

Large scale structure formation with the Schrödinger method

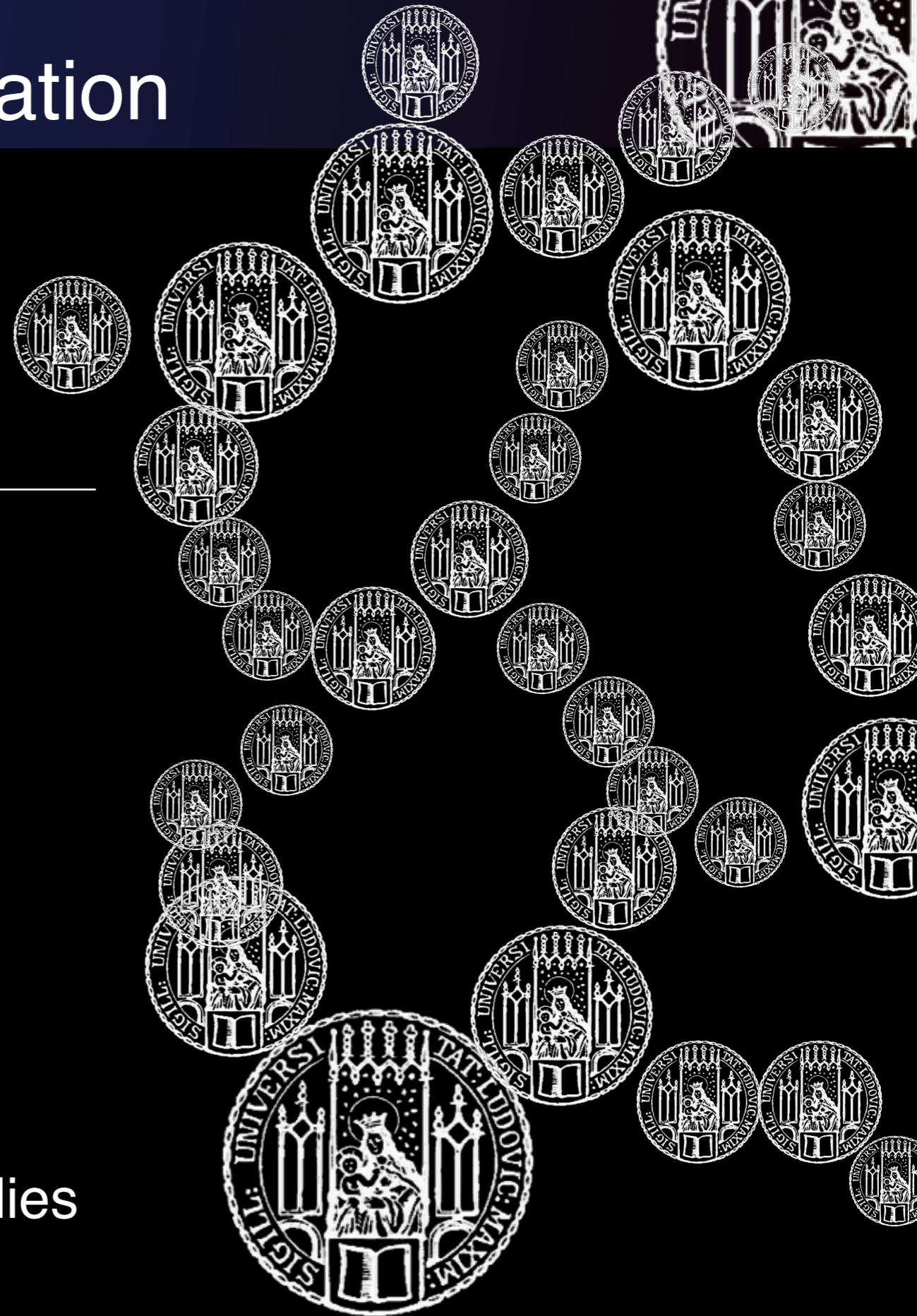
Cora Uhlemann

Arnold Sommerfeld Center, LMU
& Excellence Cluster Universe

to be continued in

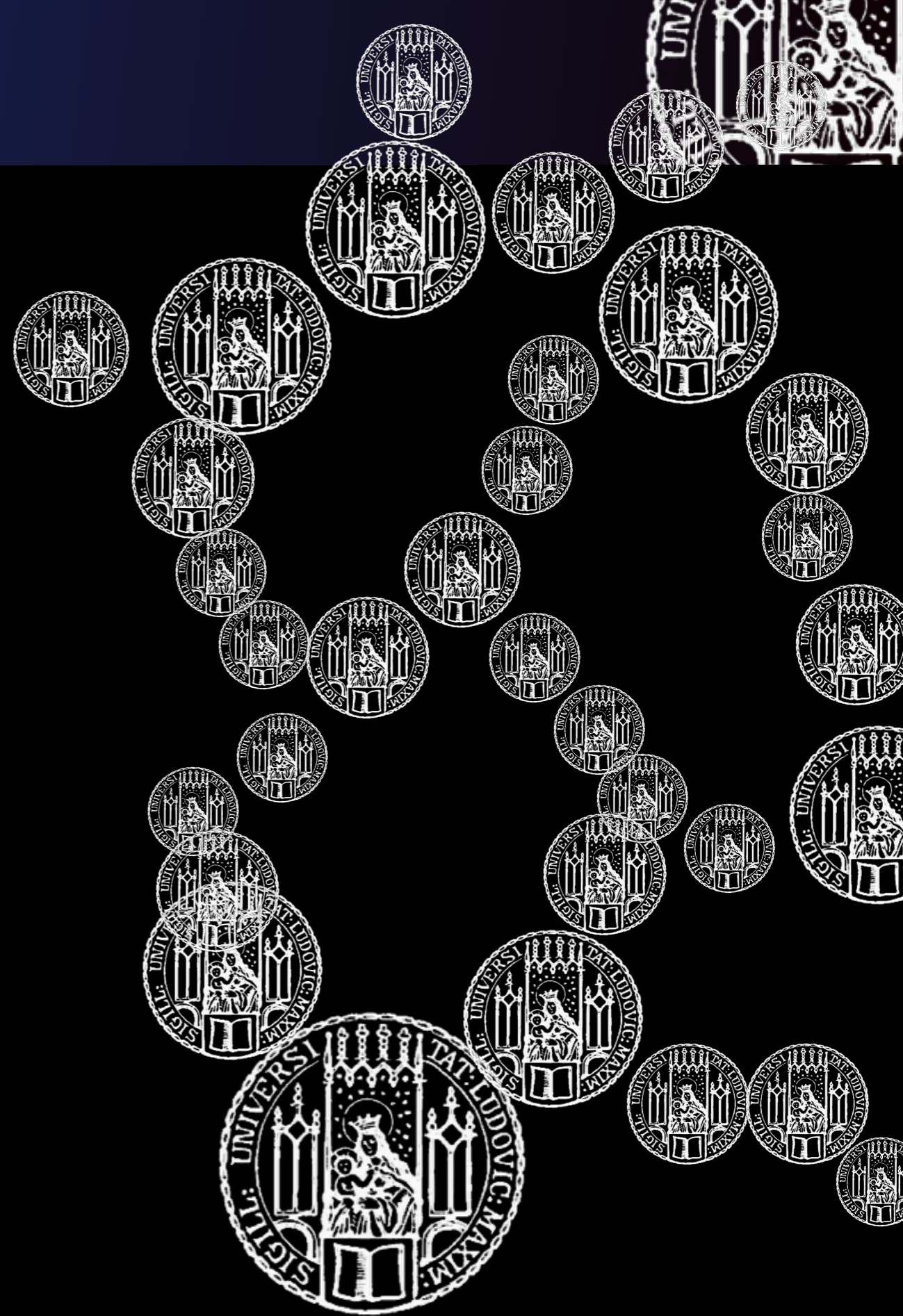
Utrecht University, Netherlands

visitor grant sponsored by
Balzan Centre for Cosmological Studies



Outline

1. Structure formation
2. Analytical description of cold dark matter
 - a. dust model
 - b. Schrödinger method
 - c. coarse-grained dust model
3. Summary



Cosmological Structure Formation



-13.8 billion years: nearly uniform initial state

Inflation

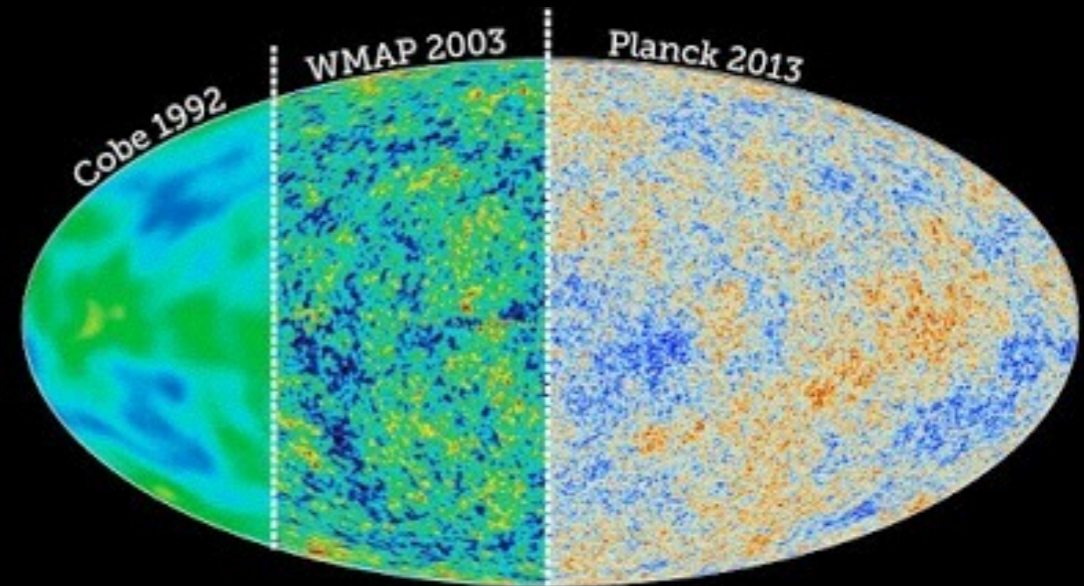
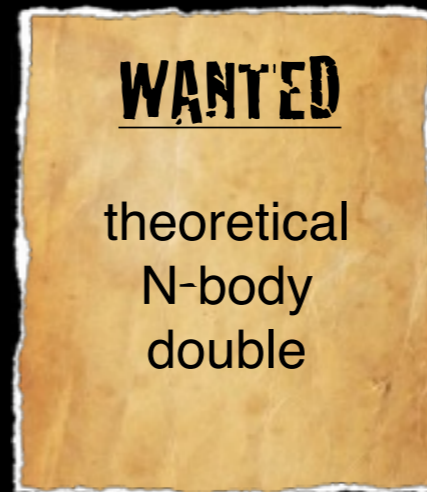
- established `boring` initial conditions
 - quantum fluctuations get amplified
 - primordial plasma cools → recombination → CMB

