

Where do AMS-02 anti-helium events come from?

Vivian Poulin

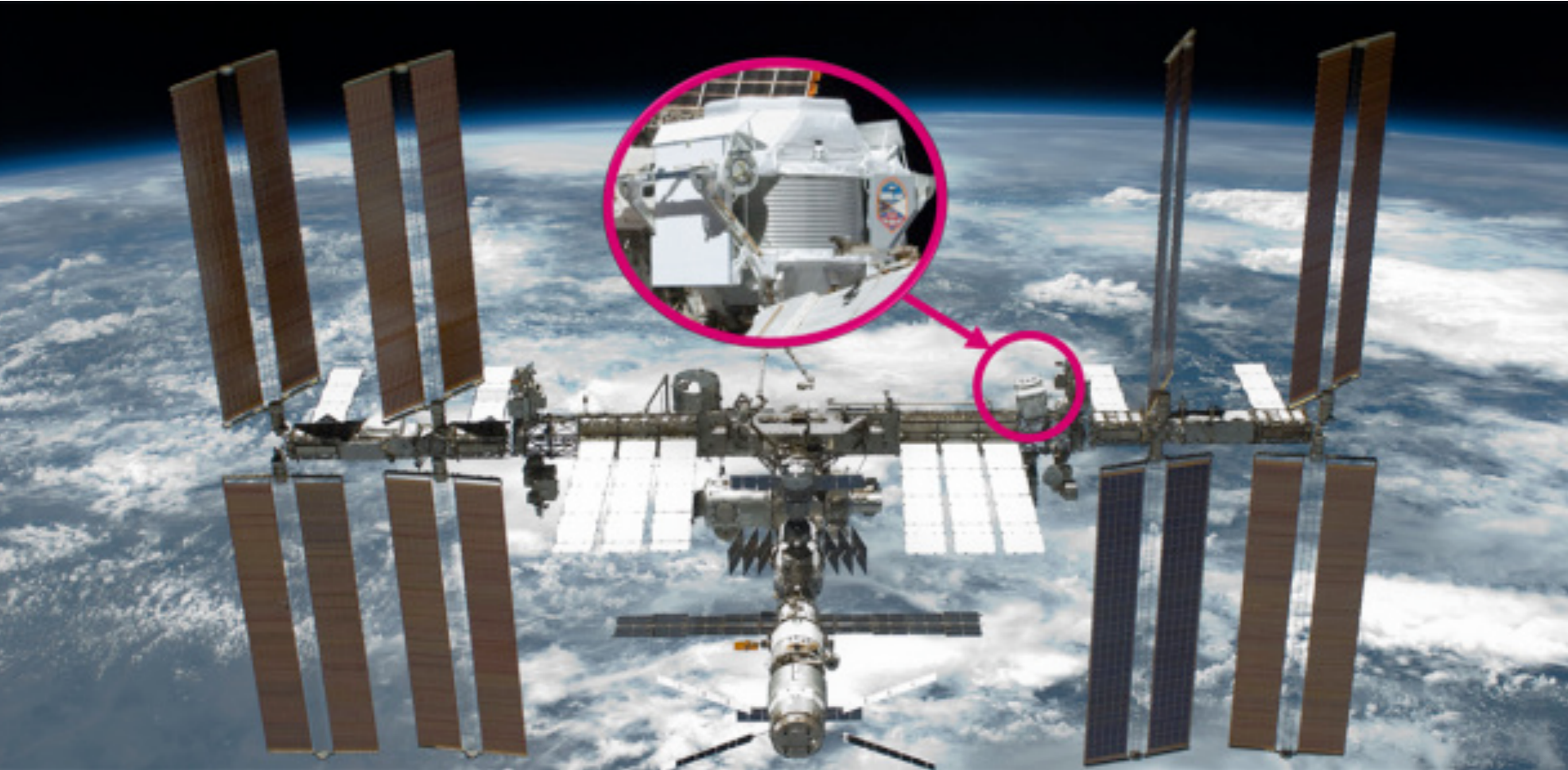
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w/ P. Salati, I. Cholis, M. Kamionkowski and J. Silk
Phys.Rev. D99 (2019) no.2, 023016

