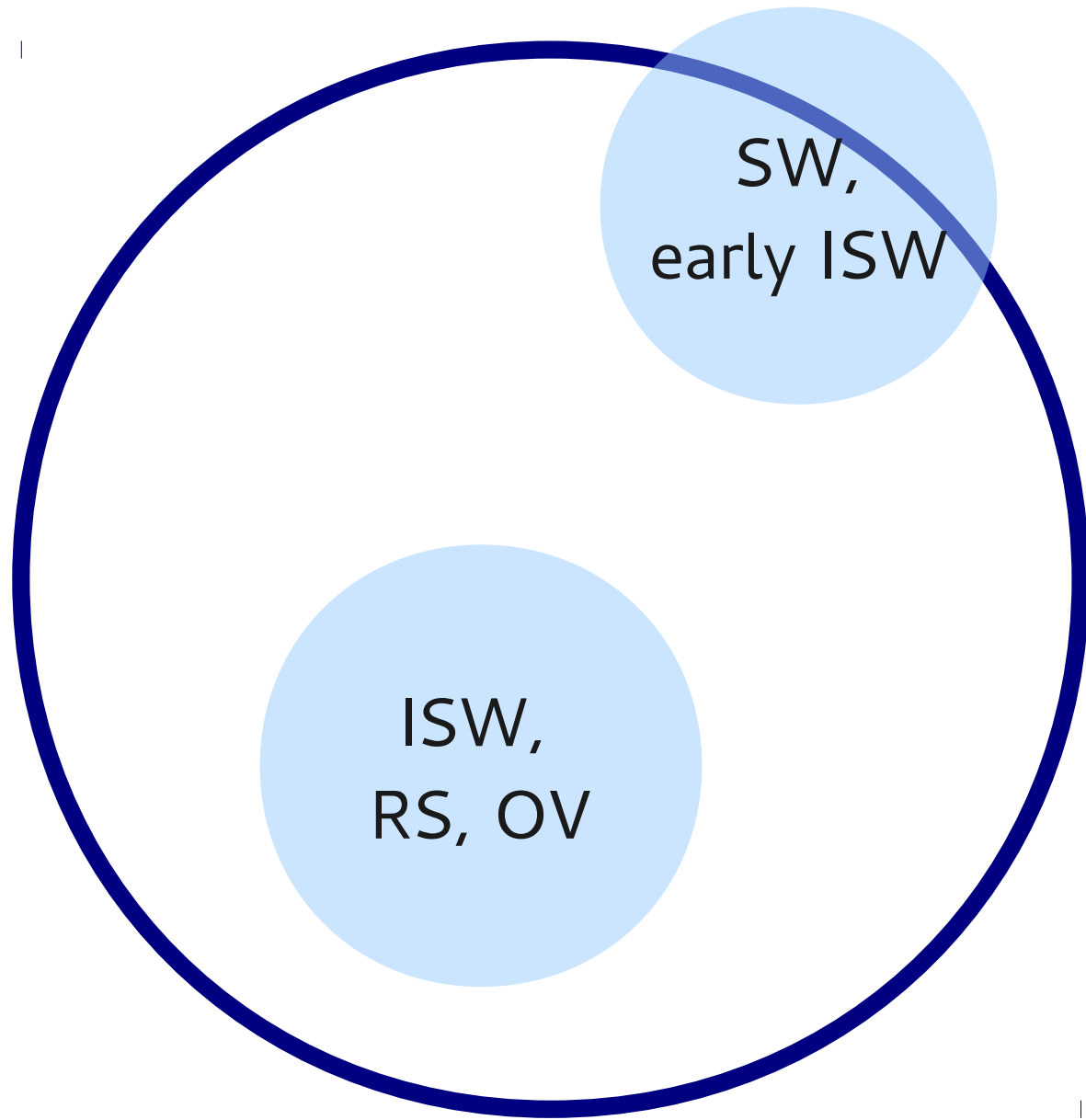


# Very asymmetric bifurcations

----- Primordial



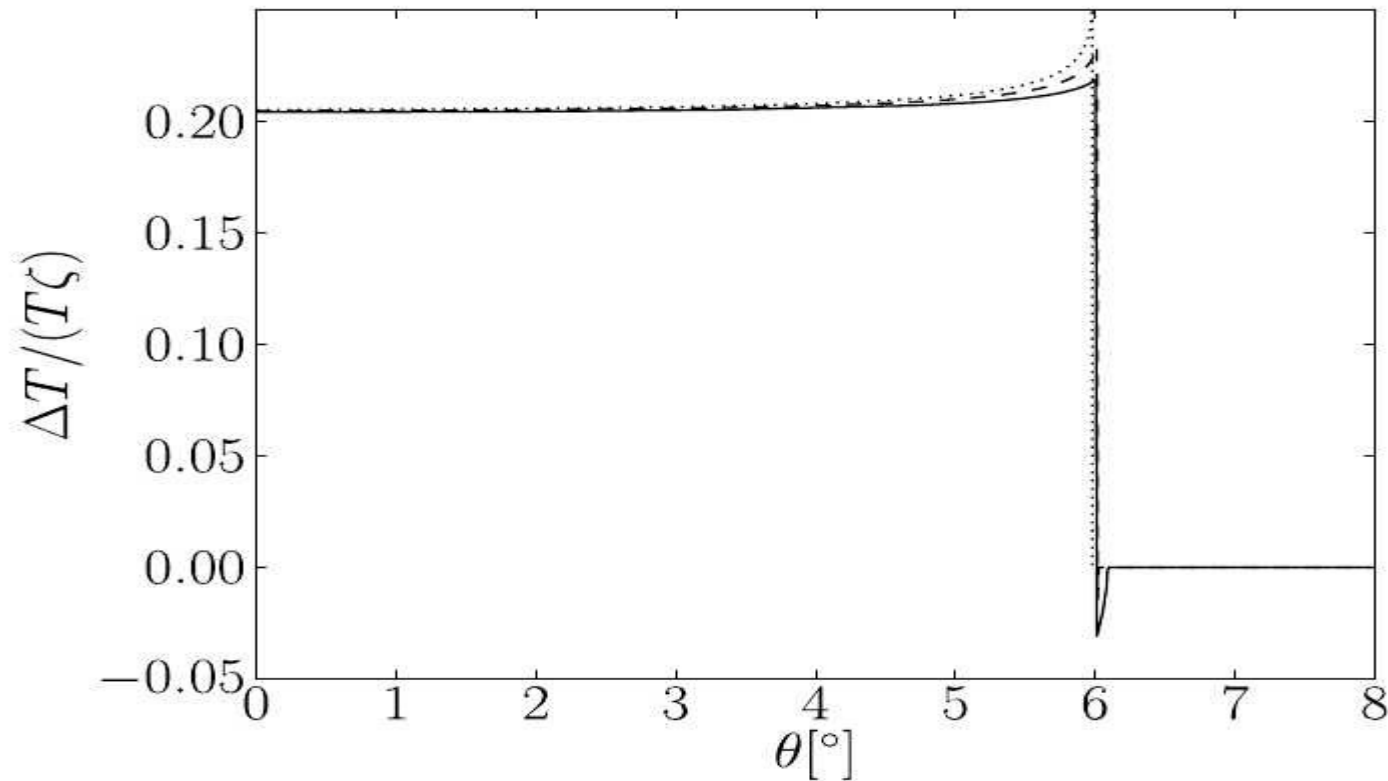




last scattering surface

# Very asymmetric bifurcations

----- SW, early ISW

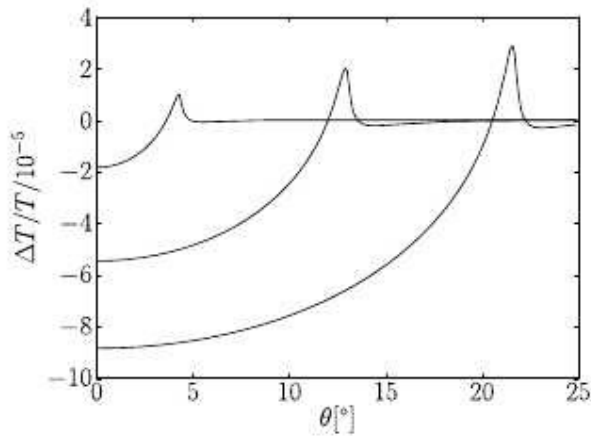


$\zeta > 0$ : under-density

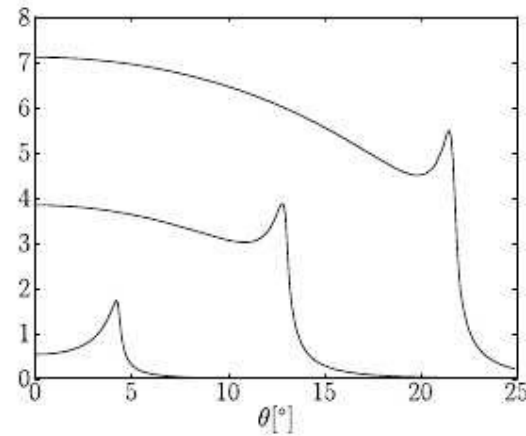
# Very asymmetric bifurcations

---- ISW,RS,OV

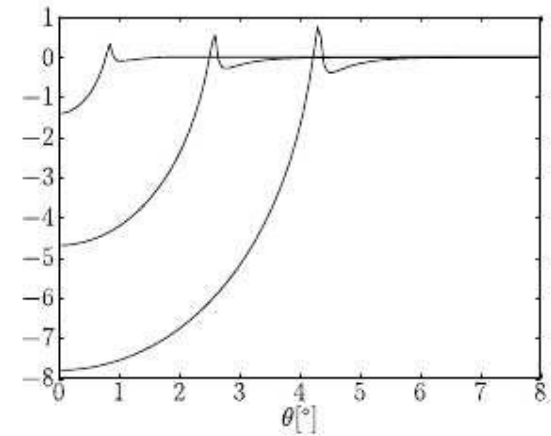
$z = 0.5$  with  $\zeta_* = 10^{-3}$



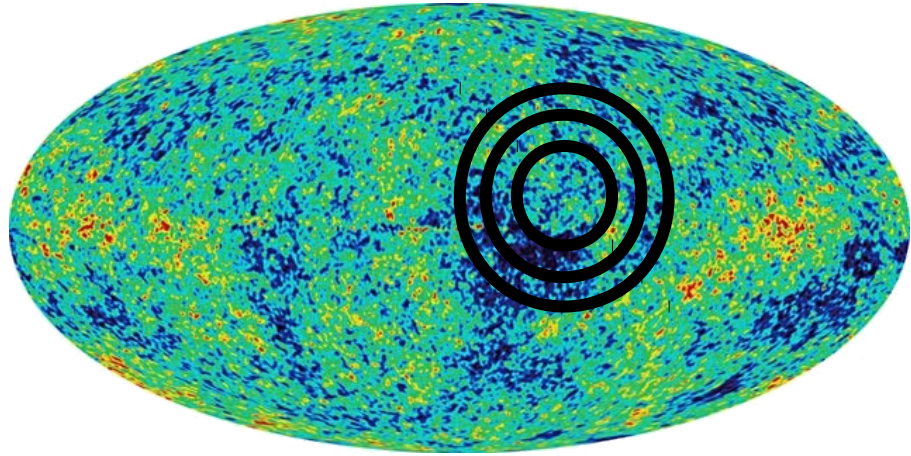
$\zeta_* = -10^{-3}$