Structure formation

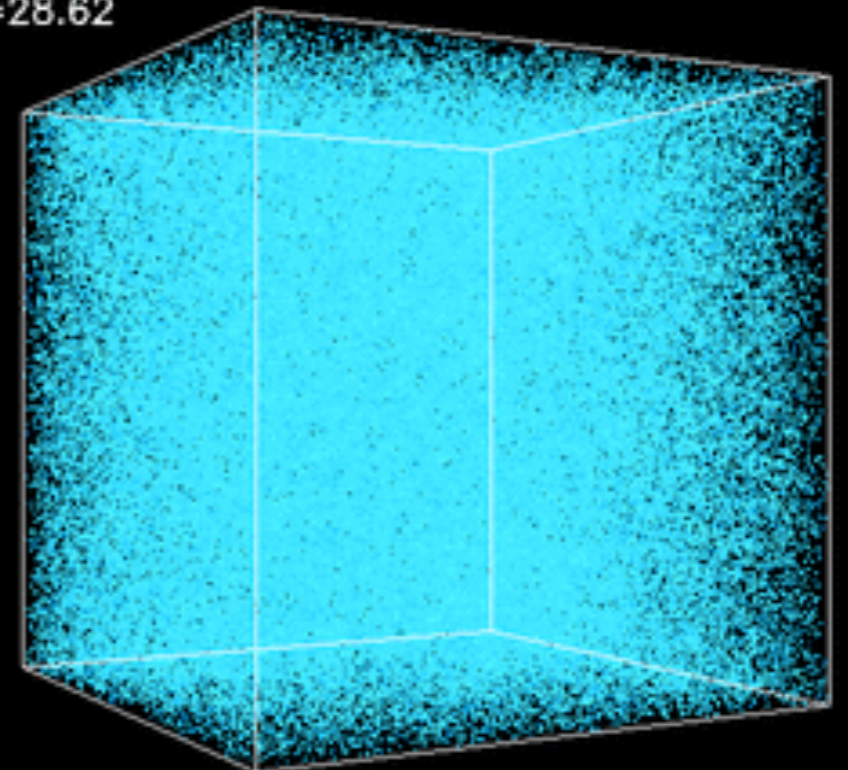
- hierarchical
- tiny over-densities act as seeds
 - congregation via gravitational instability
 - collapse into bound structures

Large scale structure: Dark Matter

- linear regime
 - ✓ analytically understood
- nonlinear stage
 - ?! N-body simulations inevitable



Z=28.62



today: rich structures in cosmic web

Kravtsov & Klypin (simulations @NCSA)



Describing Cold Dark Matter with the Schrödinger method

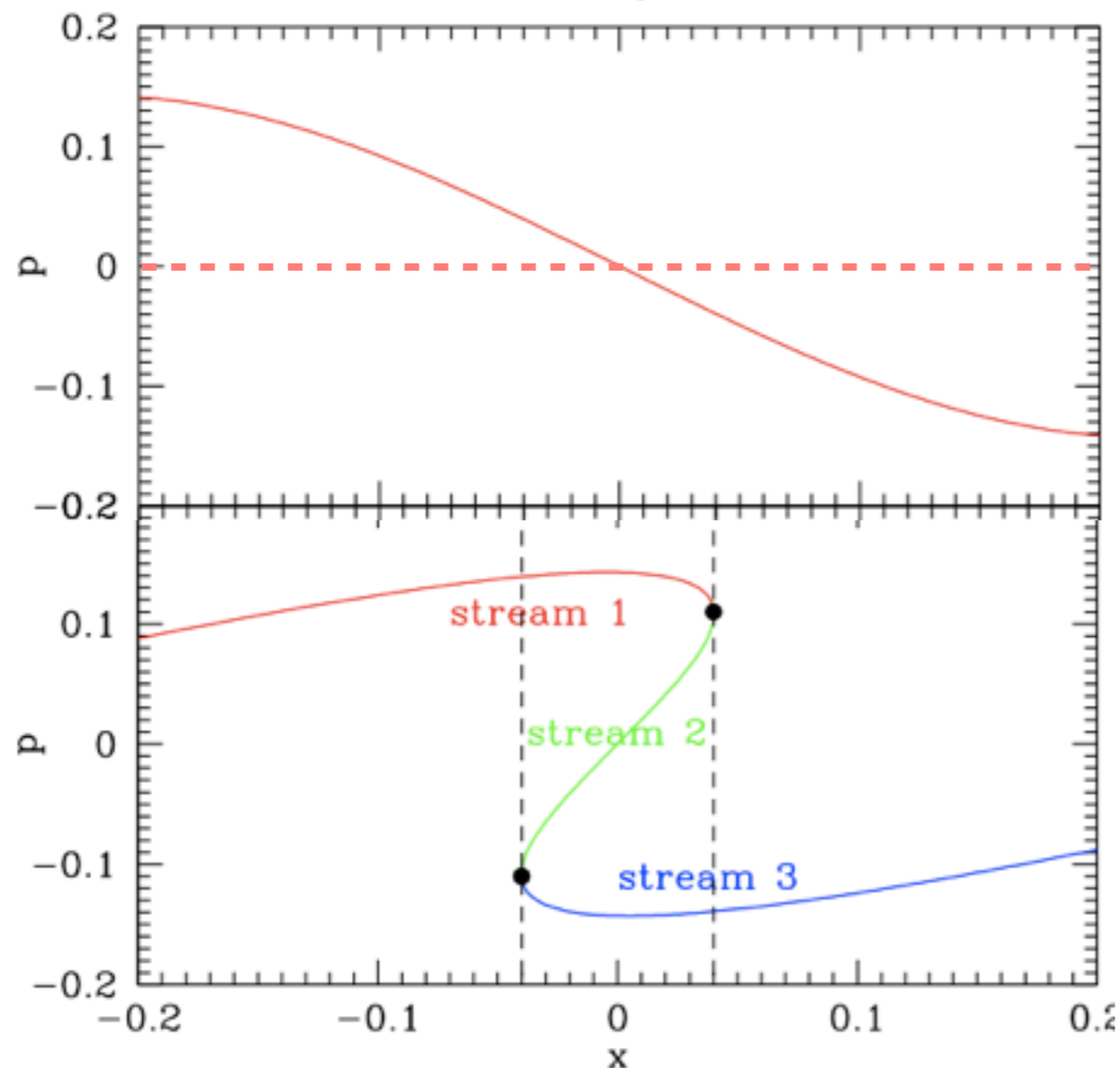
Describing Cold Dark Matter



phase space distribution function $f(t,x,p)$

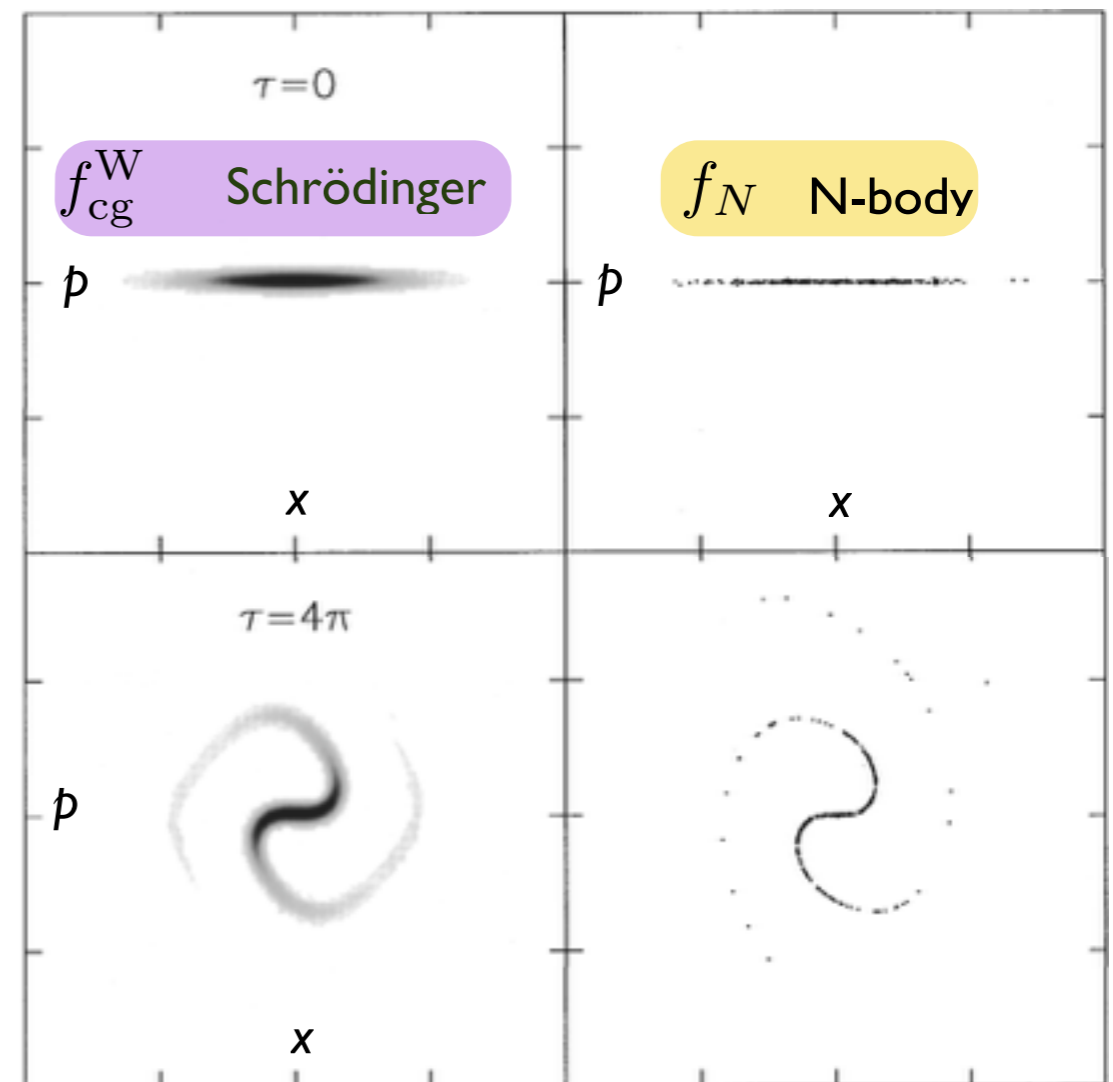
- describes number density & distribution of momenta p

theoretical expectation



Pueblas & Scoccimarro (2009, PRD 80)

numerical realization



Schrödinger method: Widrow & Kaiser (1993, ApJ 416)
Widrow (1997, PRD 55)

Describing Cold Dark Matter



N-body picture

- N non-relativistic particles
- only gravitational interaction

$$f_N(\mathbf{x}, \mathbf{p}, \tau) = \sum_{i=1}^N \delta_D(\mathbf{x} - \mathbf{x}_i(\tau)) \delta_D(\mathbf{p} - \mathbf{p}_i(\tau))$$

Vlasov - Poisson equation

$$\partial_t f_N(\mathbf{x}, \mathbf{p}, \tau) = -\frac{\mathbf{p}}{am} \nabla_x f_N + am \nabla_x V \nabla_p f_N$$

$$\Delta V(\mathbf{x}_i) = \frac{4\pi Gm}{a} \left(\sum_{j \neq i}^N \delta_D(\mathbf{x}_i - \mathbf{x}_j) - \langle n_N \rangle \right)$$

Describing Cold Dark Matter



phase space distribution function $f(\mathbf{x}, \mathbf{p}, t)$

- **N-body**: non-relativistic, only gravitationally
- **continuous**: ensemble average, no collisions

$$f_N = \sum_i \delta_D(\mathbf{x} - \mathbf{x}_i) \delta_D(\mathbf{p} - \mathbf{p}_i)$$

f

Vlasov - Poisson equation

$$\partial_\tau f(\mathbf{x}, \mathbf{p}, \tau) = -\frac{\mathbf{p}}{am} \nabla_x f + am \nabla_x V \nabla_p f$$

↑ ↑ ↑
3+3+1
variables

partial

nonlinear

gravitational potential

$$\Delta V(\mathbf{x}, \tau) = \frac{4\pi Gm}{a} (n(\mathbf{x}, \tau) - \langle n \rangle)$$

integro

number density

$$n = \int d^3p f$$

Solving is hard!