GReCO seminar series
IAP, 27/01/2020

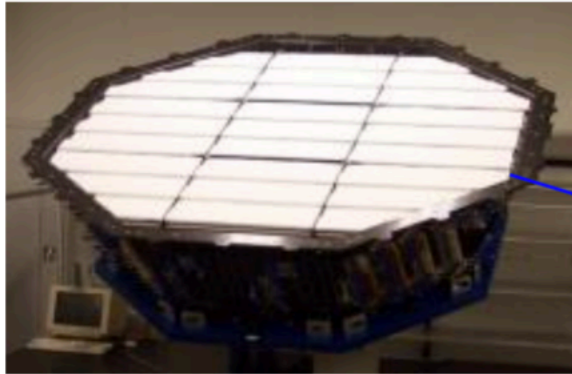


AMS: A particle physics detector on the ISS



AMS: A TeV precision, magnetic spectrometer

Transition Radiation Detector
Identify e^+ , e^-



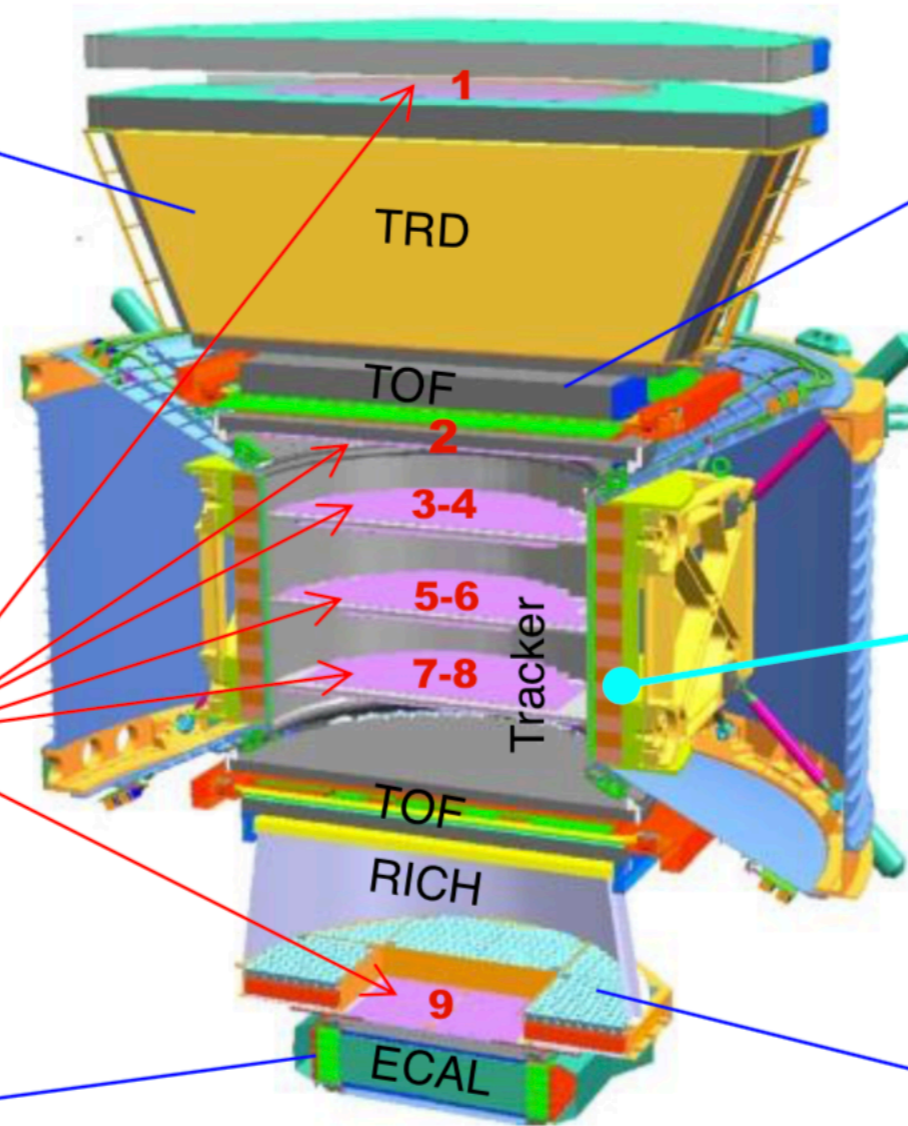
Silicon Tracker
 Z, R



Electromagnetic Calorimeter
 E of e^+ , e^-



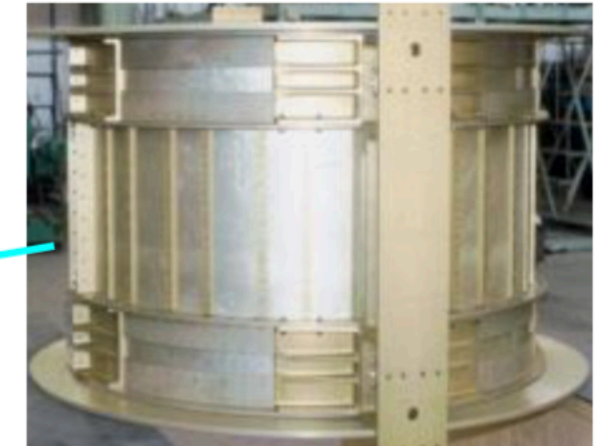
Particles and nuclei are defined by their charge (Z) and energy (E) or rigidity ($R=p/Z$)



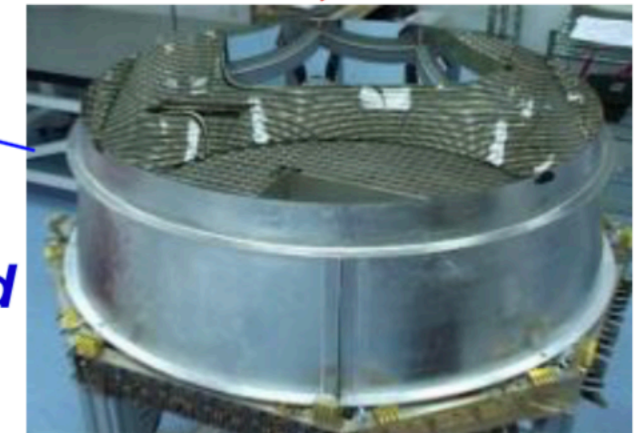
Time of Flight
 Z, E



Magnet
 $\pm Z, R$



Ring Imaging Cherenkov
 Z, E



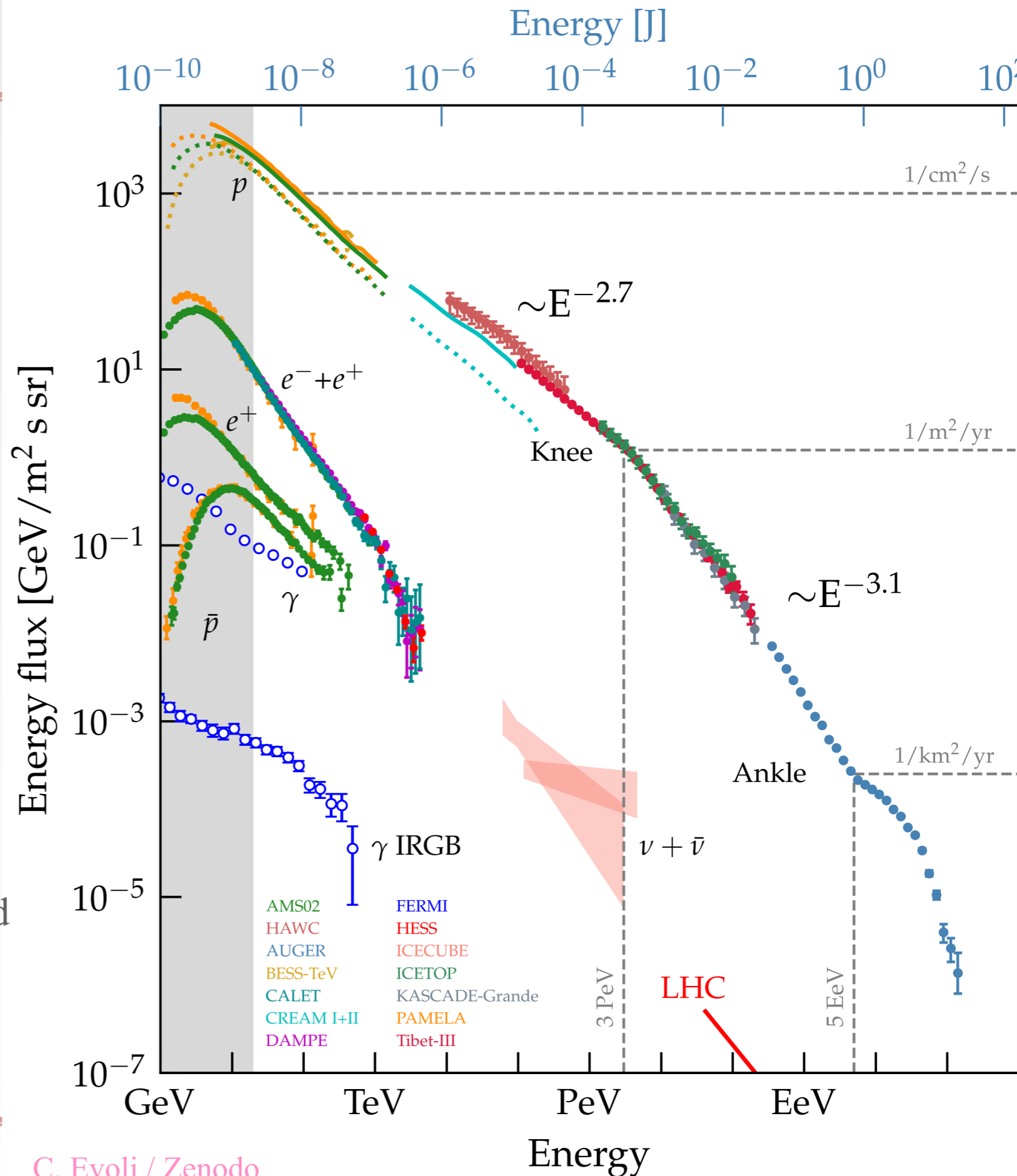
The Charge and Energy are measured independently by many detectors