$z = 10$  with  $\zeta_* = 10^{-3}$

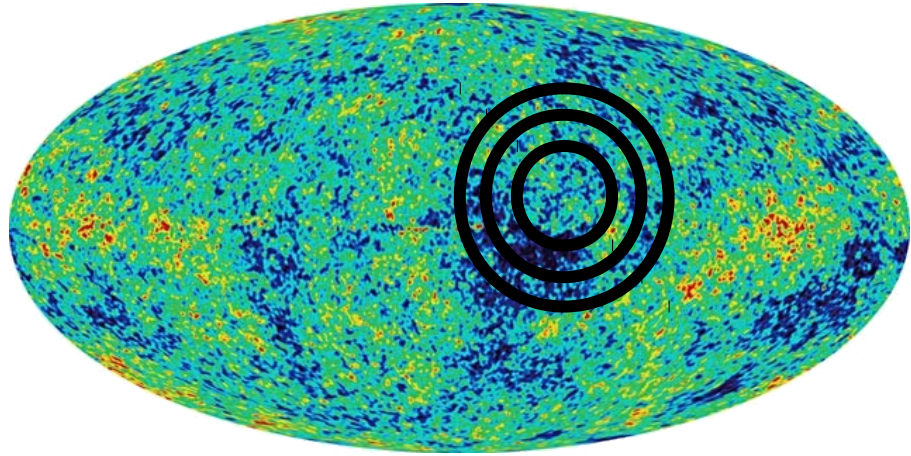


$\zeta_* > 0$ : under-density



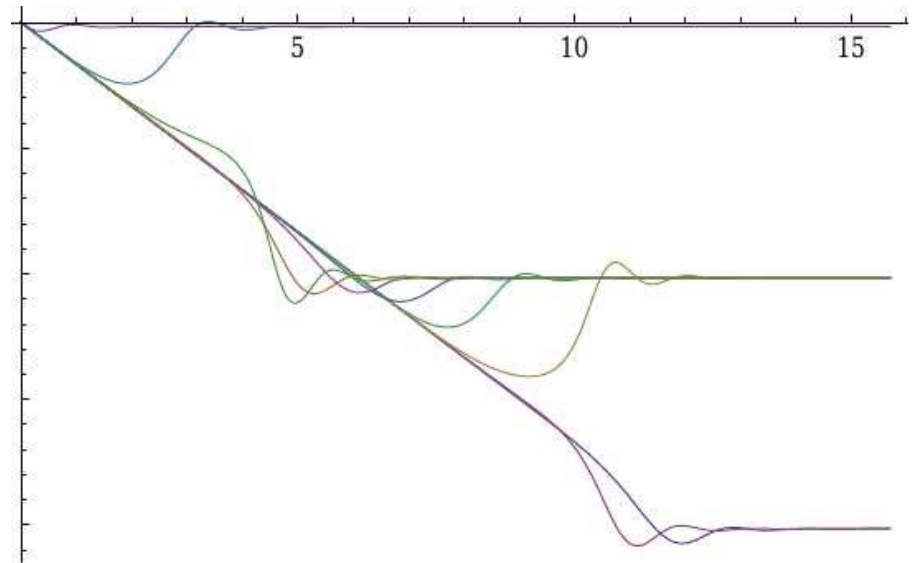
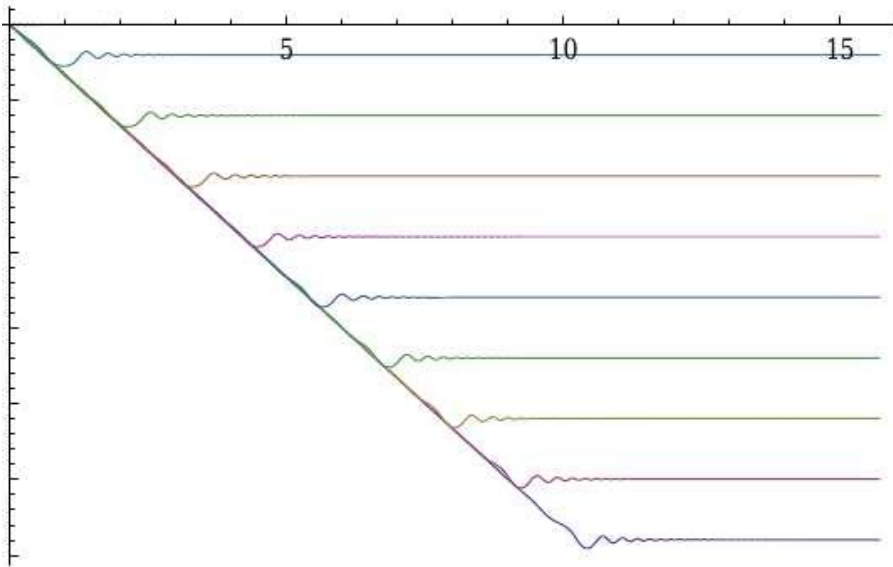
topological defect  
induced MSI

$$V = V_{sr}(\varphi) + s\chi + m^2 \Psi^a \Psi^a \cos(\chi/\chi_0)$$



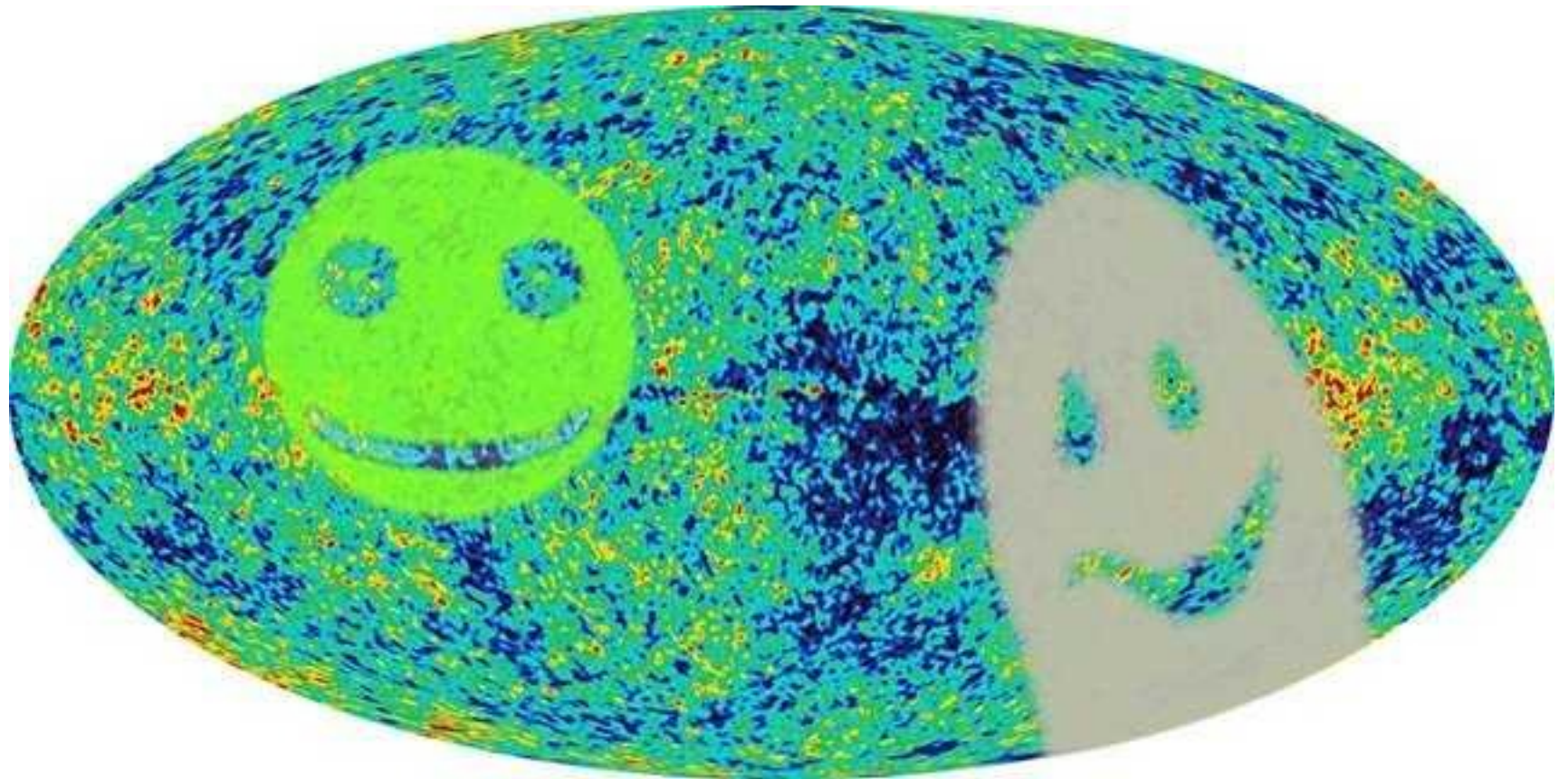
topological defect  
induced MSI

$$V = V_{sr}(\varphi) + s\chi + m^2 \Psi^a \Psi^a \cos(\chi/\chi_0)$$



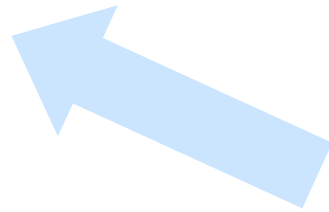
A feature in position space?