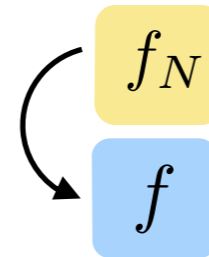
have to choose a special ansatz for $f(\mathbf{x}, \mathbf{p})$
reduce some information content

Describing Cold Dark Matter



continuous distribution function $f(\mathbf{x}, \mathbf{p}, t)$

- ensemble average
- dropping collision terms $\sim 1/N$



Vlasov - Poisson equation

$$\partial_\tau f(\mathbf{x}, \mathbf{p}, \tau) = -\frac{\mathbf{p}}{am} \nabla_{\mathbf{x}} f + am \nabla_{\mathbf{x}} V \nabla_{\mathbf{p}} f$$

gravitational potential

$$\Delta V(\mathbf{x}, \tau) = \frac{4\pi Gm}{a} (n(\mathbf{x}, \tau) - \langle n \rangle)$$

Hierarchy of Moments $M^{(n)}(\mathbf{x}) = \int d^3p p_{i_1} \dots p_{i_n} f$

- density $n(\mathbf{x})$: $M^{(0)} = n(\mathbf{x})$, velocity $\mathbf{v}(\mathbf{x})$: $M^{(1)} = n\mathbf{v}(\mathbf{x})$
- velocity dispersion $\boldsymbol{\sigma}(\mathbf{x})$: $M^{(2)} = n(\mathbf{v}\mathbf{v} + \boldsymbol{\sigma})(\mathbf{x}), \dots$ **cumulant**

$$\partial_t M^{(n)} = -\frac{1}{a^2 m} \nabla \cdot M^{(n+1)} - m \nabla V \cdot M^{(n-1)}$$

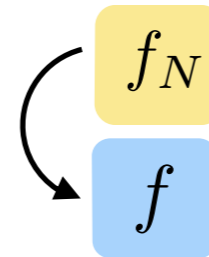
infinite **coupled hierarchy**

Describing Cold Dark Matter



phase space distribution function $f(\mathbf{t}, \mathbf{x}, \mathbf{p})$

- **N-body**: non-relativistic, only gravitationally
- **continuous**: ensemble average, no collisions



Vlasov - Poisson equation

$$\partial_\tau f(\mathbf{x}, \mathbf{p}, \tau) = -\frac{\mathbf{p}}{am} \nabla_{\mathbf{x}} f + am \nabla_{\mathbf{x}} V \nabla_{\mathbf{p}} f$$

gravitational potential

$$\Delta V(\mathbf{x}, \tau) = \frac{4\pi Gm}{a} (n(\mathbf{x}, \tau) - \langle n \rangle)$$

Hierarchy of Cumulants

- density $n(\mathbf{x})$: $C^{(0)} = \ln n(\mathbf{x})$, velocity $\mathbf{v}(\mathbf{x})$: $C^{(1)} = \mathbf{v}(\mathbf{x})$
- velocity dispersion $\sigma(\mathbf{x})$: $C^{(2)} = \sigma(\mathbf{x}), \dots$

$$\partial_t C^{(n)} = -\frac{1}{a^2 m} \left[\nabla \cdot C^{(n+1)} + \sum_{|S|=0}^n C^{(n+1-|S|)} \cdot \nabla C^{(|S|)} \right] - \delta_{n1} \cdot m \nabla V$$

consistent **truncation** $C^{(\geq n)} \equiv 0 \quad \exists n : |S| \leq 1 \wedge |S| \geq n \quad \forall S \Rightarrow n \stackrel{!}{=} 2$

Dust model



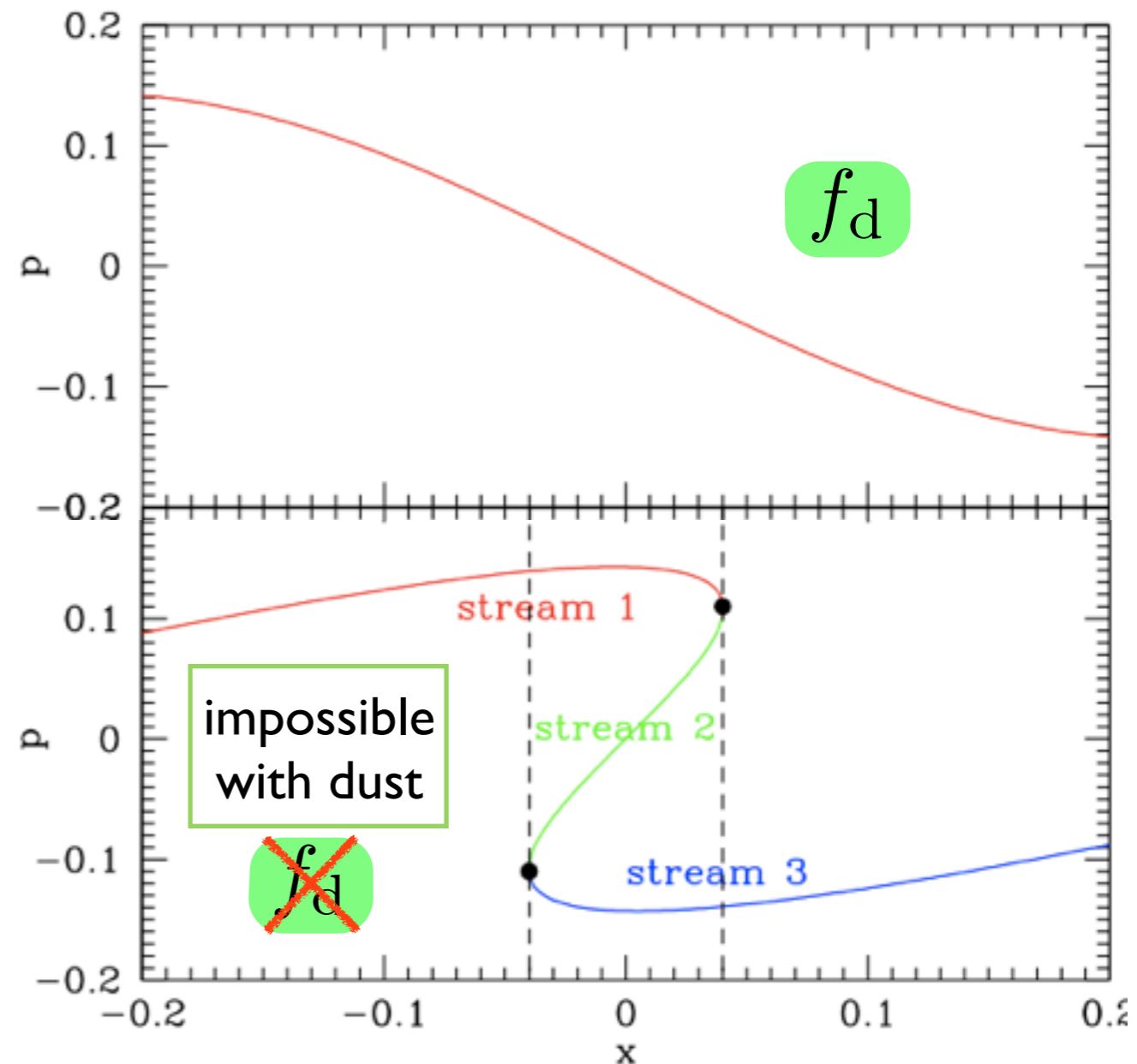
dust model

- only consistent truncation of hierarchy
- pressureless fluid: density and velocity

$$f_d(\mathbf{x}, \mathbf{p}, \tau) = n(\mathbf{x}, \tau) \delta_D^{(3)}(\mathbf{p} - \nabla \phi(\mathbf{x}, \tau))$$

$$\begin{aligned} \text{Continuity} \quad \partial_\tau n &= -\frac{1}{am} \nabla \cdot (n \nabla \phi) \\ \text{Euler} \quad \partial_\tau \phi &= -\frac{1}{2am} (\nabla \phi)^2 - amV \end{aligned}$$

- limited to **single-stream**
- no velocity dispersion, ...
- shell-crossing singularities
- **no virialization**



Schrödinger method



Schrödinger method

self-gravitating field

$$\psi = \sqrt{n} \exp\left(\frac{i}{\hbar} \phi\right)$$

Schrödinger - Poisson equation

$$i\hbar\partial_t\psi = \left[-\frac{\hbar^2}{2a^2m}\Delta + mV \right] \psi$$
$$\Delta V = \frac{4\pi G\rho_0}{a} (|\psi|^2 - 1)$$



Schrödinger method



Schrödinger method

Wigner function, constructed from self-gravitating field

$$\psi = \sqrt{n} \exp\left(\frac{i}{\hbar} \phi\right)$$

$$f_W(\mathbf{x}, \mathbf{p}) = \int \frac{d^3 \tilde{\mathbf{x}}}{(2\pi\hbar)^3} \exp\left[2\frac{i}{\hbar} \mathbf{p}' \cdot \tilde{\mathbf{x}}\right] \psi(\mathbf{x}' - \tilde{\mathbf{x}}) \bar{\psi}(\mathbf{x}' + \tilde{\mathbf{x}})$$

Schrödinger - Poisson equation

Wigner-Vlasov equation

$$\partial_t f_W = \left[\frac{\mathbf{p}^2}{2a^2 m} + mV \right] \frac{2}{\hbar} \sin\left(\frac{\hbar}{2} (\overleftarrow{\nabla}_x \overrightarrow{\nabla}_p - \overleftarrow{\nabla}_p \overrightarrow{\nabla}_x)\right) f_W$$

$$i\hbar \partial_t \psi = \left[-\frac{\hbar^2}{2a^2 m} \Delta + mV \right] \psi$$
$$\Delta V = \frac{4\pi G \rho_0}{a} (|\psi|^2 - 1)$$

Vlasov equation

$$\partial_t f = \left[\frac{\mathbf{p}^2}{2a^2 m} + mV \right] (\overleftarrow{\nabla}_x \overrightarrow{\nabla}_p - \overleftarrow{\nabla}_p \overrightarrow{\nabla}_x) f$$

Problems

- Wigner distribution function not manifestly positive
- time evolution not in good correspondence to Vlasov



Schrödinger method



Schrödinger method

- Coarse-grained Wigner function, constructed from self-gravitating field

$$\psi = \sqrt{n} \exp\left(\frac{i}{\hbar} \phi\right)$$

$$f_{cg}^W(\mathbf{x}, \mathbf{p}) = \int \frac{d^3 x' d^3 p'}{(\pi \sigma_x \sigma_p)^3} \exp\left[-\frac{(\mathbf{x} - \mathbf{x}')^2}{2\sigma_x^2} - \frac{(\mathbf{p} - \mathbf{p}')^2}{2\sigma_p^2}\right] \int \frac{d^3 \tilde{x}}{(2\pi\hbar)^3} \exp\left[2\frac{i}{\hbar} \mathbf{p}' \cdot \tilde{x}\right] \psi(\mathbf{x}' - \tilde{x}) \bar{\psi}(\mathbf{x}' + \tilde{x})$$

Schrödinger - Poisson equation

Wigner-Vlasov equation

$$\partial_t f_W = \left[\frac{\mathbf{p}^2}{2a^2 m} + mV \right] \frac{2}{\hbar} \sin\left(\frac{\hbar}{2} (\overleftarrow{\nabla}_x \overrightarrow{\nabla}_p - \overleftarrow{\nabla}_p \overrightarrow{\nabla}_x)\right) f_W$$

$$i\hbar \partial_t \psi = \left[-\frac{\hbar^2}{2a^2 m} \Delta + mV \right] \psi$$
$$\Delta V = \frac{4\pi G \rho_0}{a} (|\psi|^2 - 1)$$

Vlasov equation

$$\partial_t f = \left[\frac{\mathbf{p}^2}{2a^2 m} + mV \right] (\overleftarrow{\nabla}_x \overrightarrow{\nabla}_p - \overleftarrow{\nabla}_p \overrightarrow{\nabla}_x) f$$

Problems

- Wigner distribution function not manifestly positive
- time evolution not in good correspondence to Vlasov
- solution: add a coarse-graining $\sigma_x \sigma_p \gtrsim \hbar/2$



Schrödinger method



Schrödinger method

- Coarse-grained Wigner function, constructed from self-gravitating field

$$\psi = \sqrt{n} \exp\left(\frac{i}{\hbar} \phi\right)$$

$$f_{cg}^W(\mathbf{x}, \mathbf{p}) = \int \frac{d^3 x' d^3 p'}{(\pi \sigma_x \sigma_p)^3} \exp\left[-\frac{(\mathbf{x} - \mathbf{x}')^2}{2\sigma_x^2} - \frac{(\mathbf{p} - \mathbf{p}')^2}{2\sigma_p^2}\right] \int \frac{d^3 \tilde{x}}{(2\pi\hbar)^3} \exp\left[2\frac{i}{\hbar} \mathbf{p}' \cdot \tilde{x}\right] \psi(\mathbf{x}' - \tilde{x}) \bar{\psi}(\mathbf{x}' + \tilde{x})$$