Slide from V. Choutko

The CR spectrum

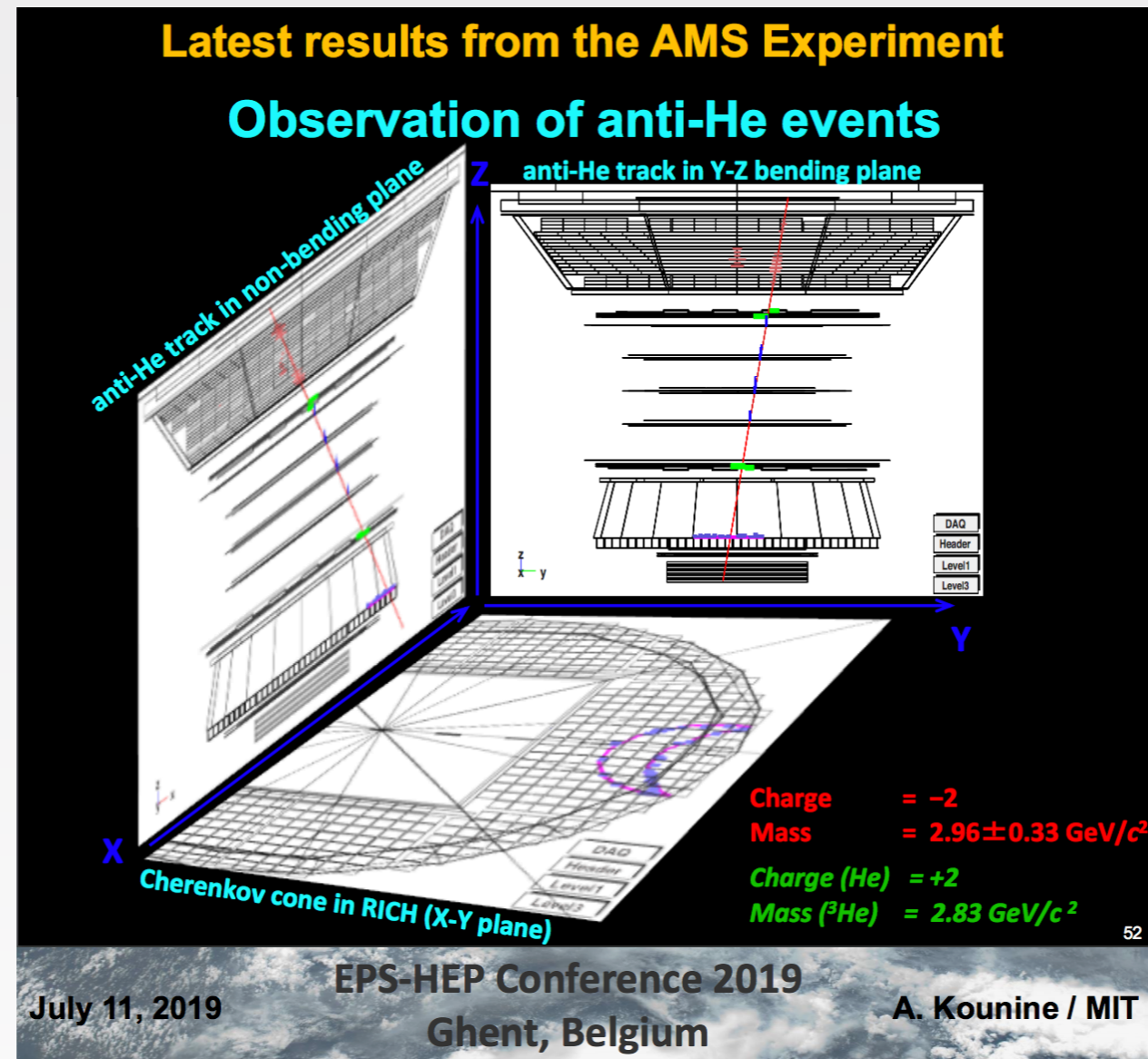
- Largely dominated by protons
 $100p : 10\text{He} : 1e^- : 10^{-2}e^+ : 10^{-3}\bar{p}$
- AMS measures at MeV-TeV energies
 \Rightarrow Galactic sources.
- AMS provides the most accurate measurement of CR fluxes to date in the MeV-TeV range.
- AMS can teach us about astrophysical sources (e.g. pulsars), Dark Matter and other exotic sources of antimatter.

V. Poulin - LUPM (CNRS)



C. Evoli / Zenodo

Did AMS see anti-helium events?



- AMS-02 might have identified $6 \overline{{}^3\text{He}}$ and $2 \overline{{}^4\text{He}}$. The event rate is $\sim 1 \overline{\text{He}}$ for 10^8He .
- Massive Monte Carlo simulations are carried out to evaluate significance.
- Current event rate is ~ 20 times above the claimed sensitivity.

Kounine, ICRC 2011

Anti-matter in the universe

- Baryon asymmetry in the universe is defined: $B \equiv \frac{n_B - n_{\bar{B}}}{s_\gamma} \sim \eta \equiv \frac{n_B}{n_\gamma} \Big|_{\text{today}}$ *e.g. Kolb&Turner's book*

- Assuming homogeneous, baryon-symmetric universe and no B-violation processes

$$\frac{n_B}{n_\gamma} = \frac{n_{\bar{B}}}{n_\gamma} \sim 10^{-17}$$

- From BBN and CMB we know $\eta \equiv \frac{n_B}{n_\gamma} \sim 6 \cdot 10^{-10} \Rightarrow B > 0$

- In our vicinity: $n_{\bar{p}} \sim 10^{-5} n_p$ in cosmic rays, most likely purely secondaries.

- Why is there so much more matter than this naive prediction? Where is the antimatter?

- AMS has detected anti-helium: How can such objects be created? anti-BBN? anti-stars?

Summary

A single ${}^4\overline{\text{He}}$ could indicate the presence of anti-objects.

I/ Anti-helium flux from standard astrophysical processes

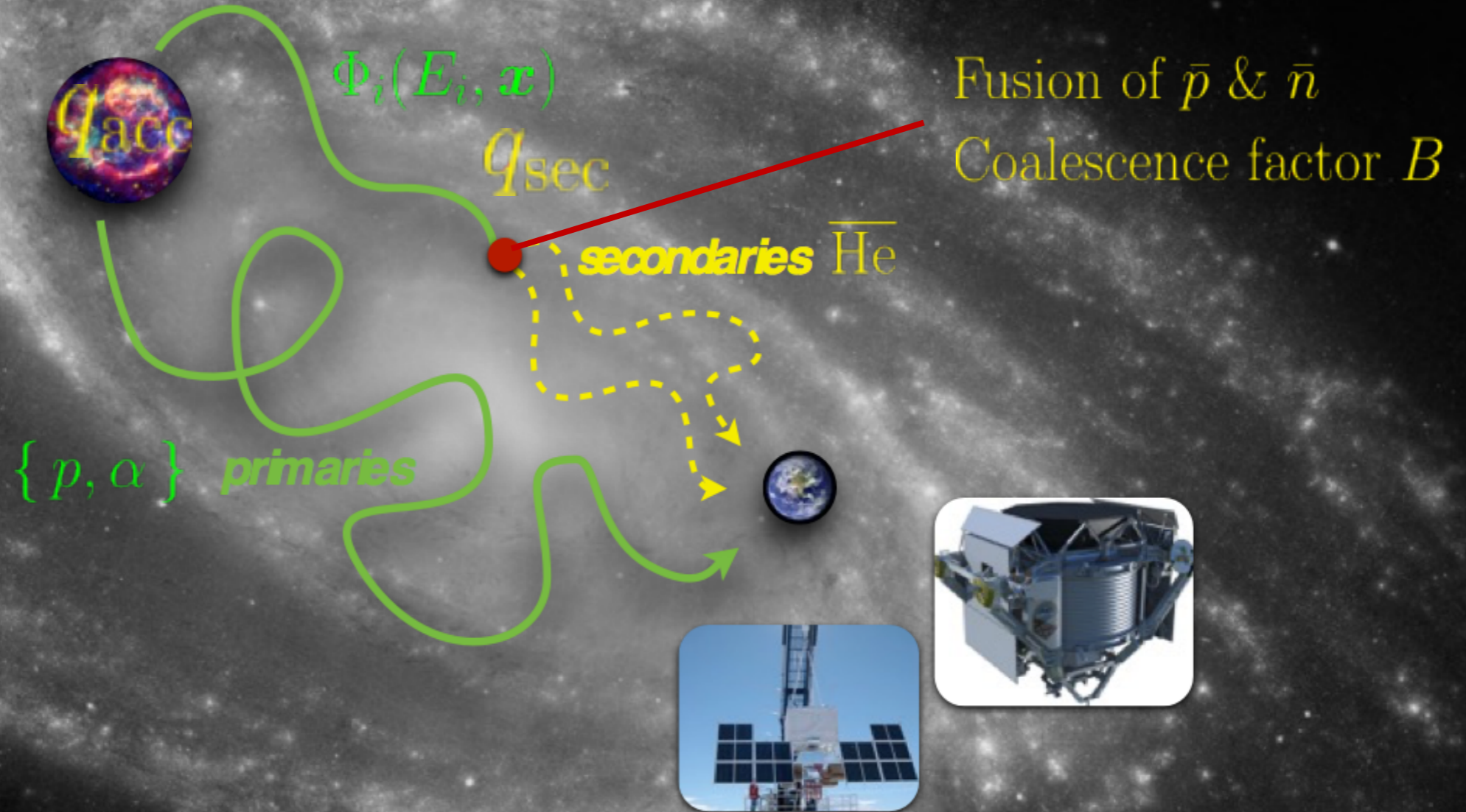
II/ Basics of baryogenesis: How to produce an anti-world?