Reminder: multi-stream inflation





# Inflation in a random potential



statistics

# Inflation in a random potential

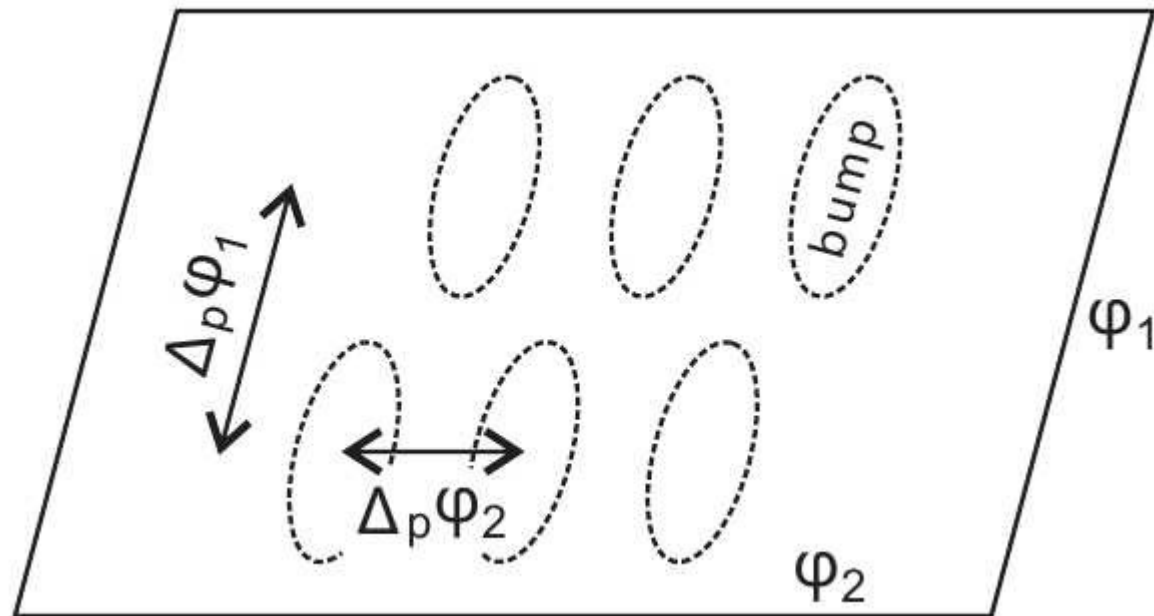
{ amplification  
random walk

isocurvature direction

# Inflation in a random potential

{ amplification  
random walk

isocurvature direction



$$\xi = \Delta_p \varphi_1 / \Delta_p \varphi_2$$

$$\ddot{\varphi}_1 + 3H\dot{\varphi}_1 + \partial_1 V(\varphi_1) + \partial_1 U(\varphi_1, \varphi_2) = 0 ,$$

$$\ddot{\varphi}_2 + 3H\dot{\varphi}_2 + \partial_2 U(\varphi_1, \varphi_2) = 0 ,$$

$$\lambda \equiv \sqrt{\langle (\partial_1 U)^2 \rangle} / |\partial_1 V| .$$

$$\delta \equiv \varphi_2^{(A)} - \varphi_2^{(B)}$$

$$\ddot{\delta} + 3H\dot{\delta} + (\partial_2^2 U)\delta = 0$$

oscillate?

or grow?

$$\partial_2^2 U \simeq \frac{\lambda \xi \partial_1 V}{\Delta_p \varphi_2} \sin \left( \frac{2\pi \dot{\varphi}_1 t}{\Delta_p \varphi_1} \right)$$

Mathieu equation,

Analog: broad resonance preheating

$$\delta \equiv \varphi_2^{(A)} - \varphi_2^{(B)}$$

$$\ddot{\delta} + 3H\dot{\delta} + (\partial_2^2 U)\delta = 0$$

$$\partial_2^2 U \simeq \frac{\lambda \xi \partial_1 V}{\Delta_p \varphi_2} \sin\left(\frac{2\pi \dot{\varphi}_1 t}{\Delta_p \varphi_1}\right)$$



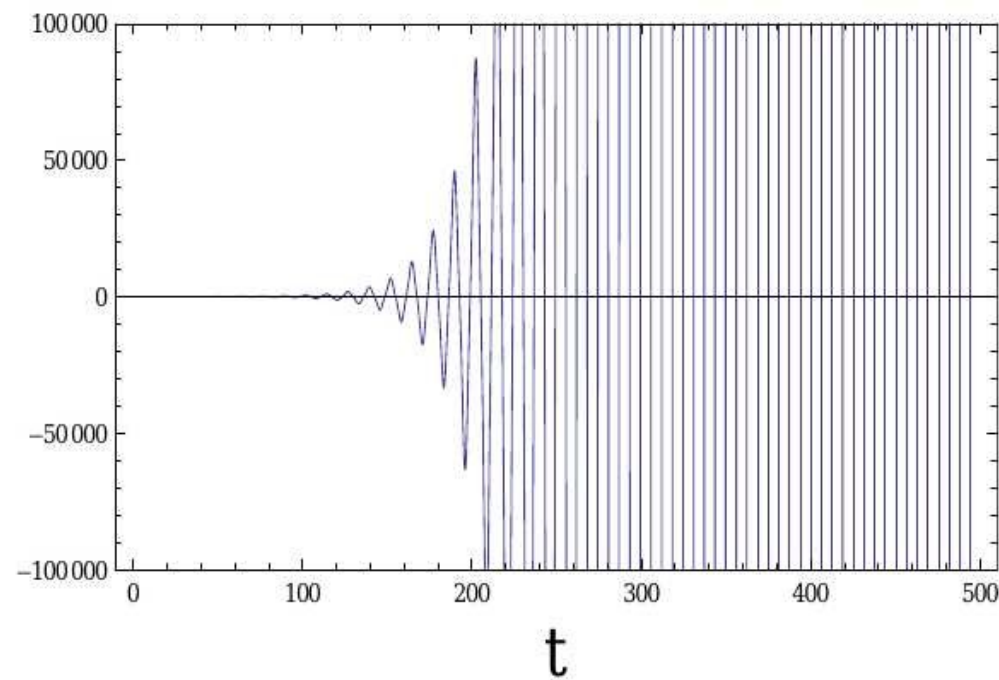
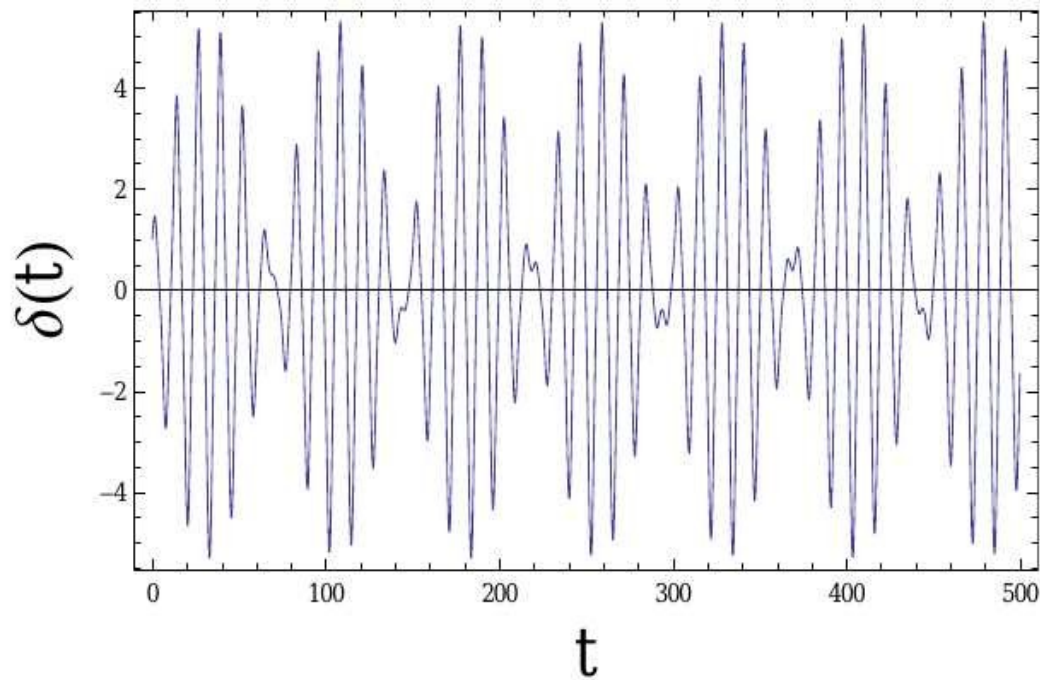
# Rough estimate: neglect friction