Schrödinger - Poisson equation

$$i\hbar \partial_t \psi = \left[-\frac{\hbar^2}{2a^2 m} \Delta + mV \right] \psi$$

$$\Delta V = \frac{4\pi G \rho_0}{a} (|\psi|^2 - 1)$$

degrees of freedom

- 2: amplitude n & phase ϕ

parameters

- coarse-graining σ_x, σ_p
 - fundamental resolution $\sigma_x \sigma_p \gtrsim \hbar/2$
- Schrödinger scale \hbar
 - degree of restriction
 - dust as special case

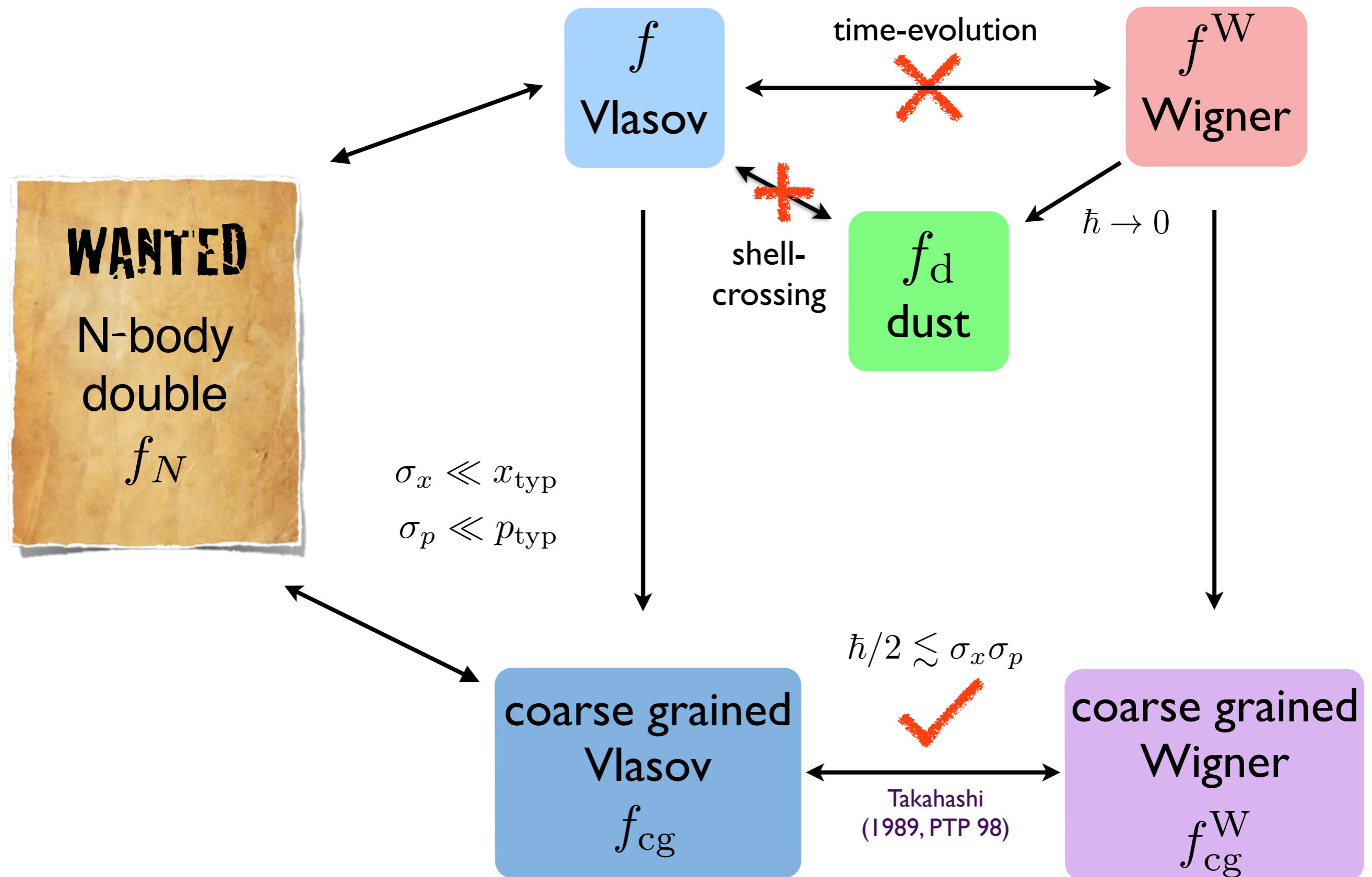
Continuity $\partial_\tau n = -\frac{1}{am} \nabla(n \nabla \phi)$