III/ Constraining the population of anti-objects in the Galaxy / Universe

I/ Anti-helium flux from standard astrophysical processes

Secondary cosmic-ray anti-helium

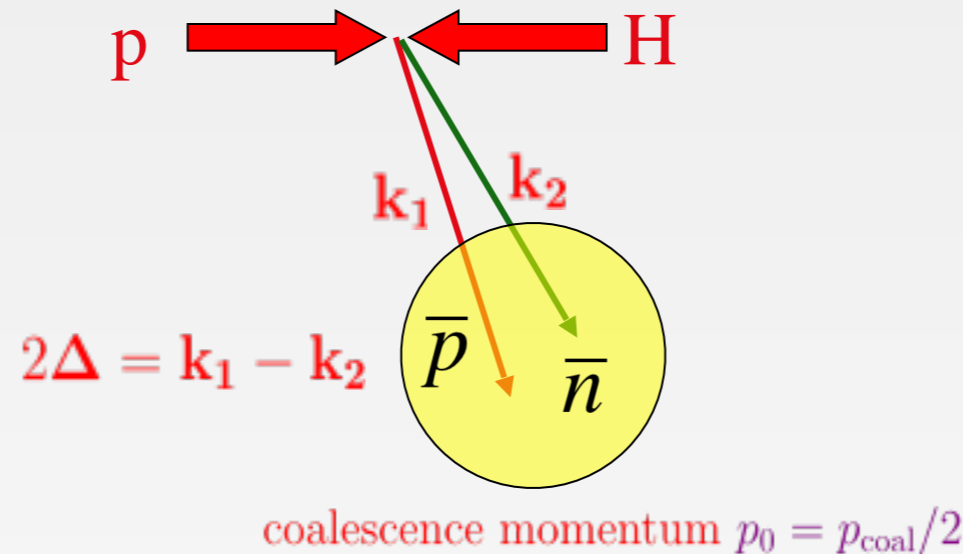
$$q_{\text{sec}}(\overline{\text{He}} | E_{\overline{\text{He}}}, \mathbf{x}) = \sum_{i \in p, \alpha} \sum_{j \in \text{H, He}} 4\pi \int dE_i \Phi_i(E_i, \mathbf{x}) n_j(\mathbf{x}) \frac{d\sigma_{ij \rightarrow \overline{\text{He}}}}{dE_{\overline{\text{He}}}}(E_i, E_{\overline{\text{He}}})$$



Solar modulation with $\phi_p^{\text{F}} \neq \phi_{\bar{p}}^{\text{F}}$

The coalescence factor

coalescence \equiv fusion of \bar{p} & \bar{n} into \bar{d} , ${}^3\overline{\text{He}}$ or ${}^4\overline{\text{He}}$



$$d^3\mathcal{N}_{\bar{d}}(\mathbf{K}) = \int d^6\mathcal{N}_{\bar{p},\bar{n}}\{\mathbf{k}_1, \mathbf{k}_2\} \times \mathcal{C}(\Delta) \times \delta^3(\mathbf{K} - \mathbf{k}_1 - \mathbf{k}_2)$$

$$B_2 = \frac{E_{\bar{d}}}{E_{\bar{p}} E_{\bar{n}}} \int d^3\Delta \mathcal{C}(\Delta) \simeq \frac{m_{\bar{d}}}{m_{\bar{p}} m_{\bar{n}}} \left\{ \frac{4}{3} \pi p_0^3 \equiv \frac{\pi}{6} p_{\text{coal}}^3 \right\}$$

Coalescence factor B_2

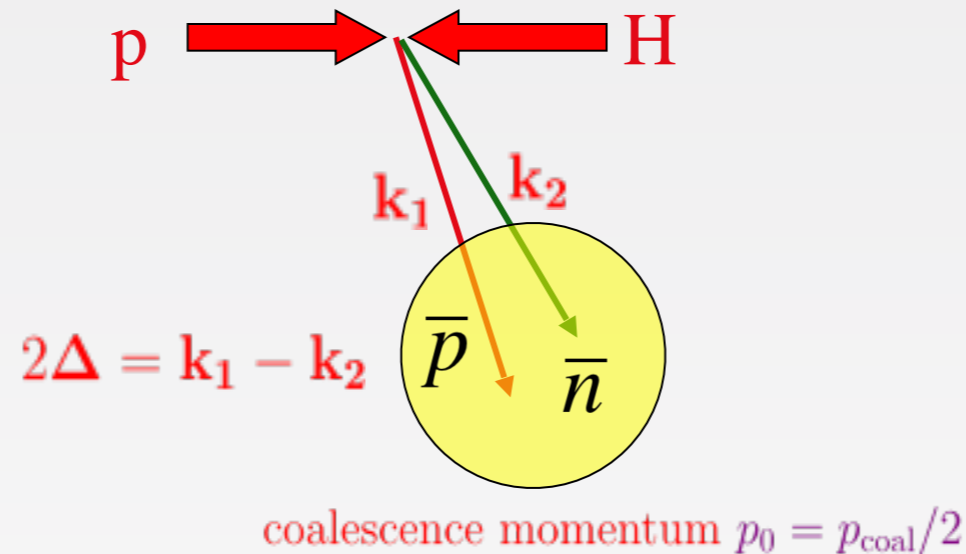
$$\frac{E_{\bar{d}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{d}}}{d^3\mathbf{K}} = B_2 \left\{ \frac{E_{\bar{p}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{p}}}{d^3\mathbf{k}_1} \right\} \left\{ \frac{E_{\bar{n}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{n}}}{d^3\mathbf{k}_2} \right\}$$

Courtesy Pierre Salati

Chardonnet, Orloff, Salati, *Phys.Lett. B*409 (1997) 313-320

The coalescence factor

coalescence \equiv fusion of \bar{p} & \bar{n} into \bar{d} , $\overline{{}^3\text{He}}$ or $\overline{{}^4\text{He}}$



Production on anti-nuclei with mass A

$$\frac{E_{\bar{A}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{A}}}{d^3\mathbf{k}_{\bar{A}}} = B_A \left\{ \frac{E_{\bar{p}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{p}}}{d^3\mathbf{k}_{\bar{p}}} \right\}^Z \left\{ \frac{E_{\bar{n}}}{\sigma_{\text{in}}} \frac{d^3\sigma_{\bar{n}}}{d^3\mathbf{k}_{\bar{n}}} \right\}^{A-Z} \quad \text{with} \quad \mathbf{k}_{\bar{p}} = \mathbf{k}_{\bar{n}} = \mathbf{k}_{\bar{A}}/A$$