$$\ddot{\delta} + 3H\dot{\delta} + (\partial_2^2 U)\delta = 0$$

$$\ddot{\delta} + \alpha \sin(\beta t)\delta = 0$$

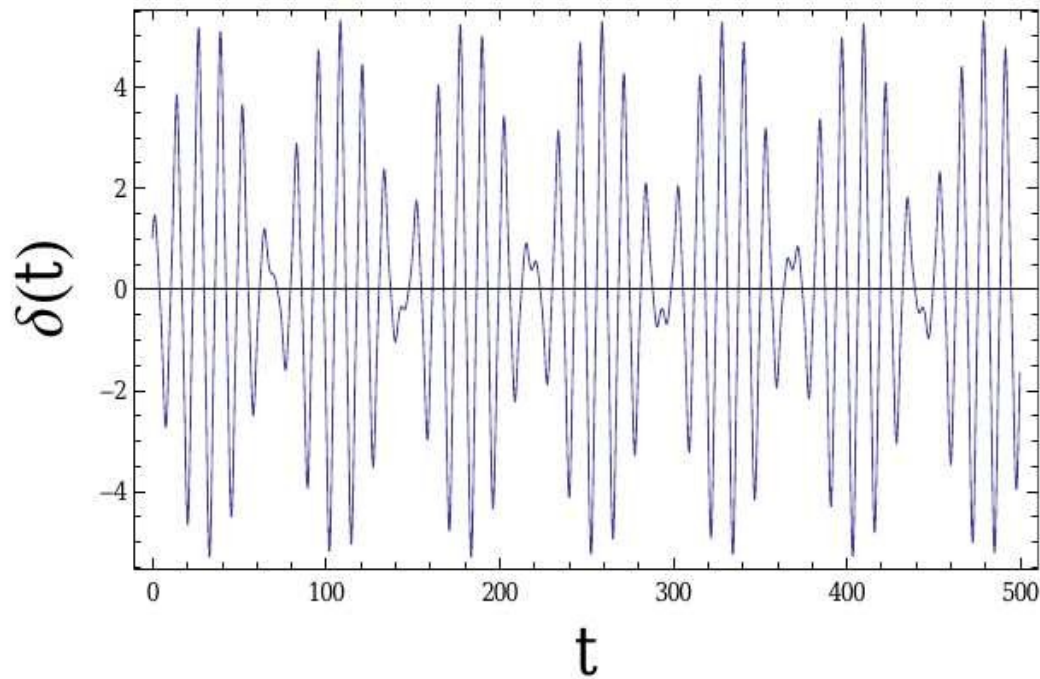
$$\alpha/\beta^2 = 0.45$$

$$\alpha/\beta^2 = 0.46$$

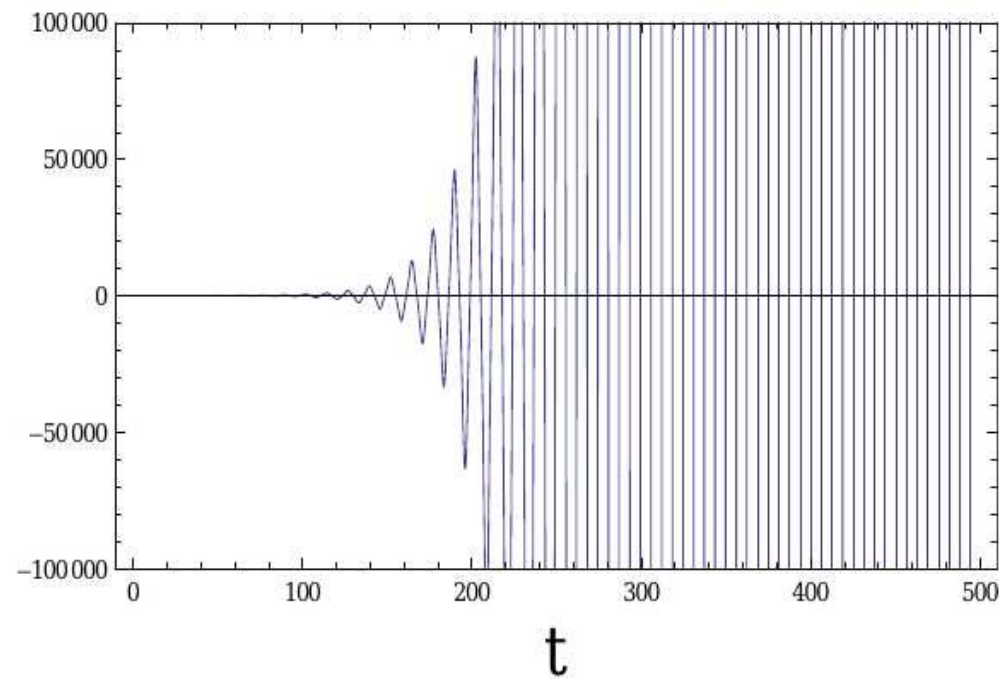


# Transition between oscillation and growing

$$\alpha/\beta^2 = 0.45$$



$$\alpha/\beta^2 = 0.46$$

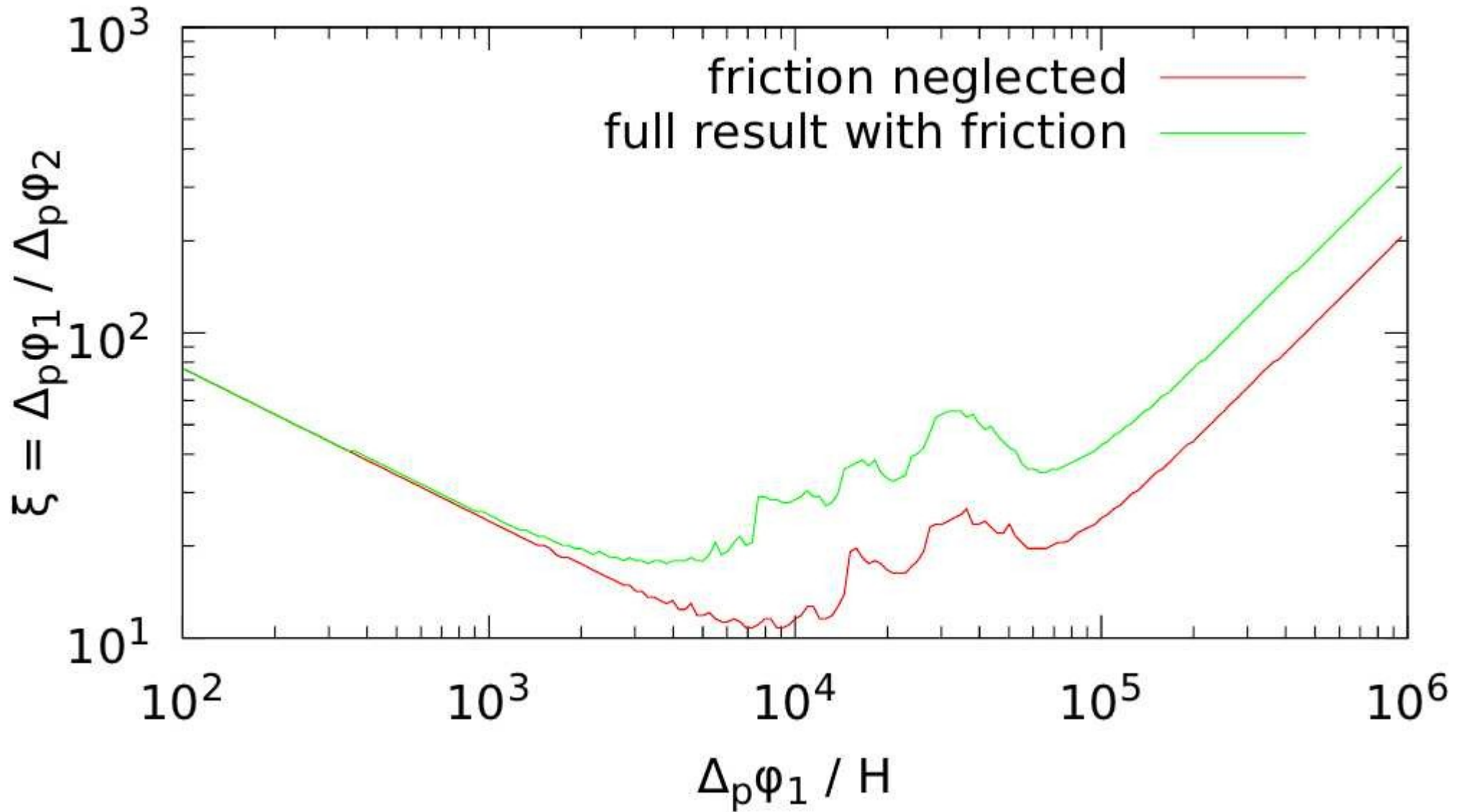


growing solution

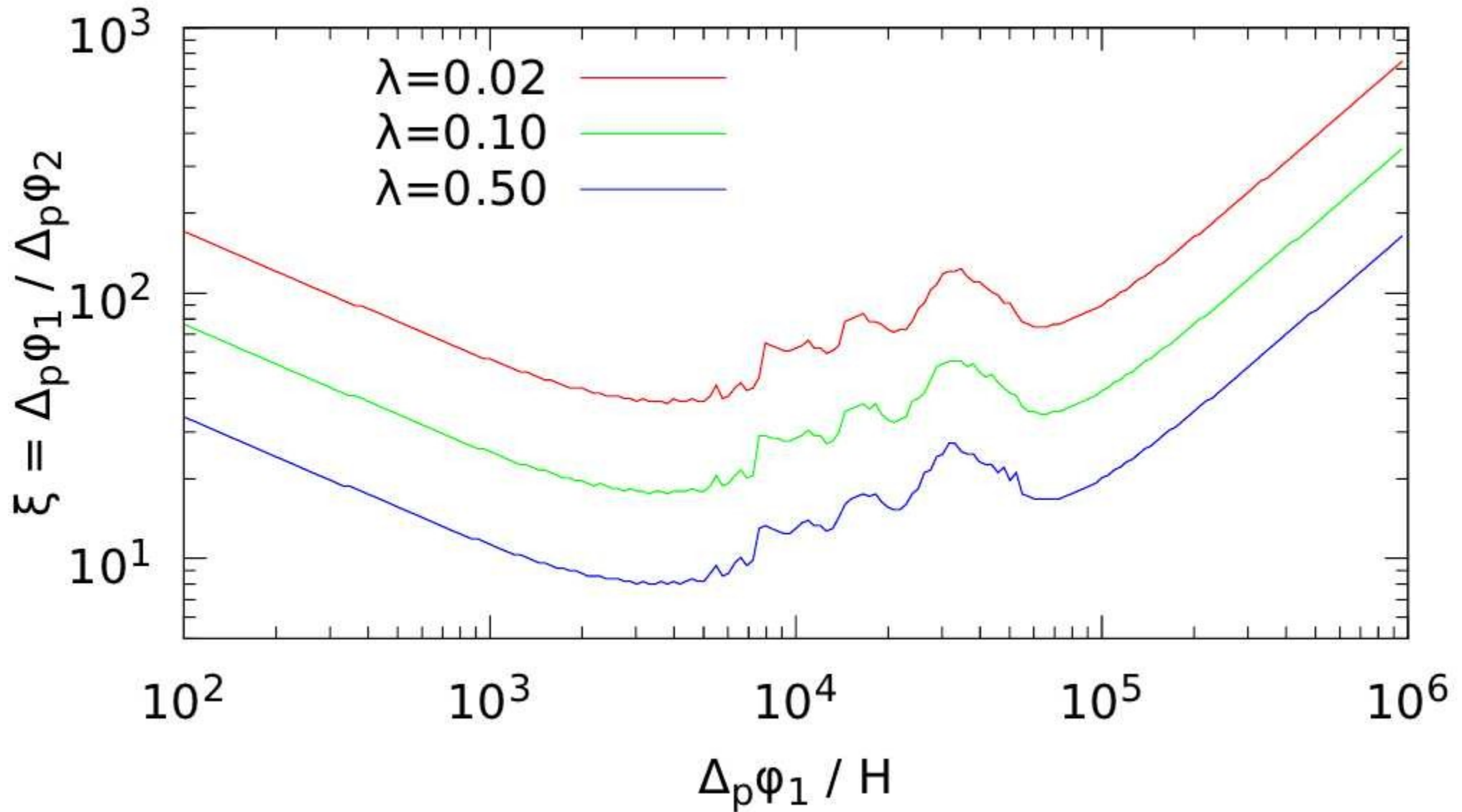
$$\frac{1}{\lambda \xi^2 P_\zeta^{1/2}} < \frac{\Delta_p \varphi_1}{H} < \frac{3 \lambda \xi^2}{2 \pi P_\zeta^{1/2}}$$