Euler $\partial_\tau \phi = -\frac{1}{2am} (\nabla \phi)^2 - amV$

quantum potential

$$+\frac{\hbar^2}{2am} \left(\frac{\Delta \sqrt{n}}{\sqrt{n}} \right)$$

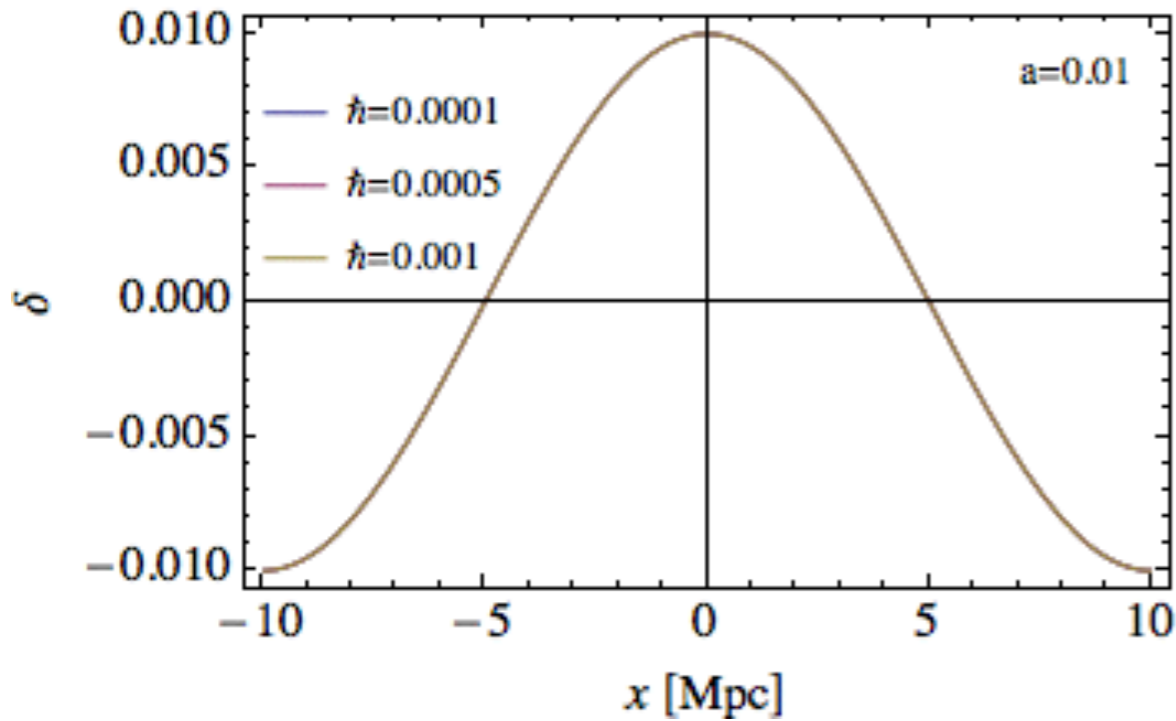
Schrödinger method at a glance



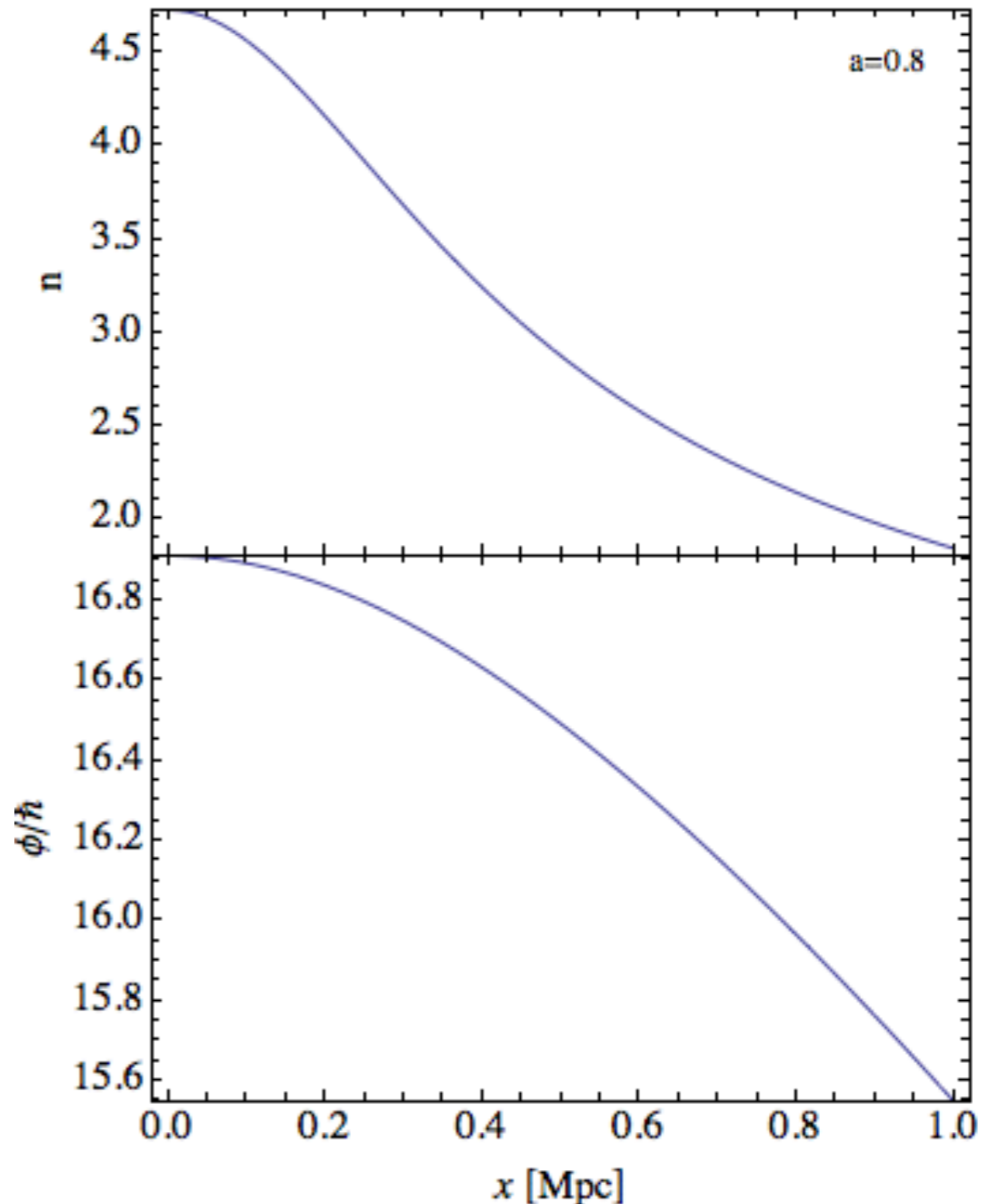
Features of Schrödinger Method



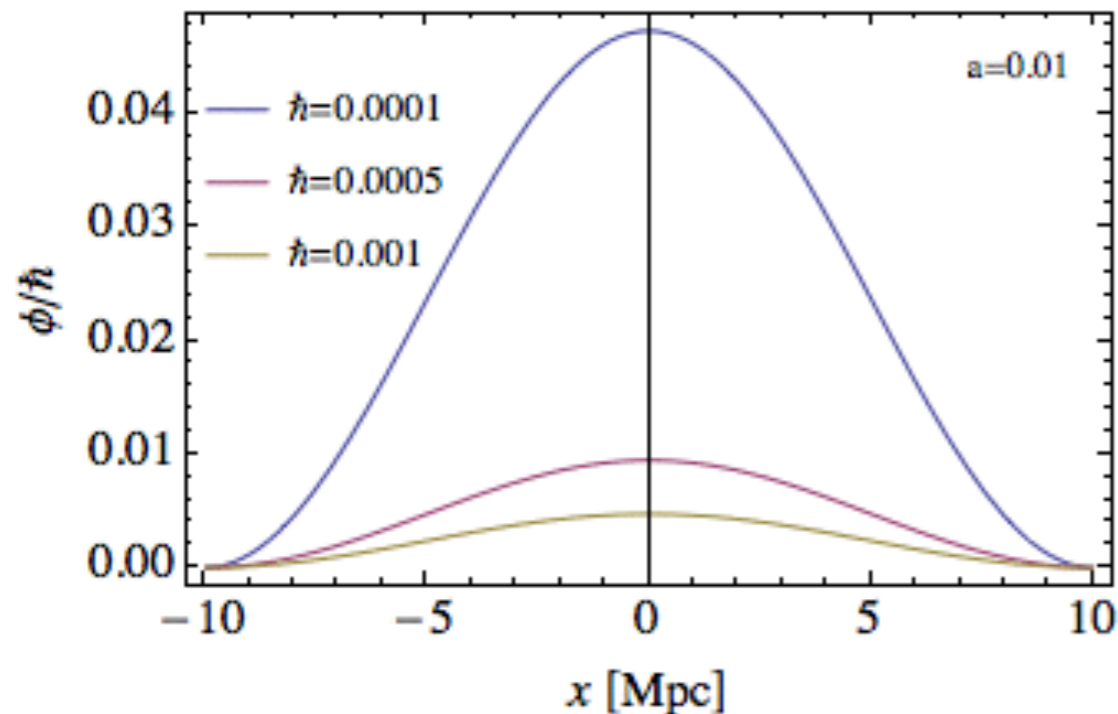
✓ prevention of shell-crossing singularities



!! $\psi = \sqrt{n}e^{i\phi/\hbar}$ free of pathologies



!?! occurrence of phase jumps



Features of Schrödinger Method



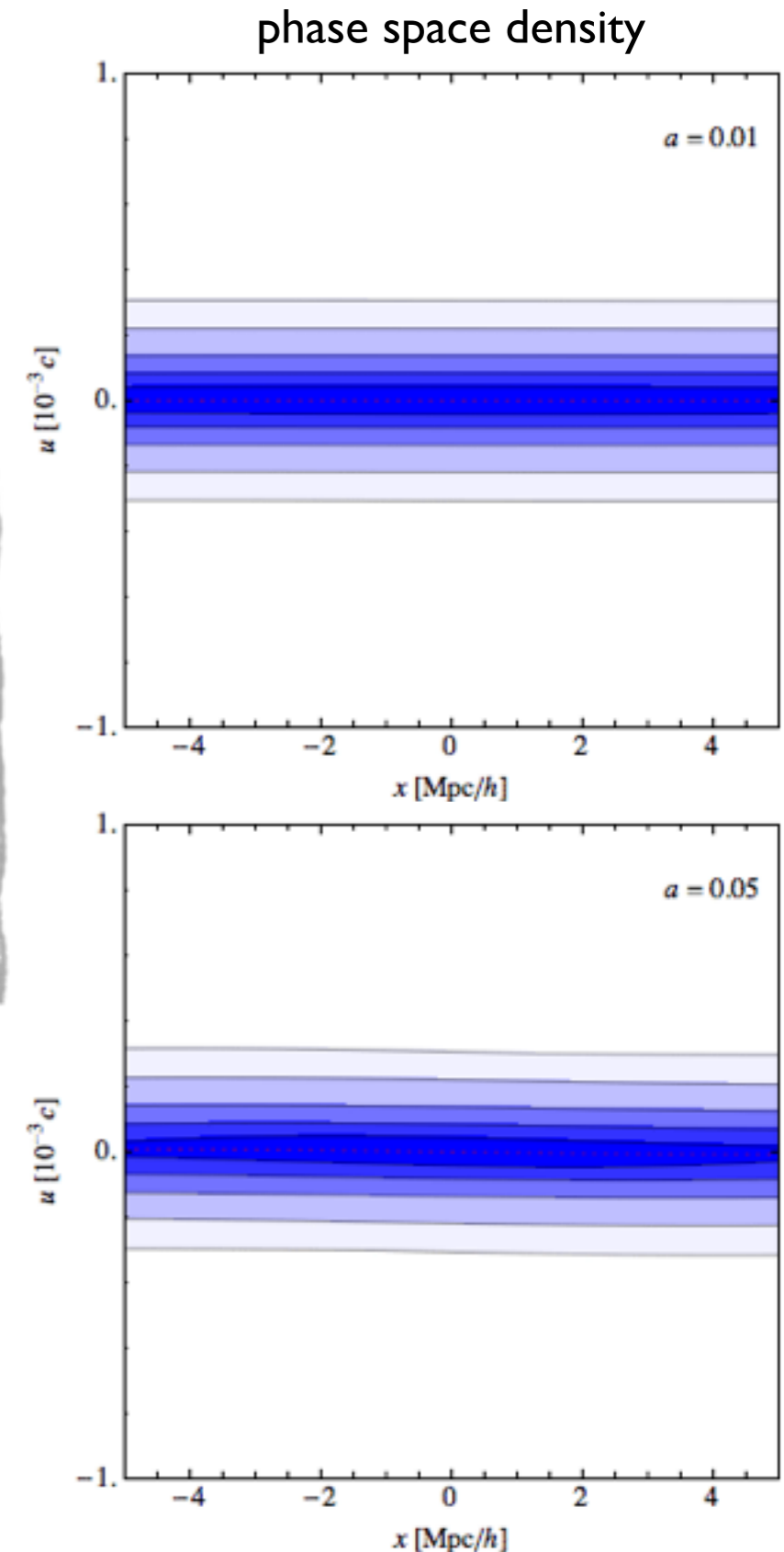
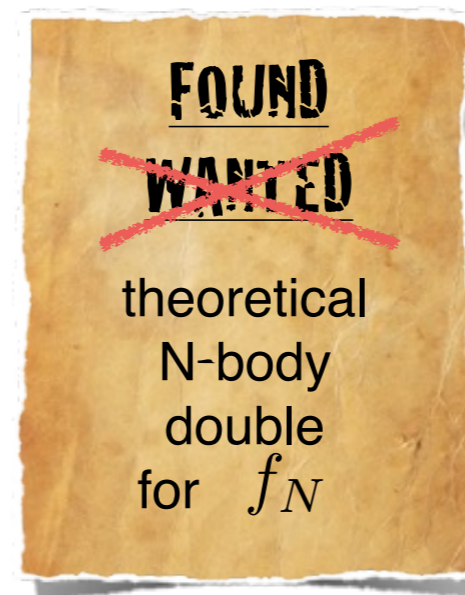
Multi-streaming

- ✗ dust model: fails at shell-crossing
- ✓ Schrödinger method: can go **beyond shell-crossing**

blue S contours: Schrödinger method
red dotted Z line: Zeldovich solution (dust model)

Virialization

- ✗ even in extended models: no **virialization**
- ✓ Schrödinger method: **bound structures like halos**



Features of Schrödinger Method



Schrödinger method

- Coarse-grained Wigner function, constructed from self-gravitating field

$$\psi = \sqrt{n} \exp\left(\frac{i}{\hbar} \phi\right)$$

$$f_{cg}^W(\mathbf{x}, \mathbf{p}) = \int \frac{d^3 x' d^3 p'}{(\pi \sigma_x \sigma_p)^3} \exp\left[-\frac{(\mathbf{x} - \mathbf{x}')^2}{2\sigma_x^2} - \frac{(\mathbf{p} - \mathbf{p}')^2}{2\sigma_p^2}\right] \int \frac{d^3 \tilde{x}}{(2\pi\hbar)^3} \exp\left[2\frac{i}{\hbar} \mathbf{p}' \cdot \tilde{x}\right] \psi(\mathbf{x}' - \tilde{x}) \bar{\psi}(\mathbf{x}' + \tilde{x})$$

Cumulants

special p-dependence

- lowest two: macroscopic density & velocity

$$\bar{n}(\mathbf{x}) = \exp\left[\frac{1}{2} \sigma_x^2 \Delta\right] n(\mathbf{x}) \quad \bar{\mathbf{v}}(\mathbf{x}) = \frac{1}{am\bar{n}(\mathbf{x})} \exp\left[\frac{1}{2} \sigma_x^2 \Delta\right] (n \nabla \phi)(\mathbf{x})$$

- higher cumulants given self-consistently
evolution equations fulfilled automatically

$$C^{(0)} = \ln n, \quad C^{(1)} = \nabla \phi$$

$$C^{(n+2)} = -\frac{\hbar^2}{4} \nabla \nabla C^{(n)}$$

closure of hierarchy

CU, Kopp & Haugg (2014, PRD 90, 023517)