Coalescence factor B_A

$$B_A = \frac{m_A}{m_p^Z m_n^{A-Z}} \left\{ \frac{\pi}{6} p_{\text{coal}}^3 \right\}^{A-1}$$

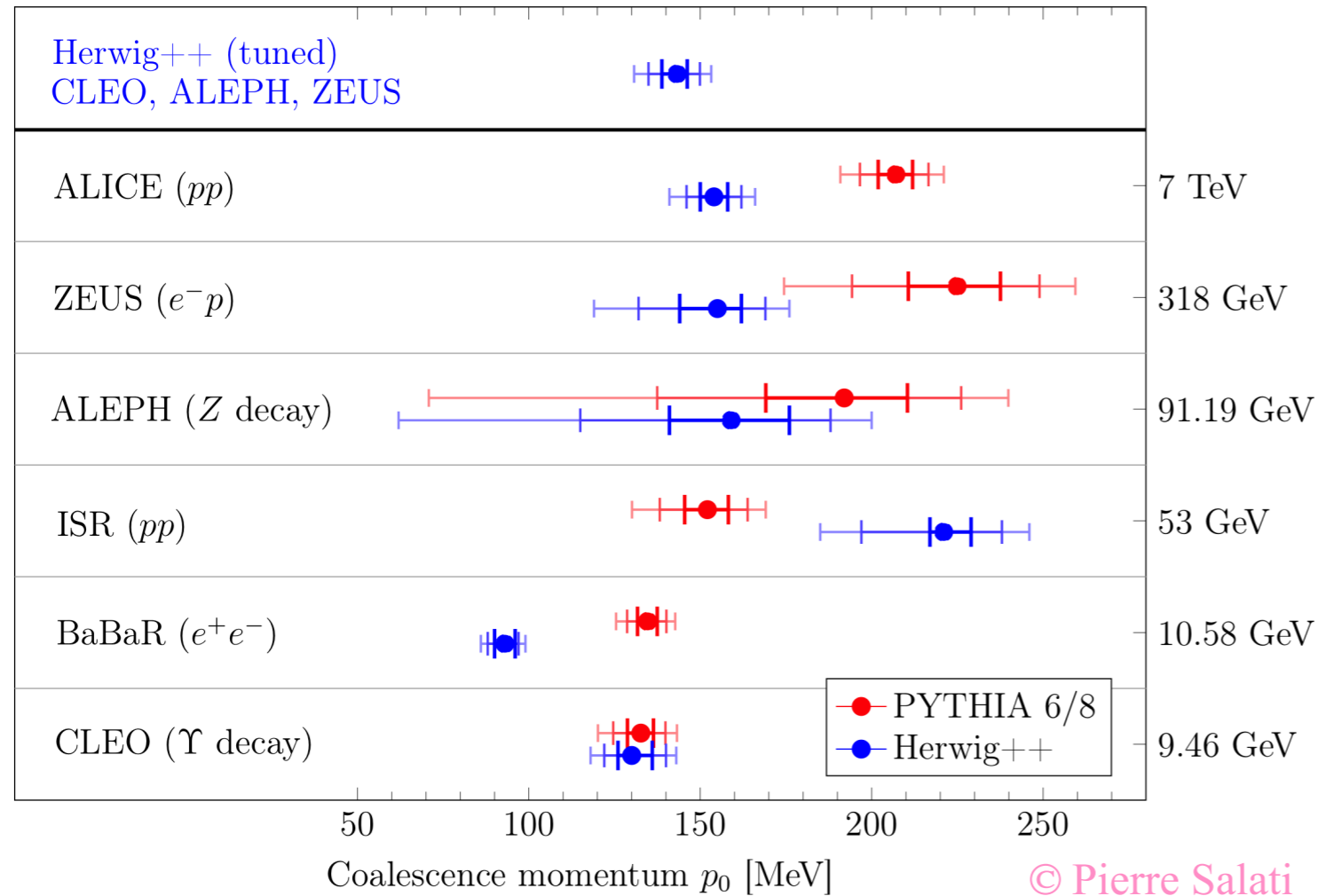
Courtesy Pierre Salati

Chardonnet, Orloff, Salati, *Phys.Lett. B409* (1997) 313-320

Determination of the coalescence momentum

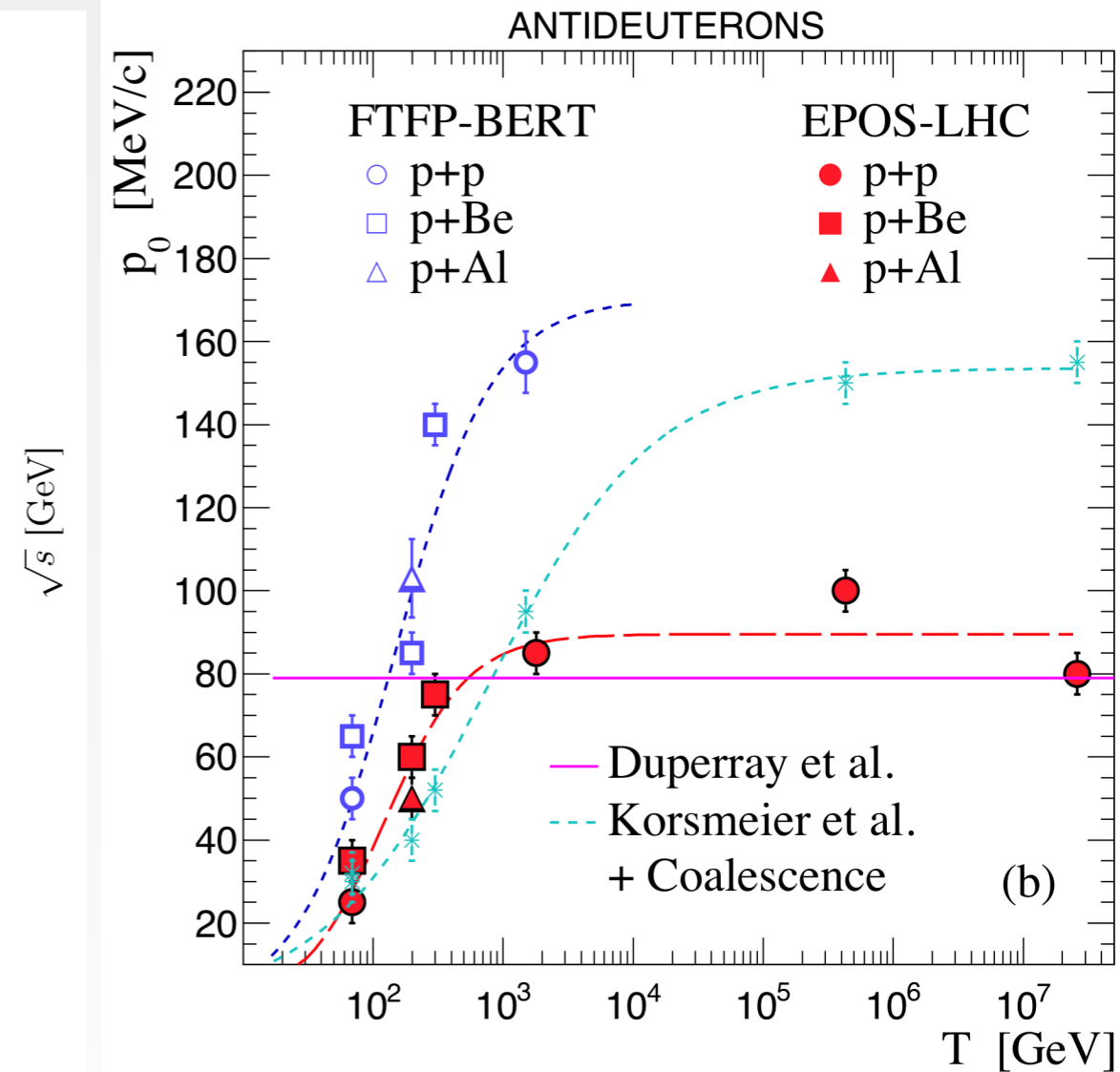
- Monte Carlo simulations show different results depending on simulator / data sets / \sqrt{s}