large exponent

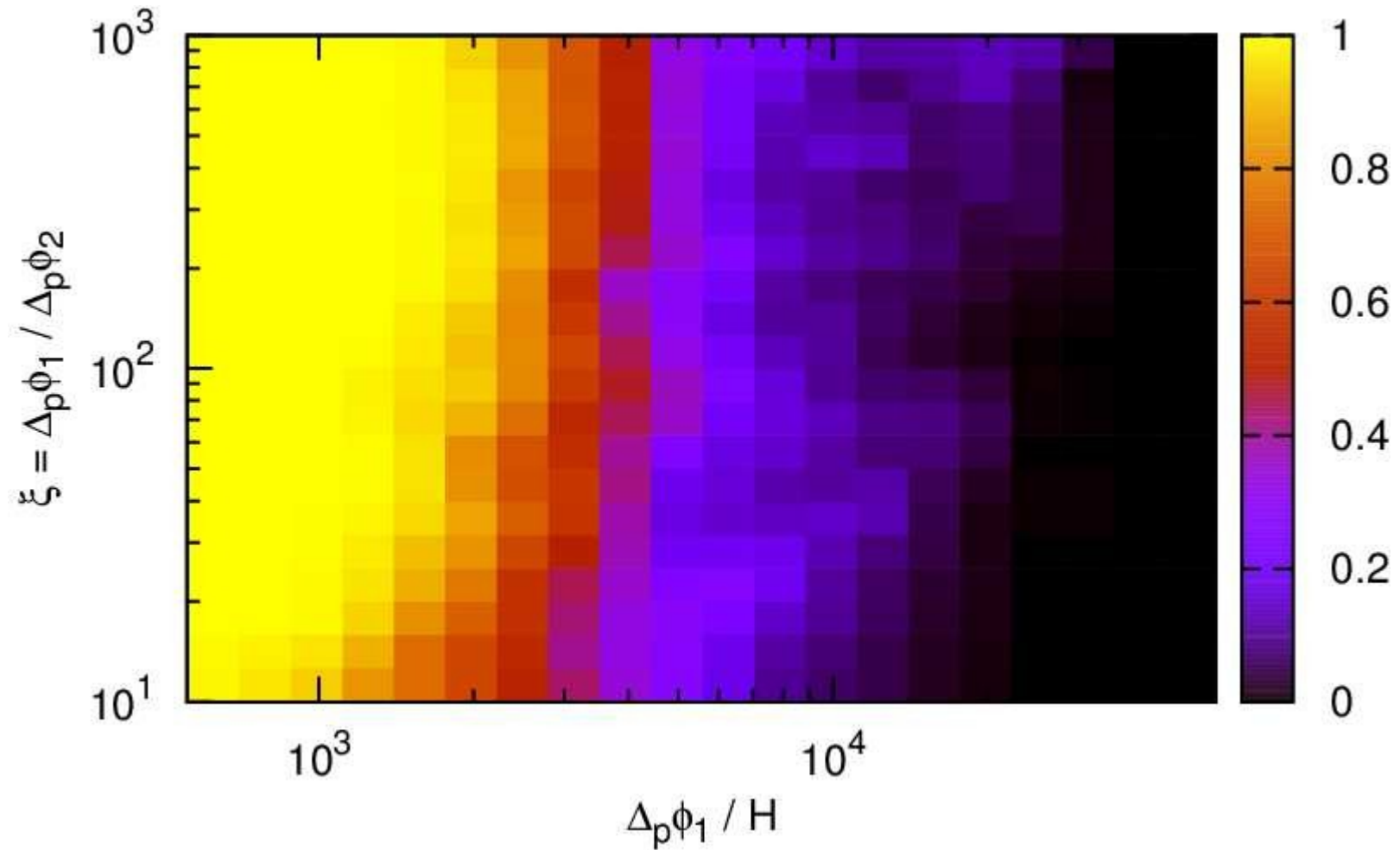
with friction



with friction



# Completely random potential

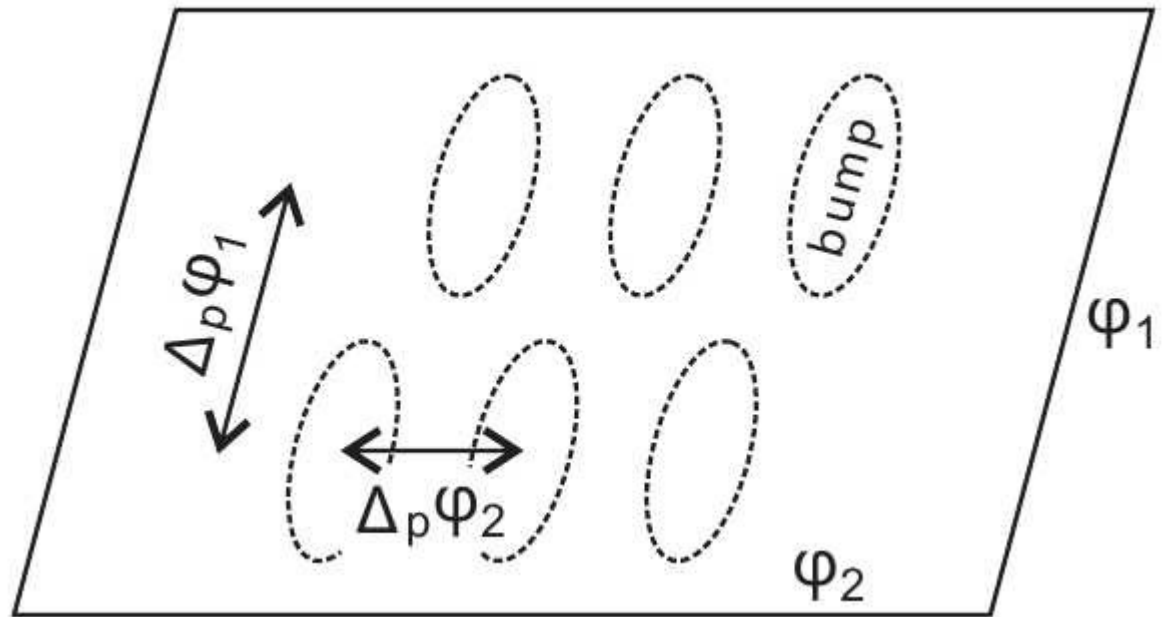




# Inflation in a random potential

{ amplification  
random walk

isocurvature direction



$$\Delta t = \Delta_p \varphi_1 / \dot{\varphi}_1$$

$$\Delta_q \varphi_2 = P_\zeta^{1/4} \sqrt{\frac{H \Delta_p \varphi_1}{2\pi}} = \Delta_q^H \varphi \sqrt{\frac{\Delta_p \varphi_1}{\Delta_c^H \varphi_1}}$$

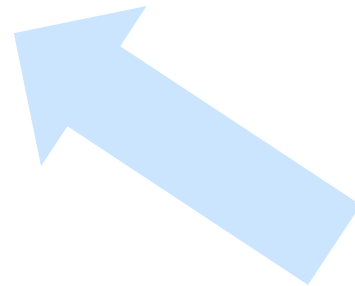
$$P_{\text{bifur}}(\Delta t) = \Delta_q \varphi_2 / \Delta_p \varphi_2$$

$$\Delta_p \varphi_1 = \Delta_p \varphi_2$$

$$P_{\text{bifur}}(N) \begin{cases} \simeq \tilde{N} P_{\text{bifur}} \\ \simeq \frac{2}{3} \tilde{N}^{3/2} P_{\text{bifur}} \end{cases}$$

$$\begin{cases} \leq 20H \\ \leq 100H \end{cases}$$

$\times (6N)^{1/2}$  for the whole  
observable universe:  
about  $160H$  or  $800H$



bifurcation along  
a world line

$$\Delta_p \varphi_1 = \Delta_p \varphi_2$$

$$\left\{ \begin{array}{l} \leq 20H \\ \leq 100H \end{array} \right.$$



# Chain inflation:

Freese, Spolyar (2004)

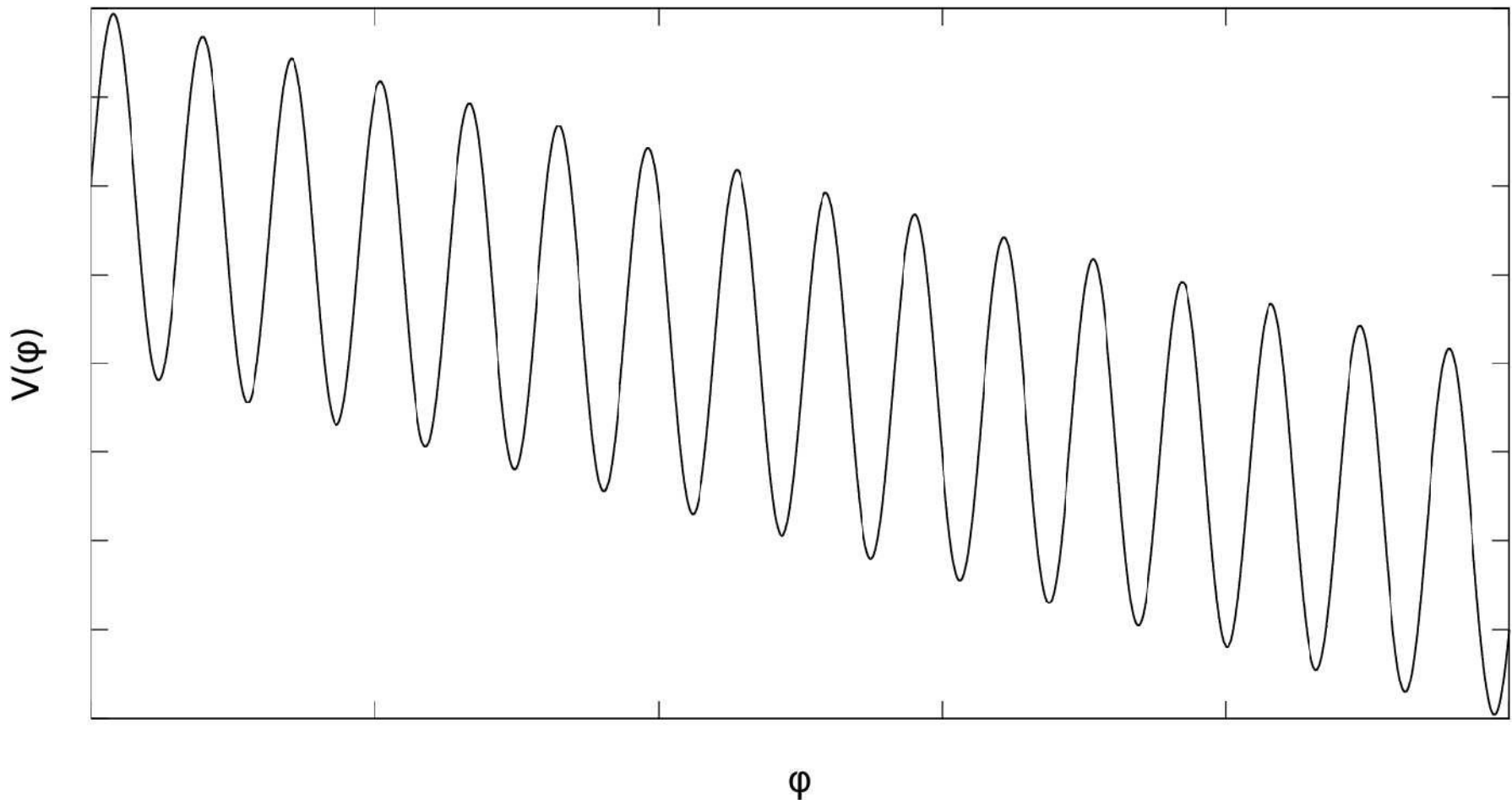
See also: Liu, Feldstein, Tweedie, Chialva,  
Danielsson, Huang, ...