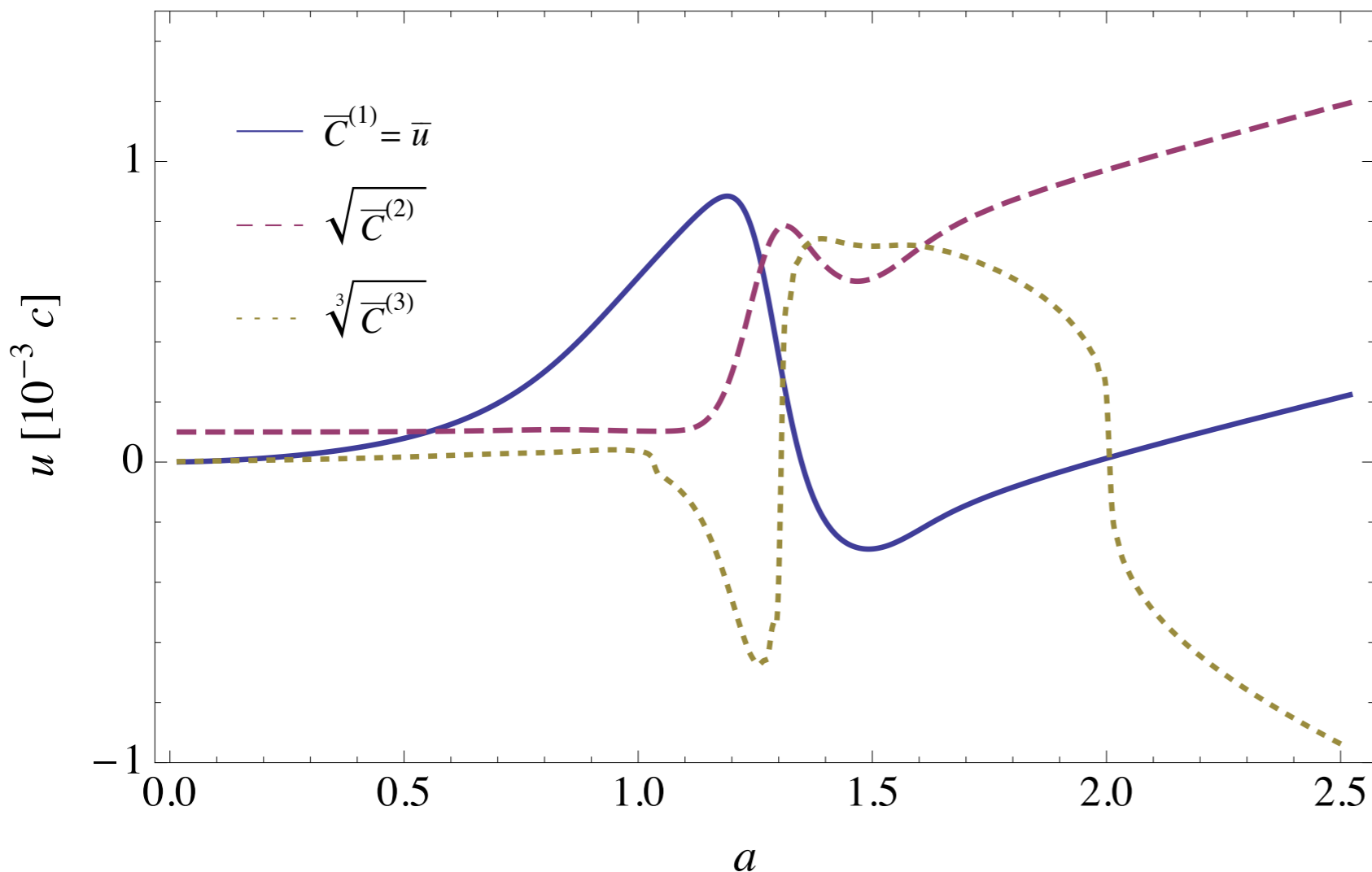
add coarse-graining on top of that

Features of Schrödinger Method

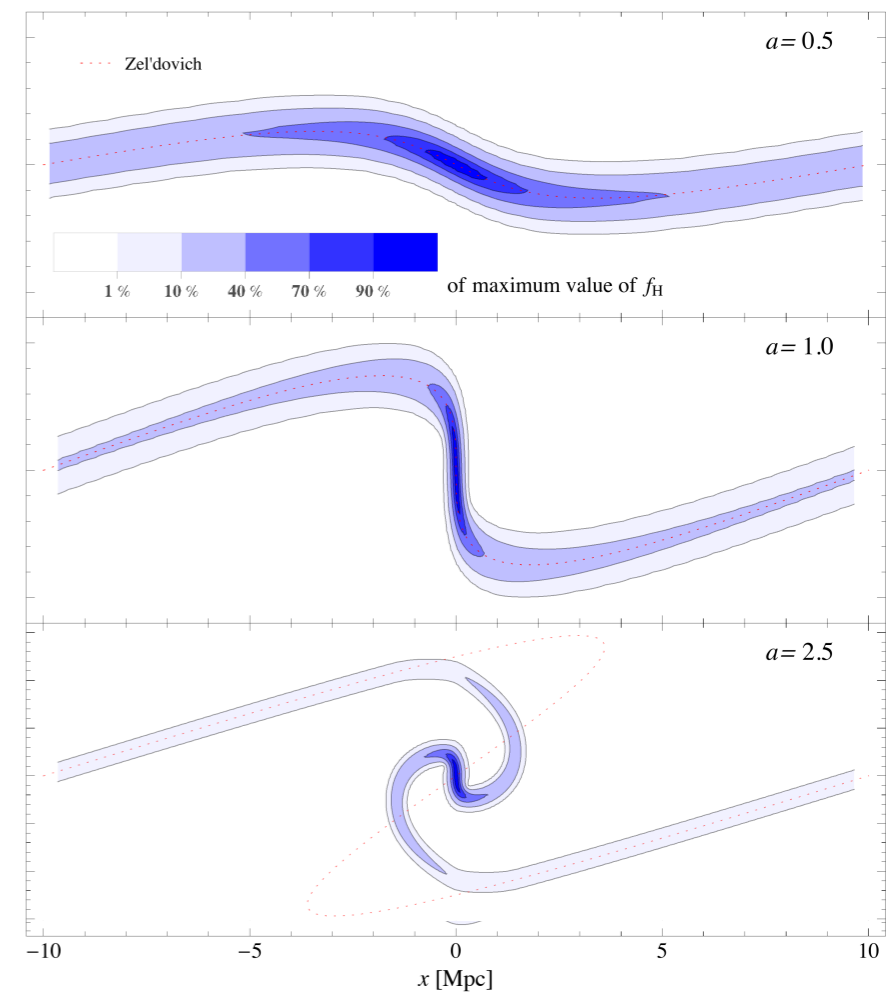


Multi-streaming

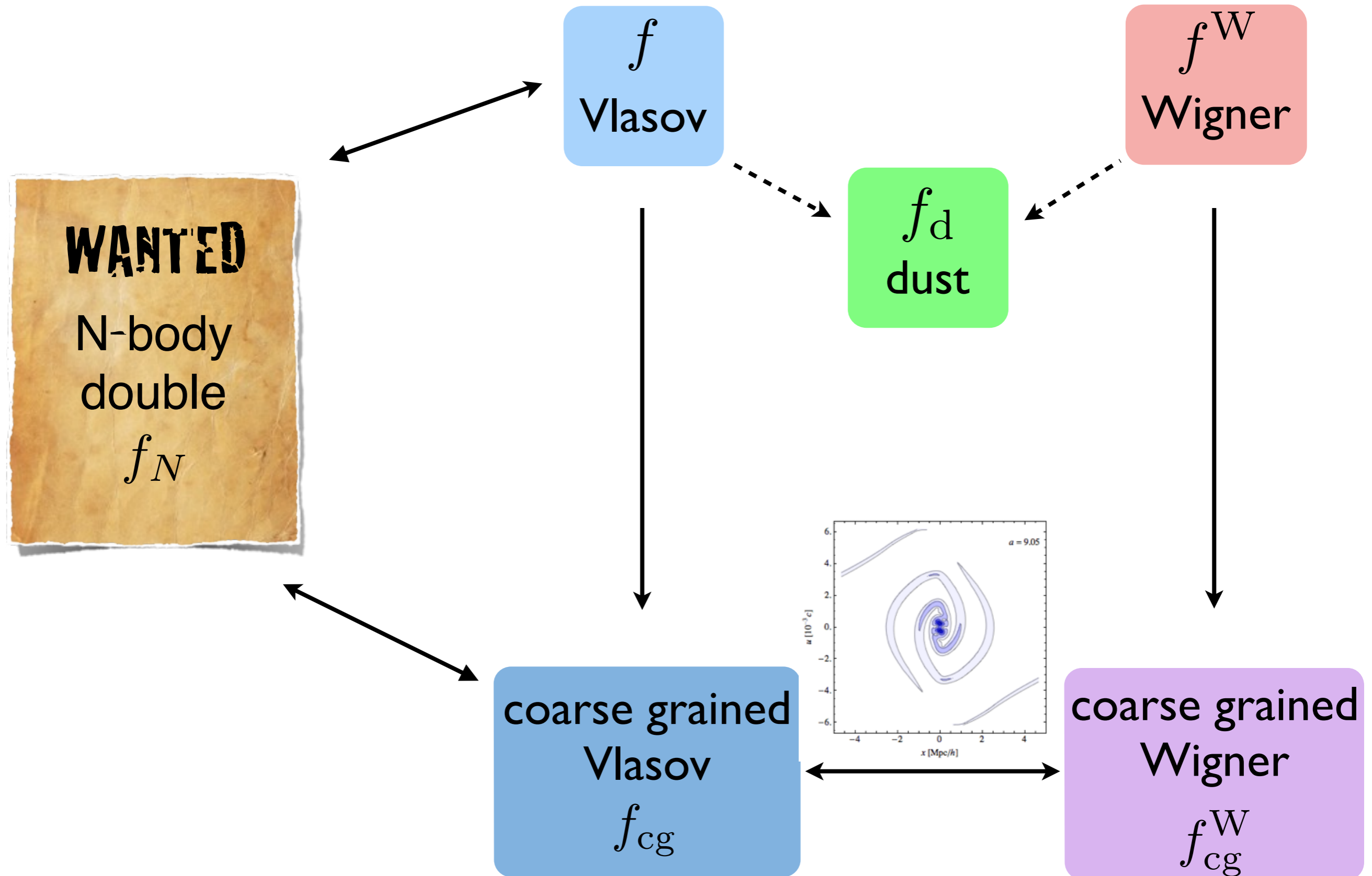
- higher cumulants encode multi-streaming effects
- during shell-crossing: higher moments sourced dynamically



Schrödinger method: cumulants at $x = -0.5$ Mpc:
all equally important after shell crossing