Fitting p_0 to data on \bar{d} production



Ibarra & Wild, JCAP 1302 (2013)

Dal & Raklev, PRD89 (2014)

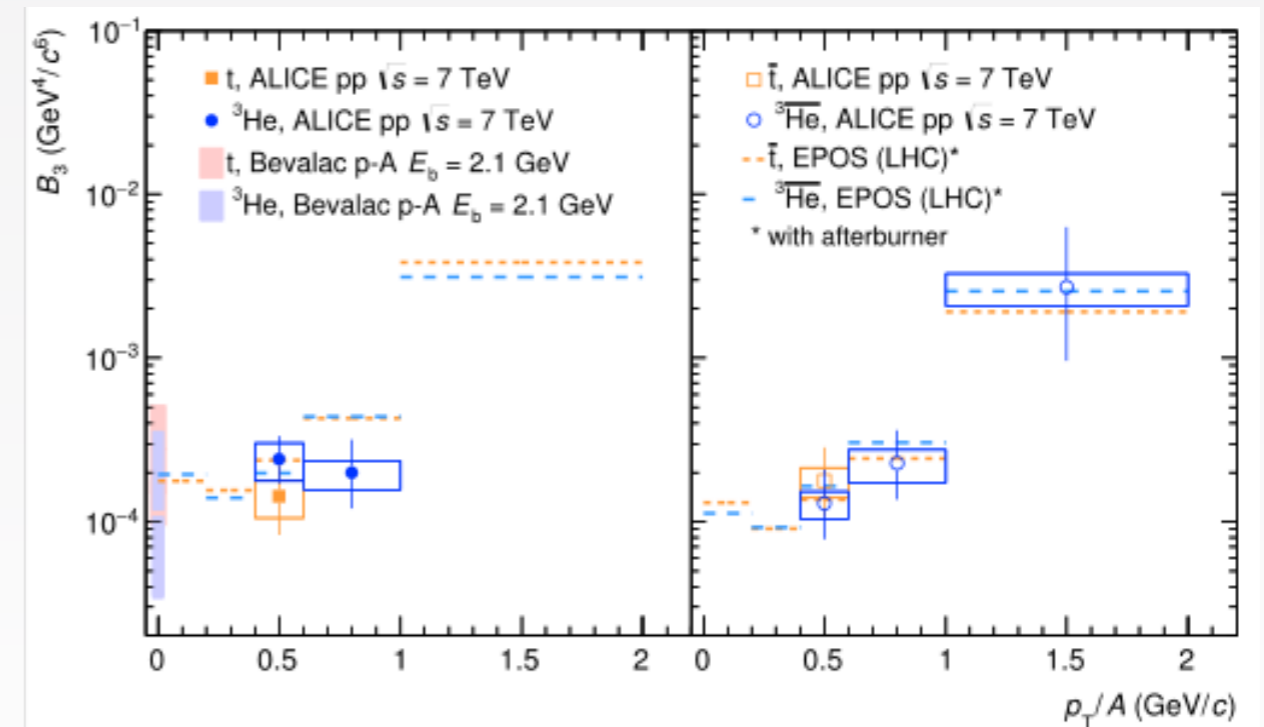
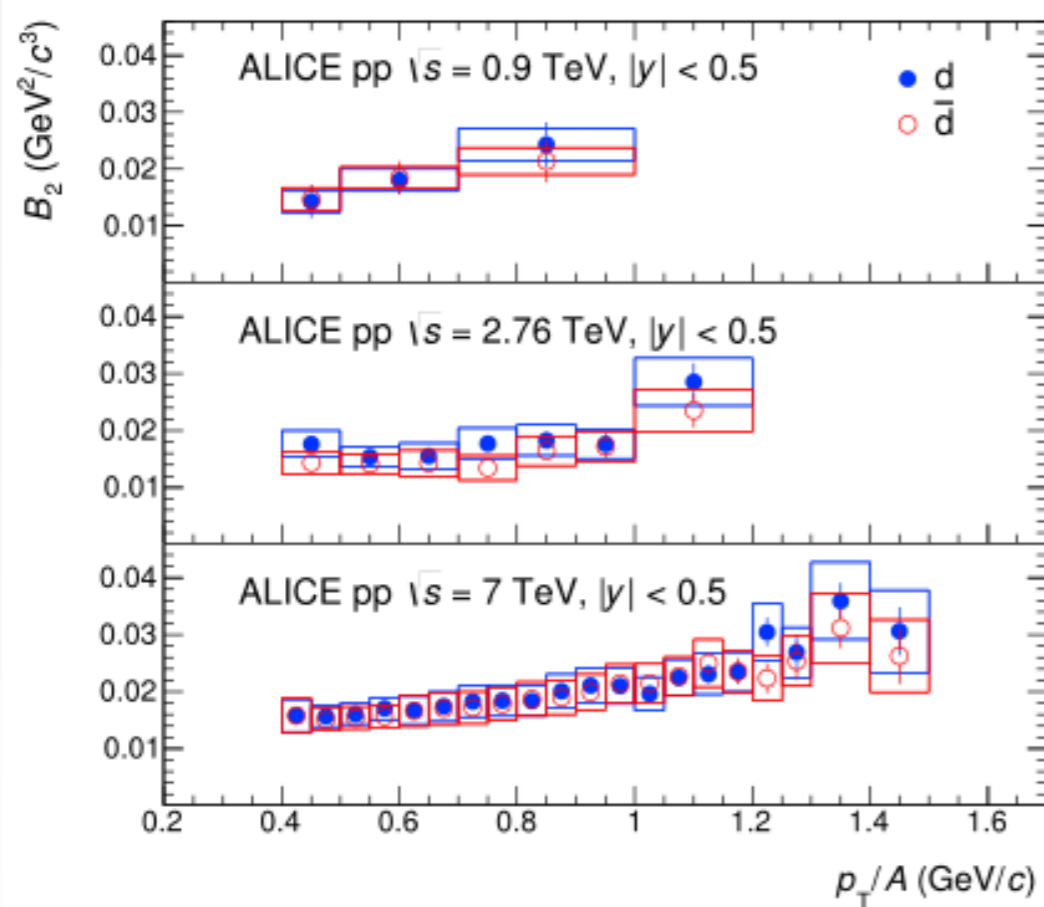
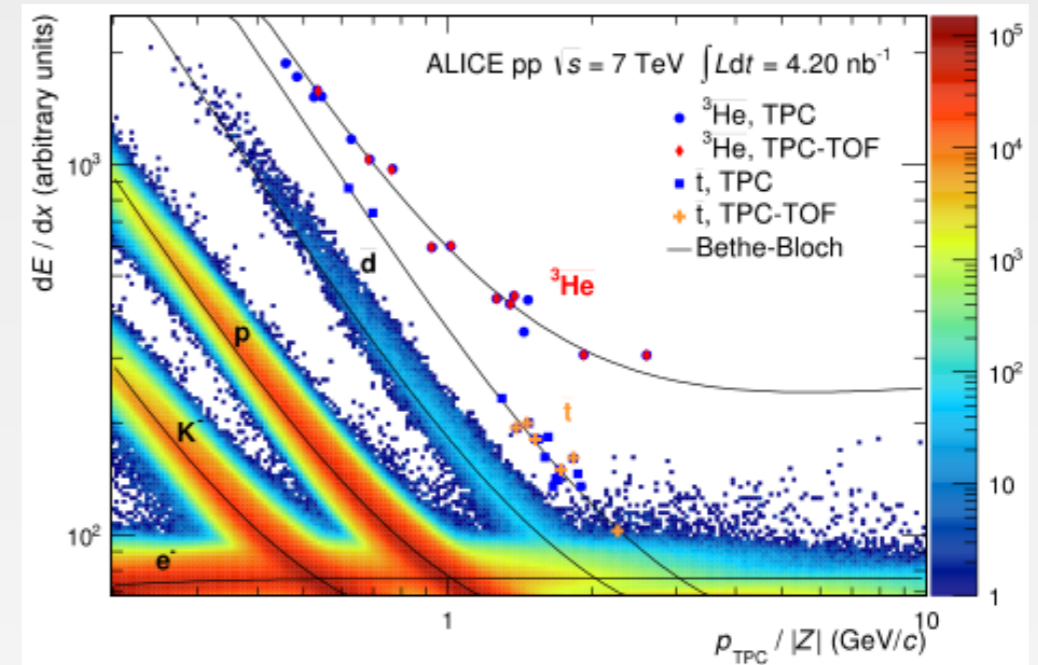