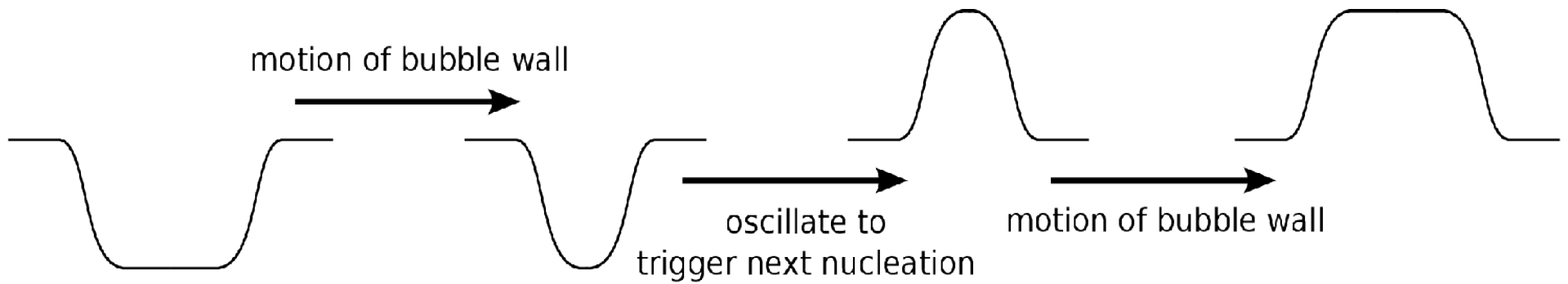
Perturbations: first full, analytical result





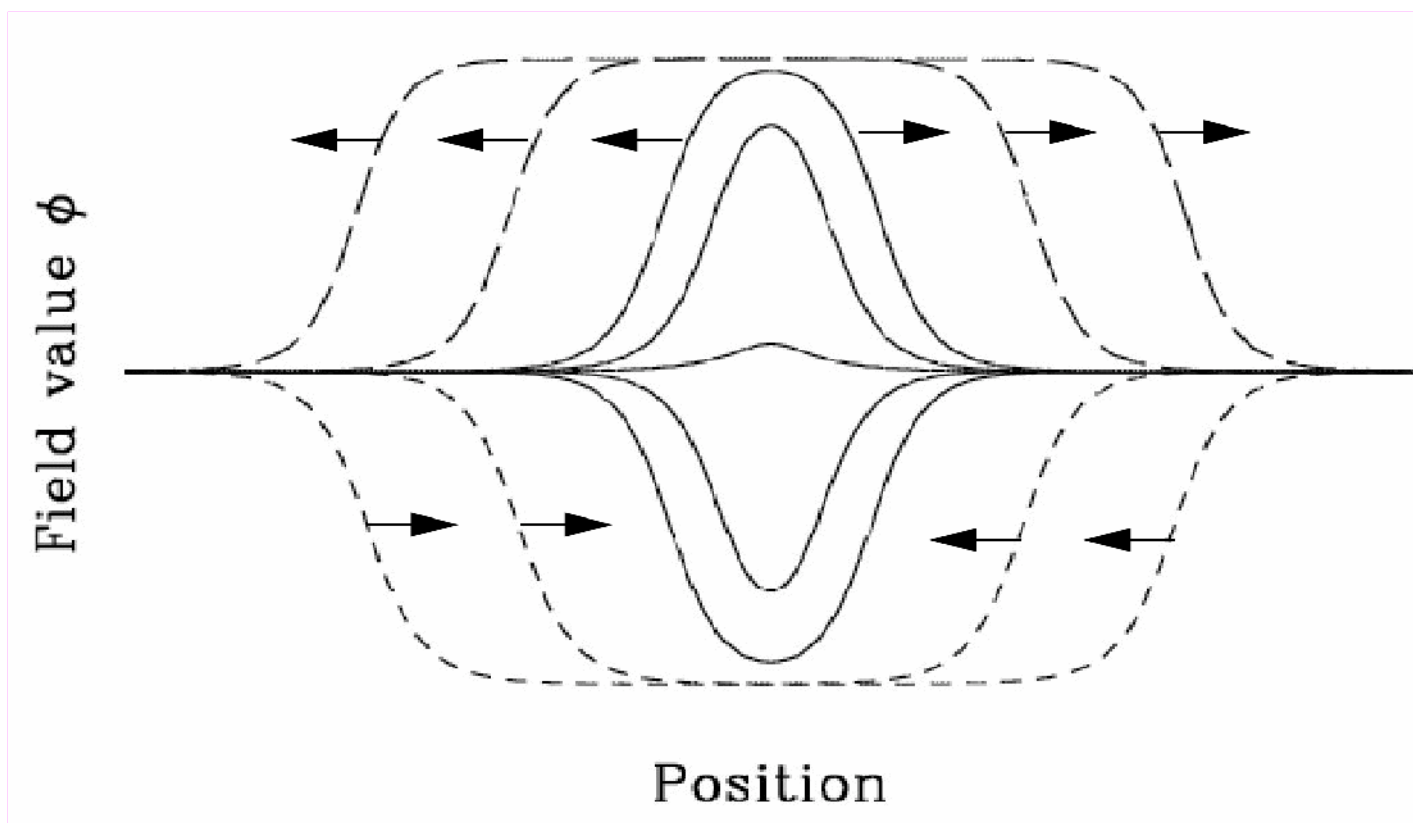
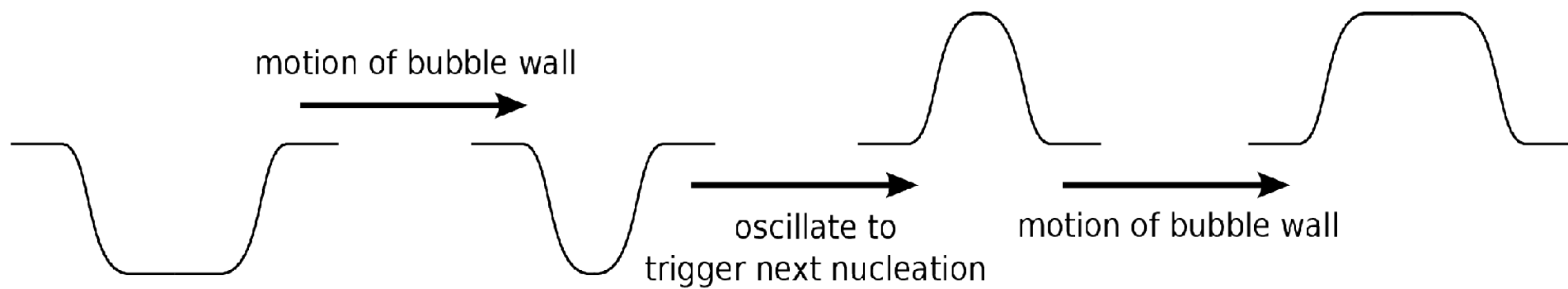
# Simplest first single field chain inflation





New effect:

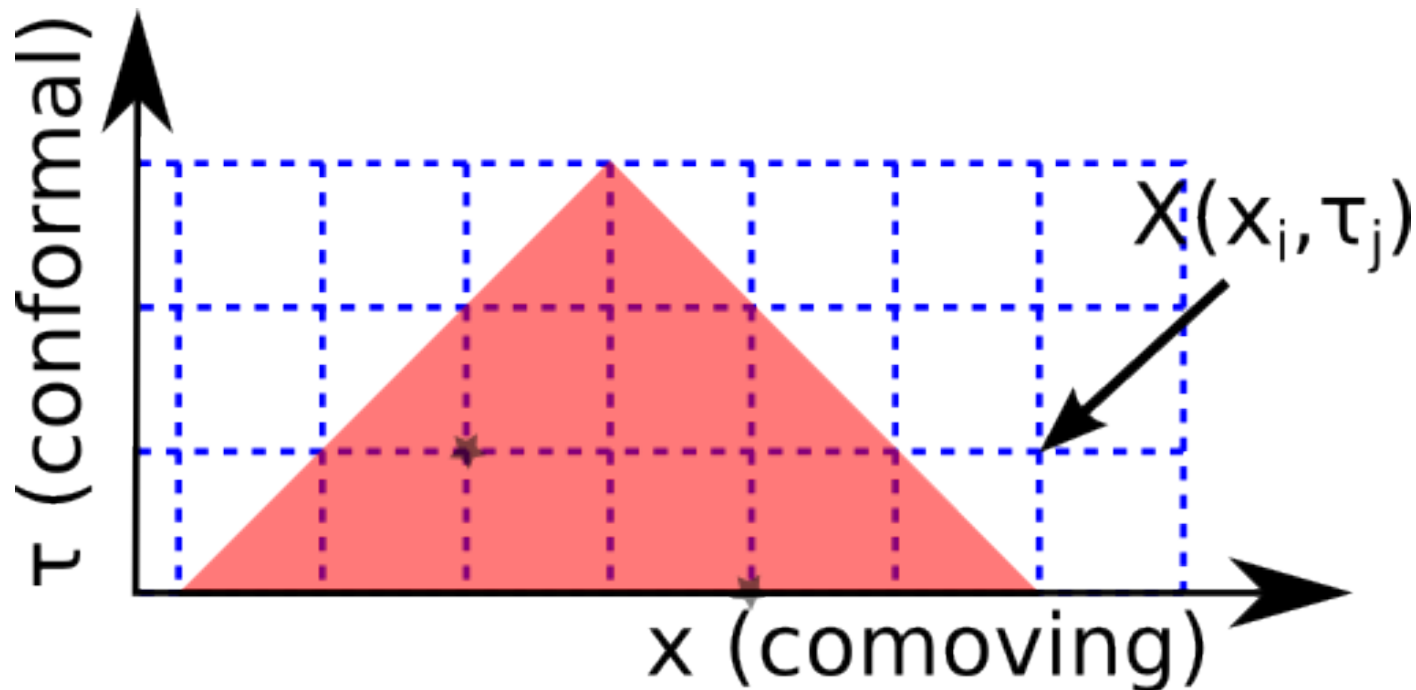
walls collide into new bubble

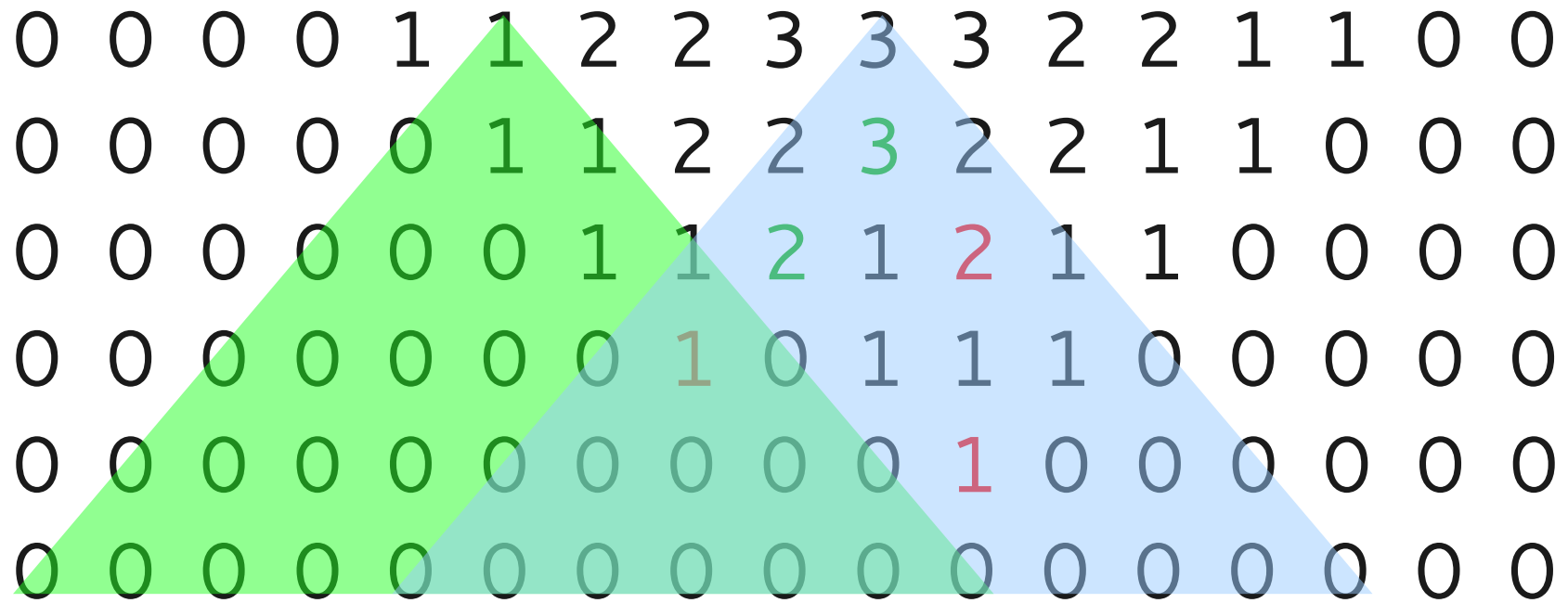


# Background dynamics

$$\varphi(x, \tau) = \sum X(x_i, \tau_j)$$

$$X \sim \text{Pois}(\Gamma H^{-4} \tau^{-4} d^3x d\tau)$$



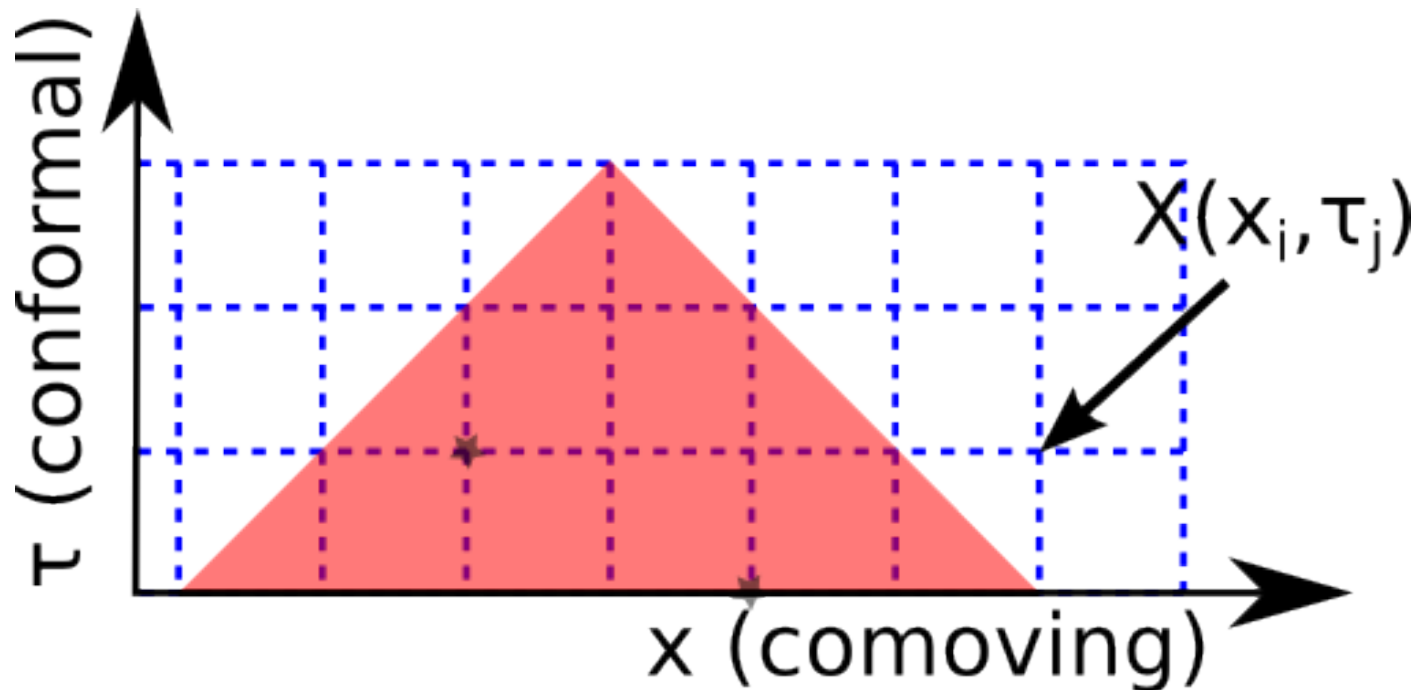


spontaneous  
from bubble wall

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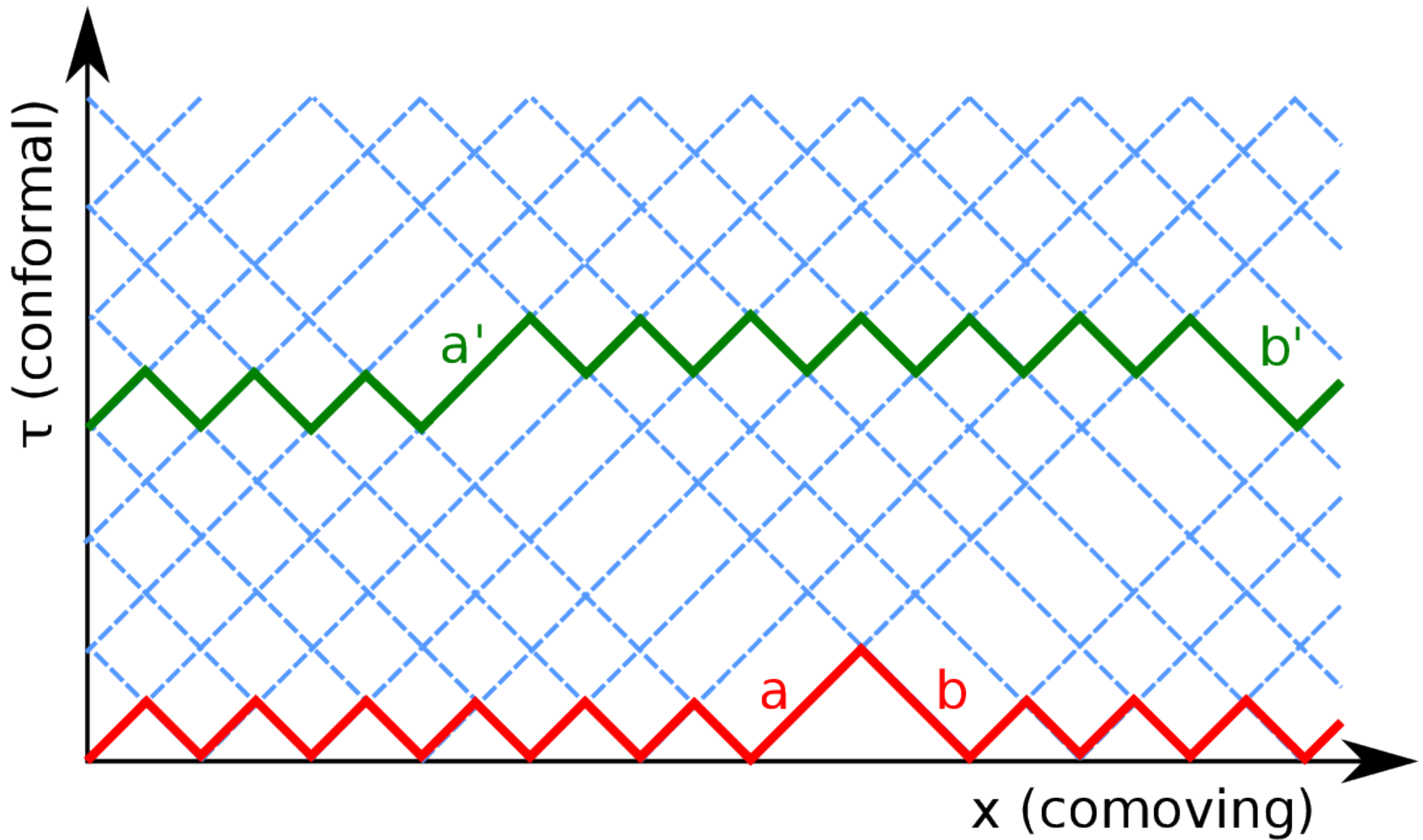
$$X \sim \text{Pois}(\Gamma H^{-4} \tau^{-4} d^3x d\tau)$$

$$X+Y \sim \text{Pois}(\lambda_X + \lambda_Y)$$

$$\varphi(x, \tau) \sim \text{Pois}(\varphi_i + 4\pi\Gamma H^{-3}(t-t_i)/3)$$

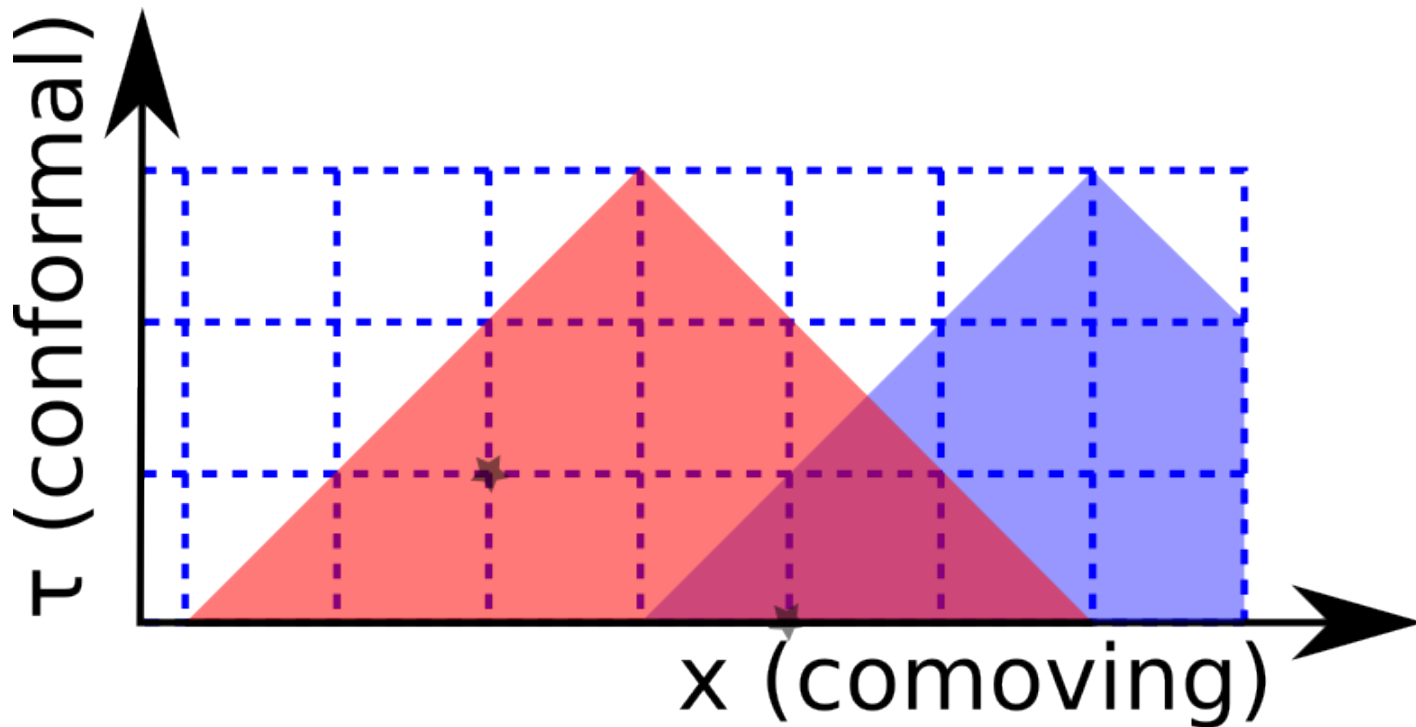
$$\langle \varphi(x, \tau) \rangle = \varphi_i + 4\pi\Gamma H^{-3}(t-t_i)/3$$