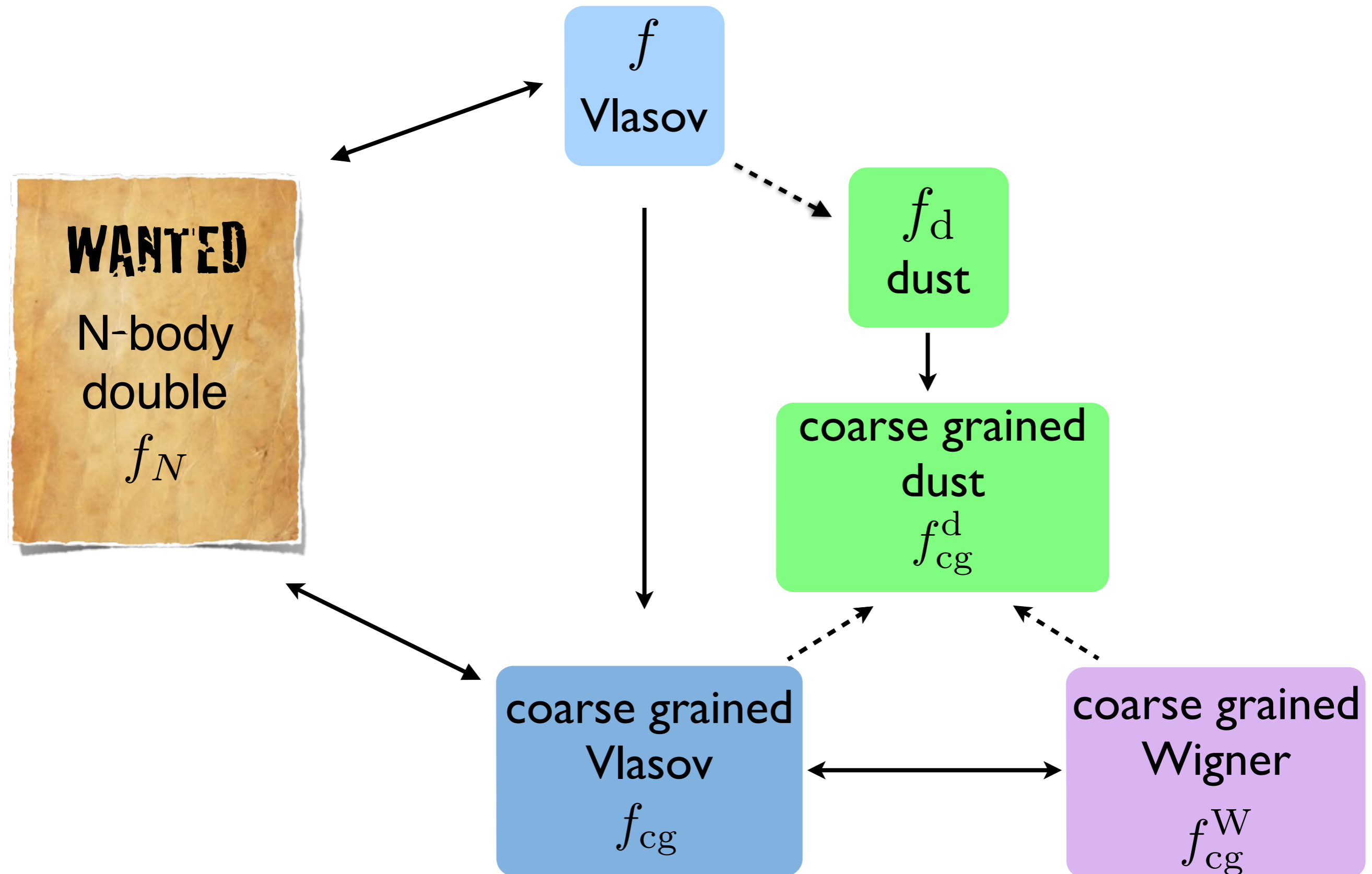
Schrödinger method at a glance





Dark Matter power spectrum with the Coarse-grained dust model

Coarse-grained perturbation theory



Eulerian Perturbation Theory



Dust model

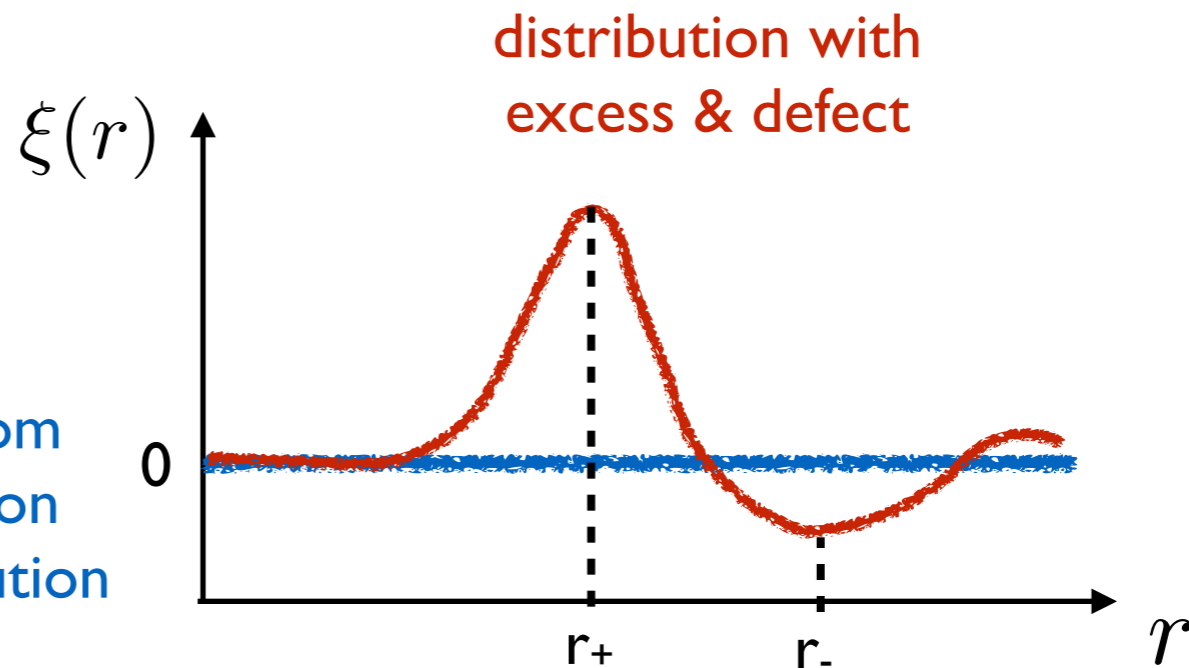
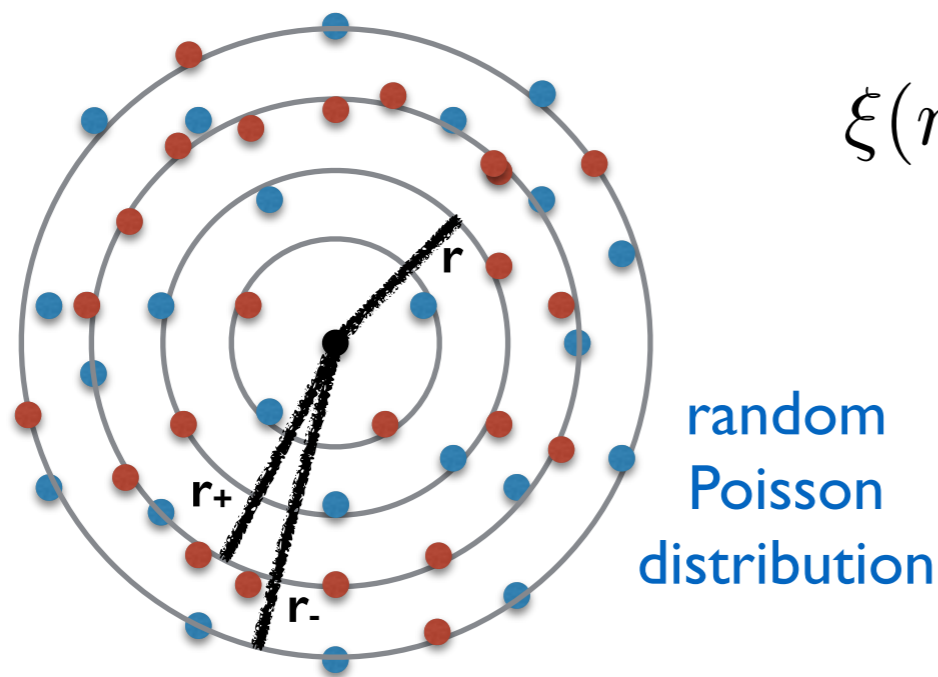
- express fluid equations in terms of $\delta = n - 1$ and $\theta = \nabla \cdot \mathbf{v} \propto \Delta\phi$ (no vorticity)
- perturbative expansion: separation ansatz (fastest growing mode)

$$\delta(\tau, \mathbf{k}) = \sum_{n=1}^{\infty} a^n(\tau) \delta_n(\mathbf{k}) \quad \theta(\tau, \mathbf{k}) = \mathcal{H}(\tau) \sum_{n=1}^{\infty} a^n(\tau) \theta_n(\mathbf{k})$$

Correlation function

- 2-point correlation: excess probability of finding 2 objects separated by r

$$dP = n[1 + \xi(r)]dV \quad \text{homogeneity \& isotropy: } \xi(\mathbf{r}) = \xi(r)$$



Eulerian Perturbation Theory

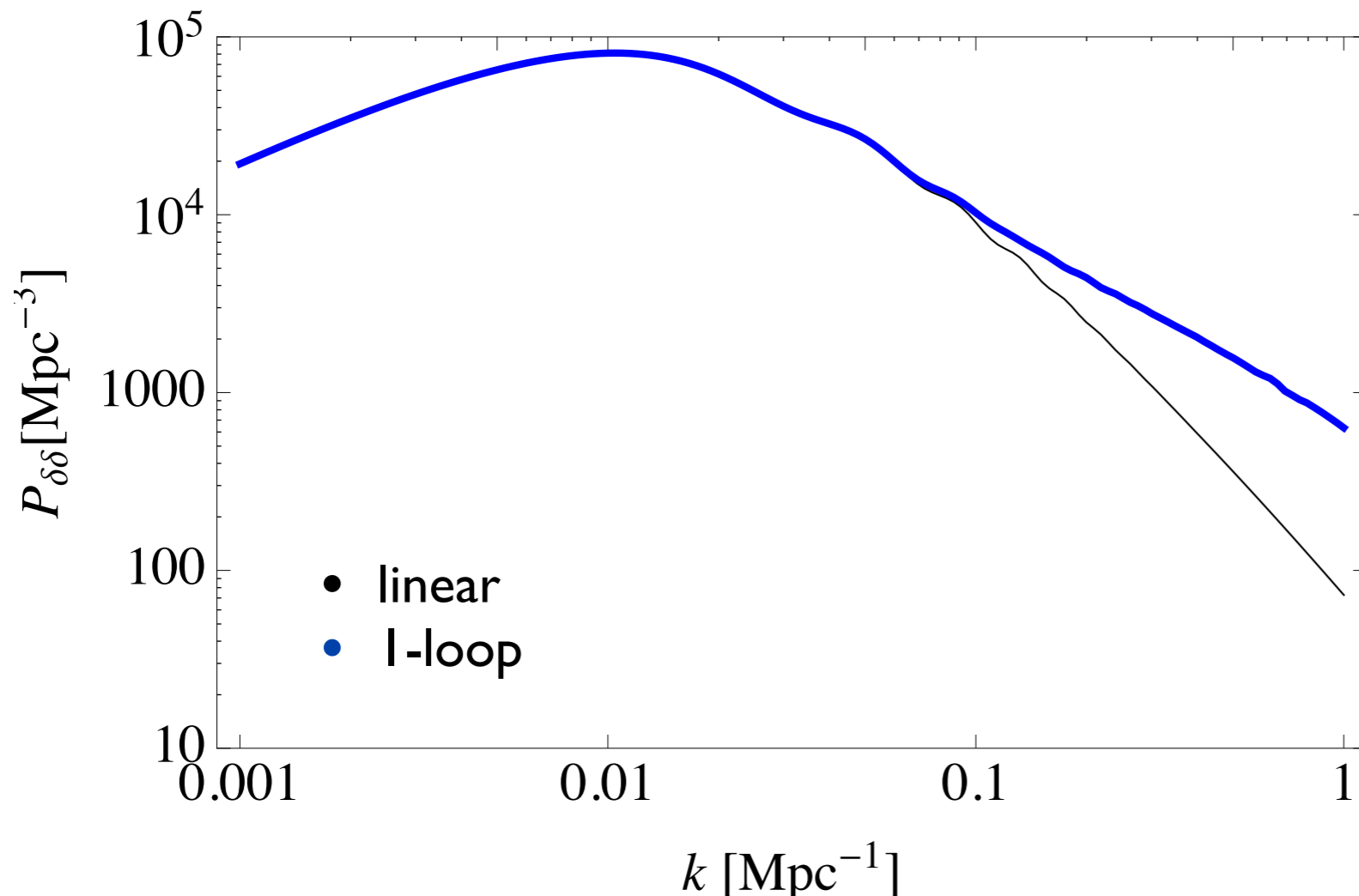


Dust model

- express fluid equations in terms of $\delta = n - 1$ and $\theta = \nabla \cdot \mathbf{v} = \Delta\phi$ (no vorticity)
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Density power spectrum



correlation function
=
FT of power spectrum

$$\xi(r) = \frac{1}{2\pi^2} \int dk k^2 P(k) \frac{\sin(kr)}{kr}$$

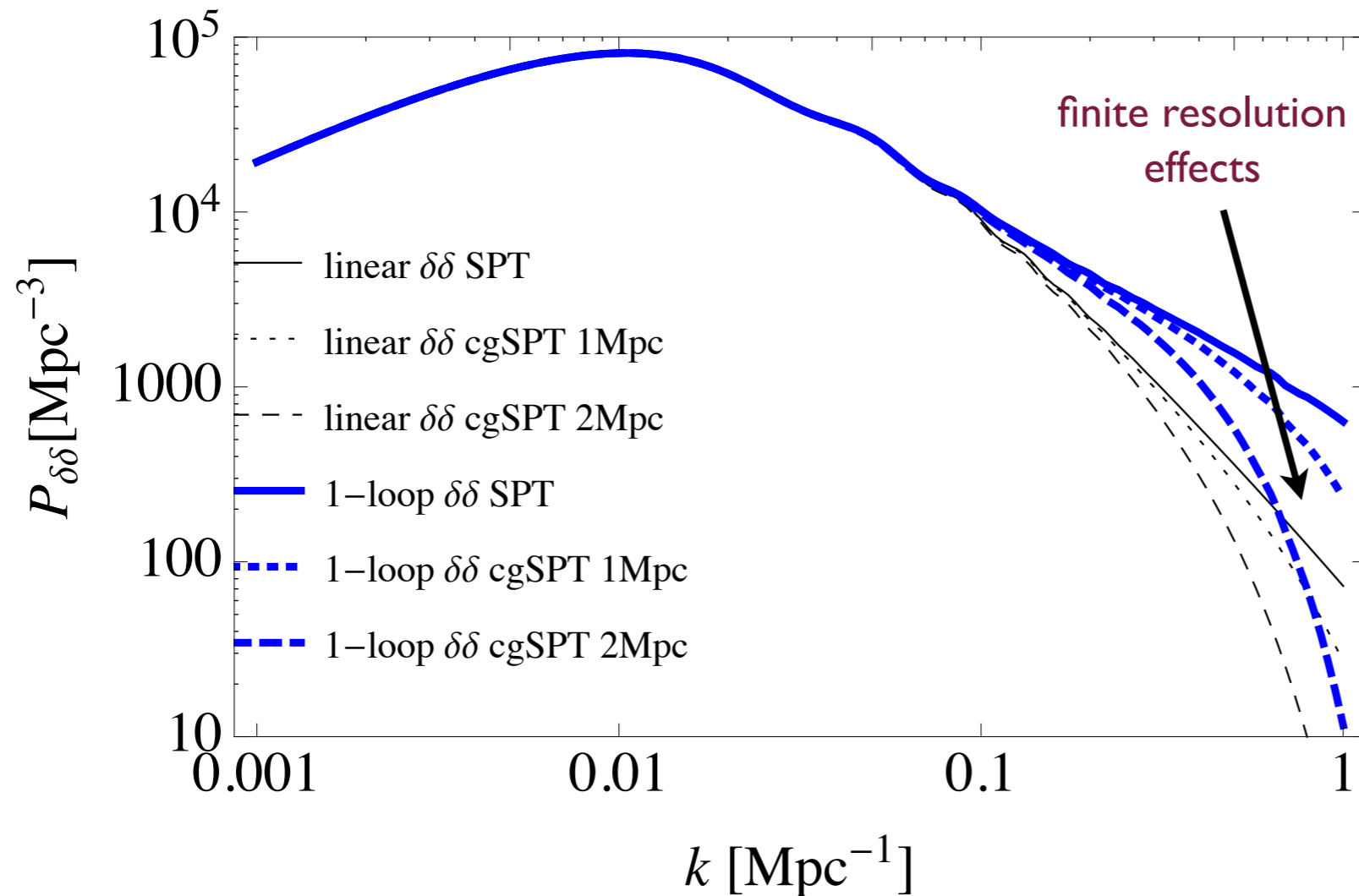
Eulerian Perturbation Theory



Coarse grained dust model

- consider only σ_x correction in Schrödinger method
- in 1st order: smoothing of input power spectrum