Gomez-Coral ++ PRD98 (2018)

Alice can measure the coalescence factor

- Collaborations (e.g. Alice) provide us with measurements of the coalescence factor B

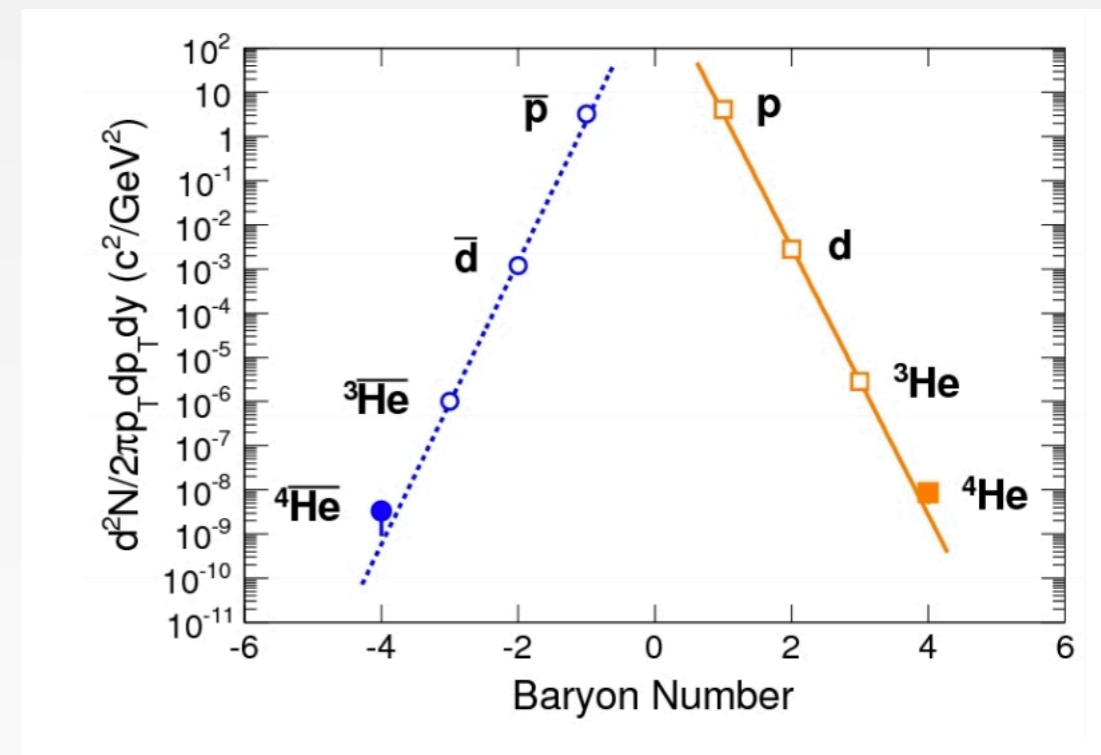
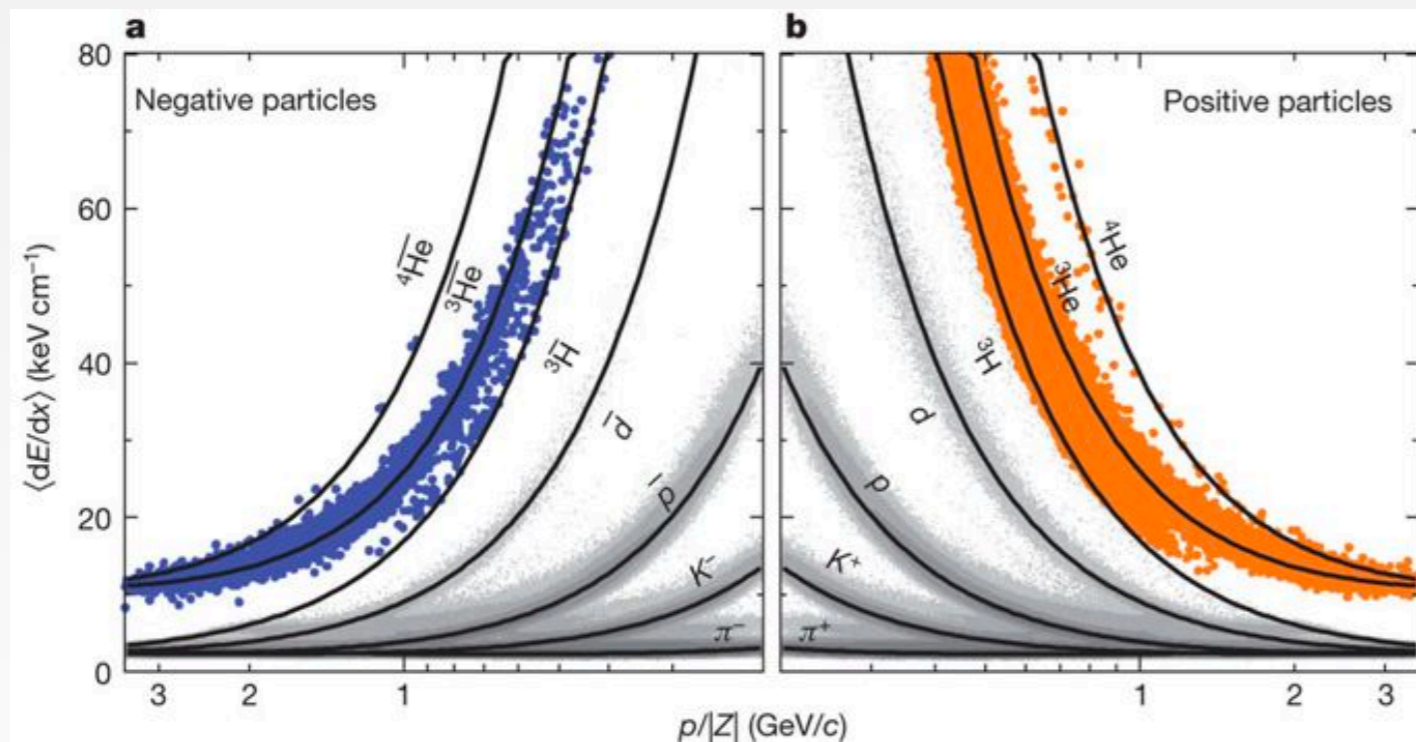
$$E_A \frac{d^3 N_A}{dp_A^3} = B_A \left(E_p \frac{d^3 N_p}{dp_p^3} \right)^Z \left(E_n \frac{d^3 N_n}{dp_n^3} \right)^N$$



Acharya++ PRDC97 (2018)

What about ${}^4\overline{\text{He}}$?