# Perturbation theory



$$\langle \delta\varphi(x, \tau) \delta\varphi(x + r, \tau) \rangle_c = \int_{-\infty}^{\tau} d\eta V_2(\eta) \frac{\Gamma}{H^4 \eta^4}$$

$$V_2(\eta) \equiv \begin{cases} \frac{4}{3}\pi(\tau - \eta - r/2)^2(\tau - \eta + r/4) & \text{when } r < 2(\tau - \eta) \\ 0 & \text{when } r \geq 2(\tau - \eta) \end{cases}$$



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$$\langle \delta\varphi_{\mathbf{k}_1}(\tau) \delta\varphi_{\mathbf{k}_1}(\tau) \rangle = (2\pi)^3 \delta^3(\mathbf{k}_1 + \mathbf{k}_2) \int_0^{\infty} dr \frac{4\pi r \sin(kr)}{k} \langle \delta\varphi(0, \tau) \delta\varphi(r, \tau) \rangle$$

$$\langle \delta\varphi_{\mathbf{k}_1}(\tau) \delta\varphi_{\mathbf{k}_1}(\tau) \rangle = (2\pi)^3 \delta^3(\mathbf{k}_1 + \mathbf{k}_2) \frac{8\pi^3 \Gamma}{3H^4 k^3}$$

$$\langle \zeta_{\mathbf{k}_1} \zeta_{\mathbf{k}_2} \rangle = (2\pi)^3 \delta^3(\mathbf{k}_1 + \mathbf{k}_2) \frac{3\pi H^4}{2\Gamma k^3} \quad P_\zeta = \frac{3H^4}{4\pi\Gamma}$$

consistency checks

$$V = V_0 - A\phi + V_1 \sin(2\pi\phi/\Delta\phi)$$

①  $\lambda = V'''' < 2\pi$  everywhere

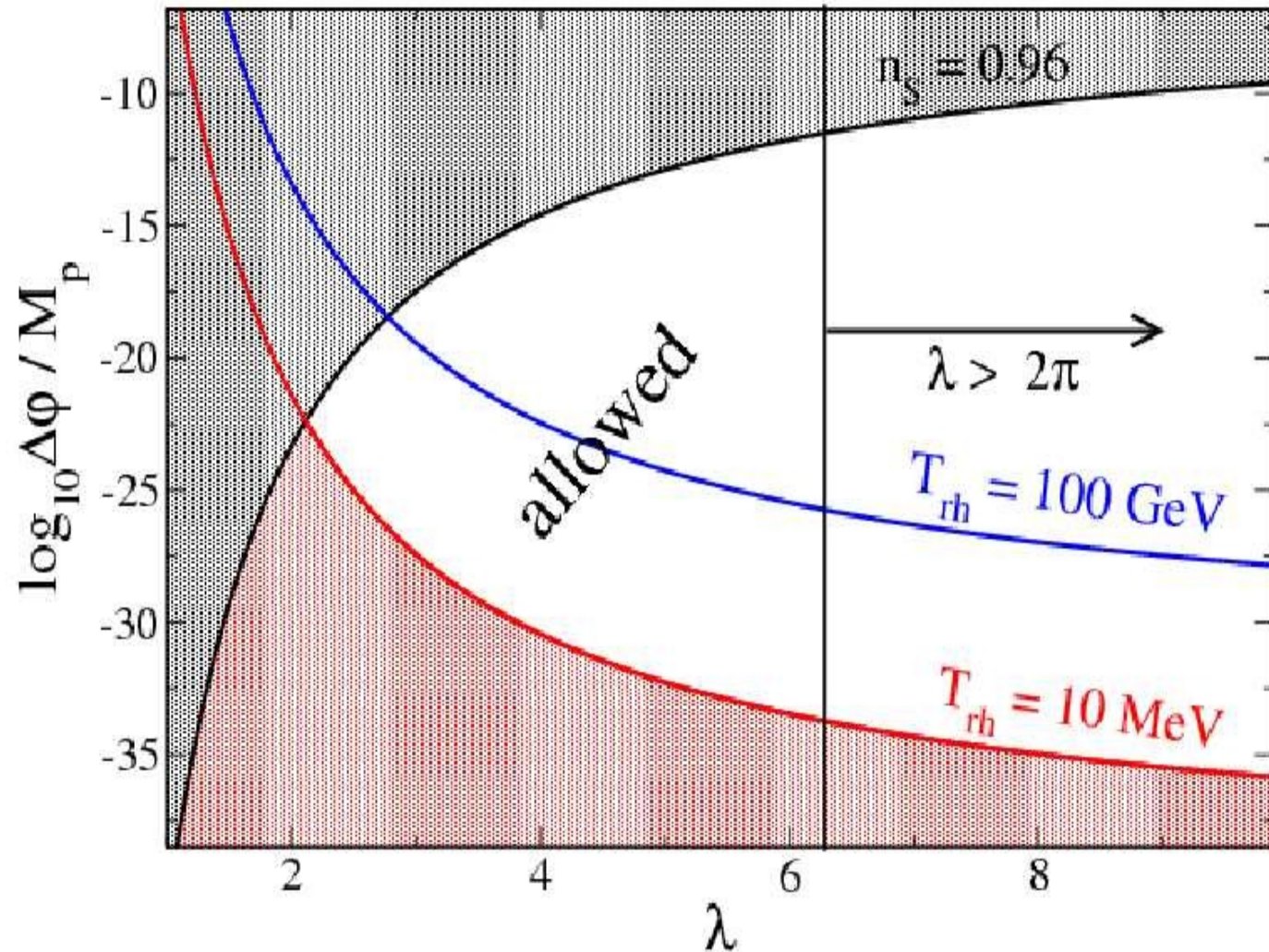
=====

②  $\phi$  stuck

③ small bubble

④  $P_\zeta$  and  $n_s$

⑤  $T_{\text{reh}}$







**The potential may get burnt !**

# (non-) Gaussianity

$$X \sim \text{Pois}(\lambda), \langle (X - \langle X \rangle)^3 \rangle = \lambda$$

$$\langle \delta\varphi^3 \rangle = \int_{-\infty}^{\tau} d\eta V_3(\eta) \frac{\Gamma}{H^4 \eta^4}$$

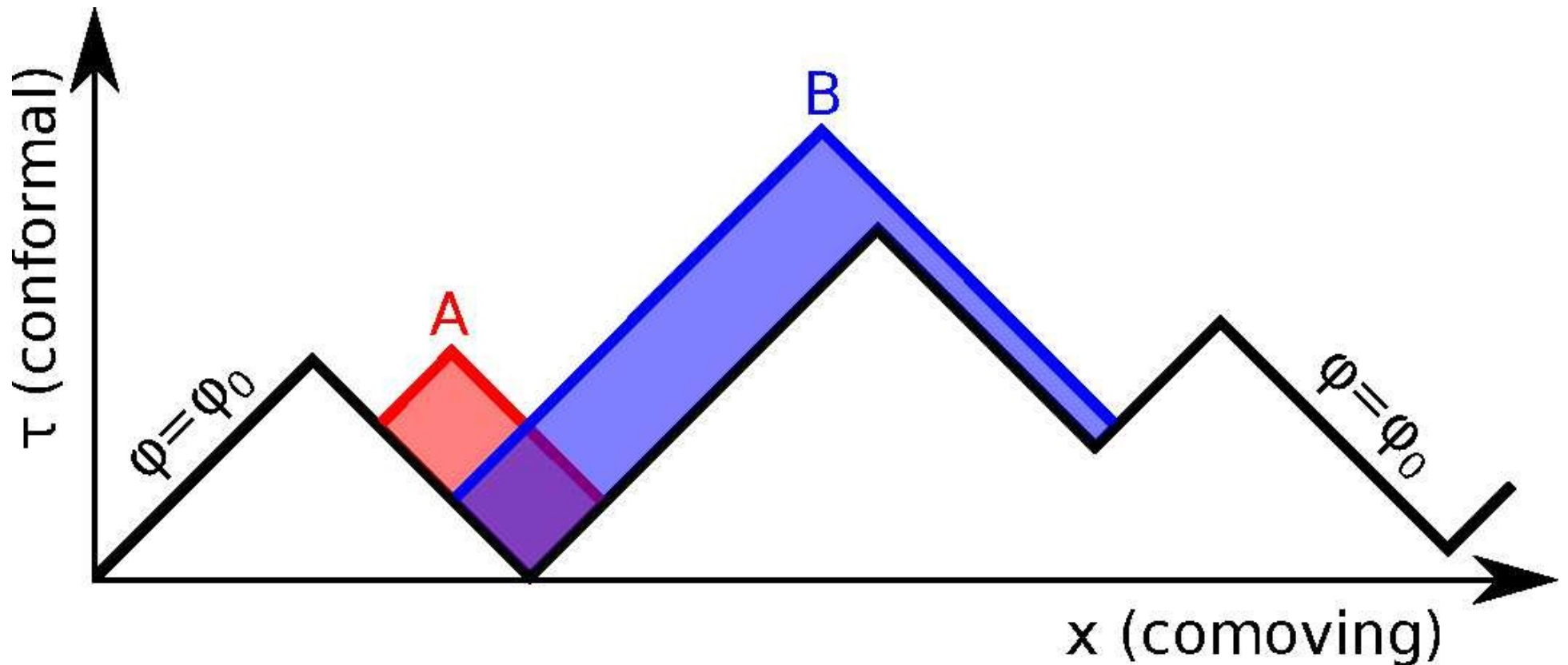
$$\langle \zeta^3 \rangle = (2\pi)^3 \delta^3(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) \frac{3}{5} f_{NL} P_{\zeta}^2$$

$$f_{NL} \sim g_{NL} \sim h_{NL} \sim i_{NL} \sim \dots \sim \mathcal{O}(1)$$

Multiple fields:

same power spectrum

fewer minimas needed







**Merci!**

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