Density power spectrum



trivial effect:

density power spectrum
gets smoothed

Eulerian Perturbation Theory

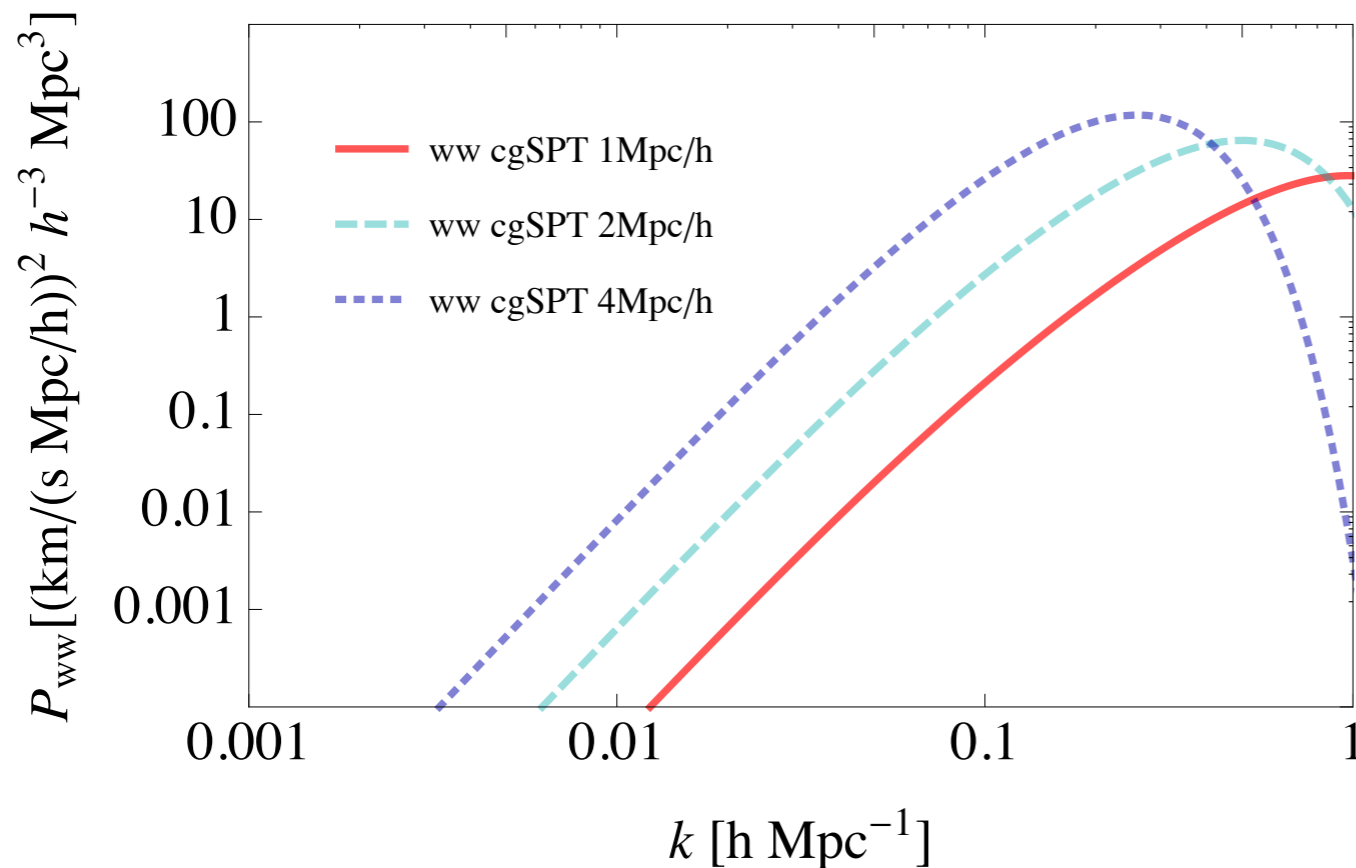


Coarse grained dust model

- similar to dust, but mass-weighted velocity $\bar{v} := \frac{\overline{nv}}{\bar{n}}$
- large scale vorticity $\bar{w} := \nabla \times \bar{v} \neq 0$!

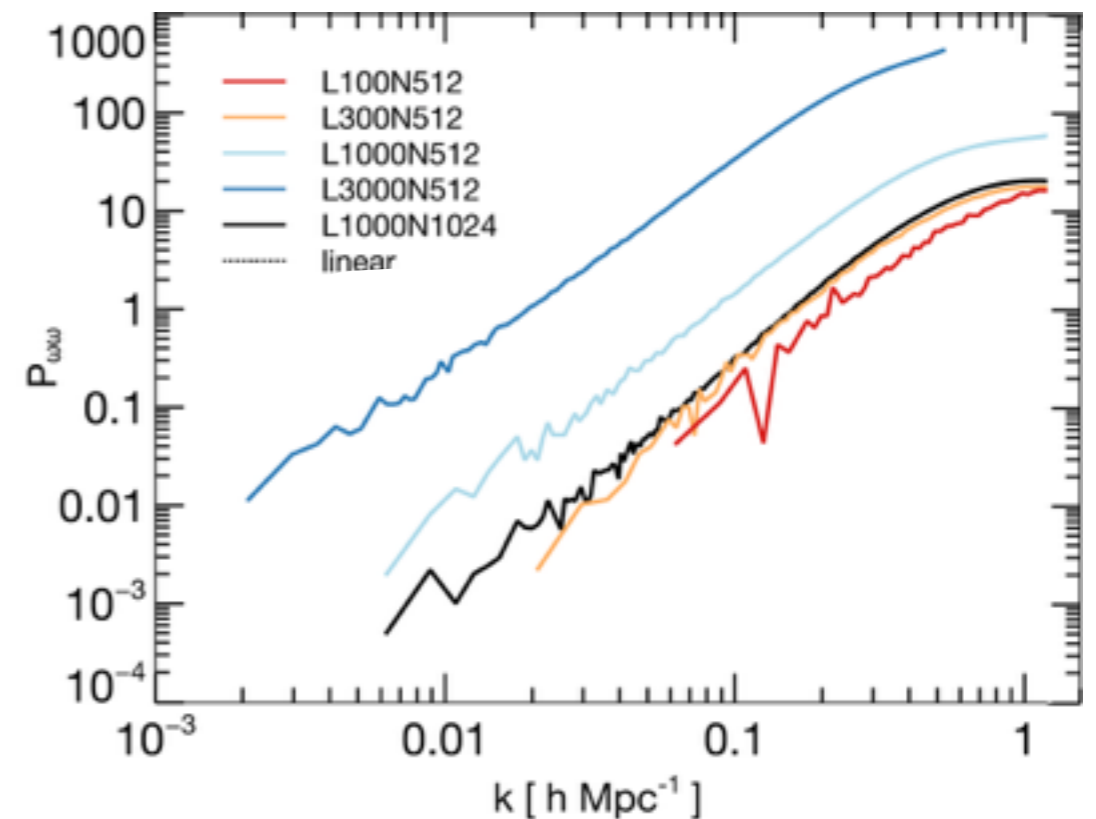
Vorticity power spectrum $P_{ww}(k)$

CU & Kopp (2015, PRD 91, 084010, arXiv: 1407.4810)



corresponding N-body data

Hahn, Angulo & Abel (2014, arXiv: 1404.2280)



Conclusion & Prospects

Schrödinger method

- models CDM using a self-gravitating scalar field
- analytical tool to access nonlinear stage of structure formation
 - describes multi-streaming & allows for virialization

CU, Kopp, Haugg (2014, PRD 90, 023517, arXiv: 1403.5567)

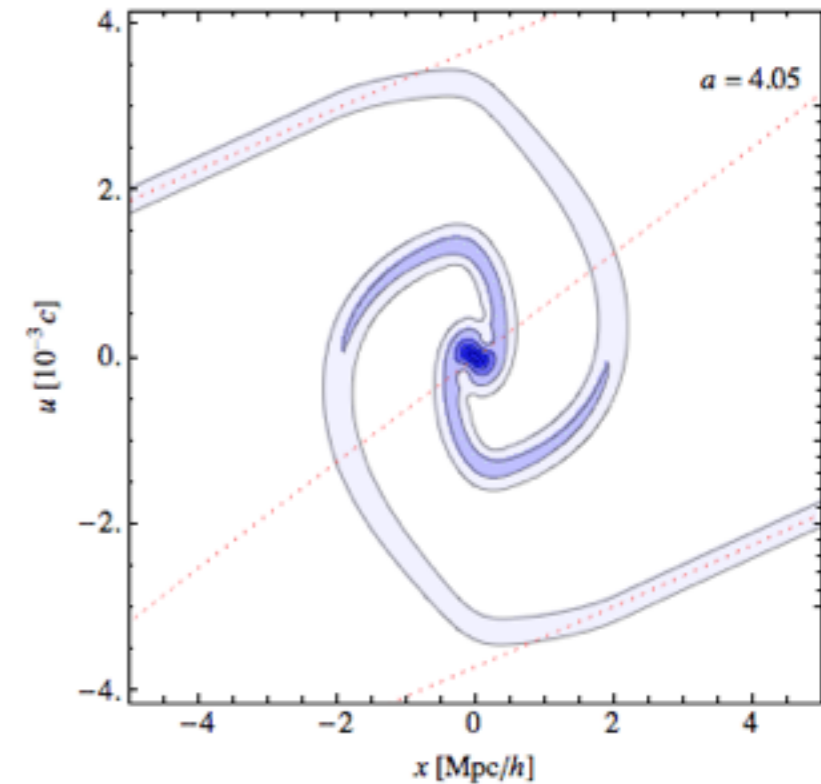
Coarse-grained dust model

- mass-weighted velocity via Gaussian smoothing
 - vorticity compatible with N-body

CU & Kopp (2015, PRD 91, 084010, arXiv: 1407.4810)

Prospects

- Schrödinger method in terms of a filtering of the distribution function
 - approximate hierarchy closures instead of exact truncation
 - comparison of moments 1D Schrödinger vs. 1D Vlasov-Poisson
 - disentangle limitations of pressureless fluid & perturbation theory
- consider the effect of varying phase-space resolution \hbar
- understand universal density profiles of halos (NFW)





Thank You for Your Attention

Questions?