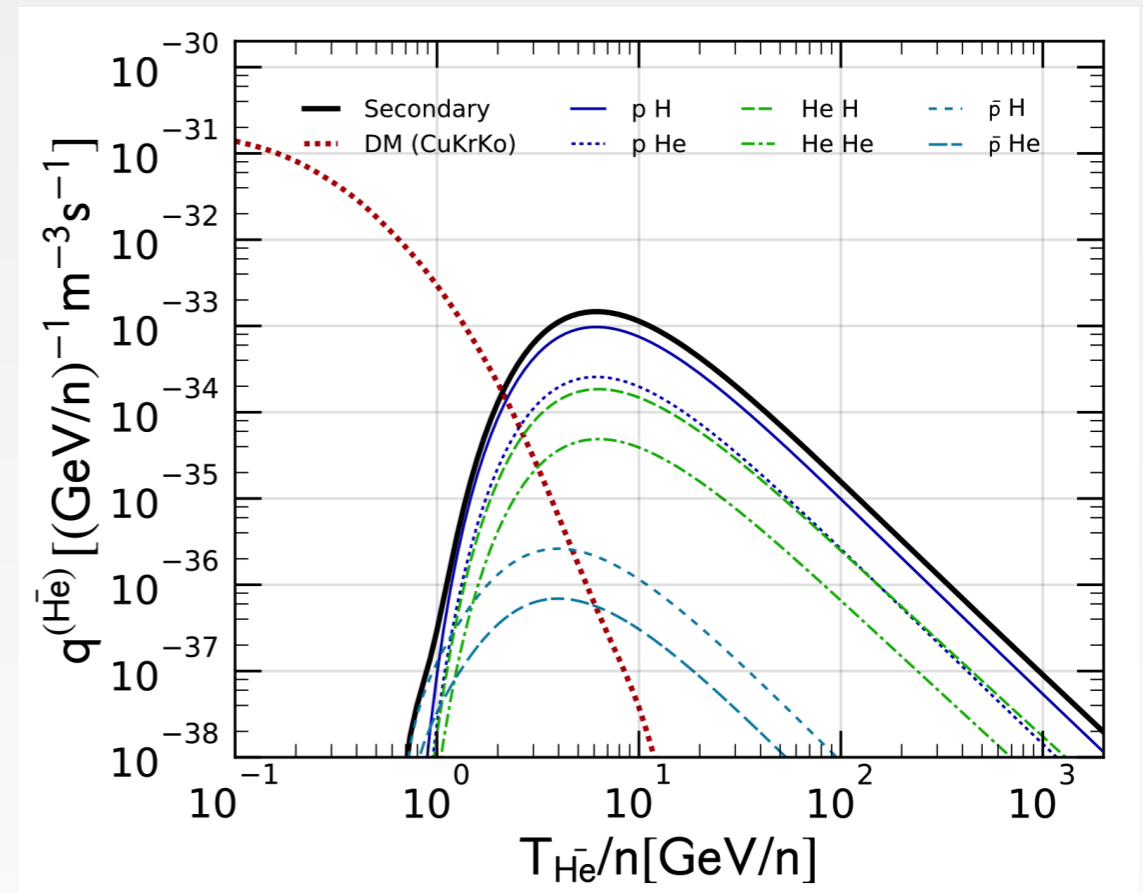
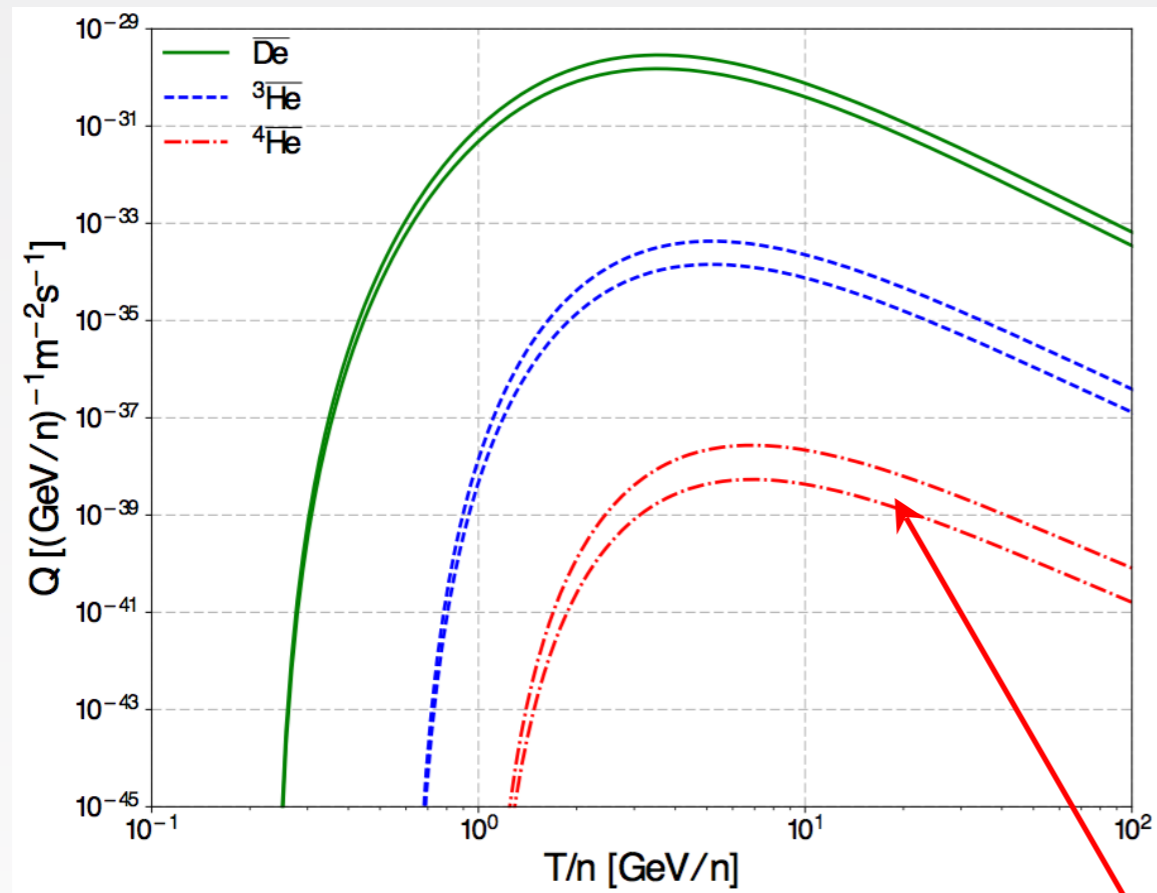
- First measurement of ${}^4\overline{\text{He}}$ by the STAR collaboration in Au-Au collision at $\sqrt{s} = 200$ GeV/n
- ${}^4\overline{\text{He}}/{}^3\overline{\text{He}} \simeq 10^{-3}$



STAR Collaboration, Nature 473 (2011)

Source term for production by spallation

$$q_{\text{sec}}(\bar{\text{He}} | E_{\bar{\text{He}}}, \mathbf{x}) = \sum_{i \in \text{p}, \alpha} \sum_{j \in \text{H}, \text{He}} 4\pi \int dE_i \Phi_i(E_i, \mathbf{x}) n_j(\mathbf{x}) \frac{d\sigma_{ij \rightarrow \bar{\text{He}}}}{dE_{\bar{\text{He}}}}(E_i, E_{\bar{\text{He}}})$$



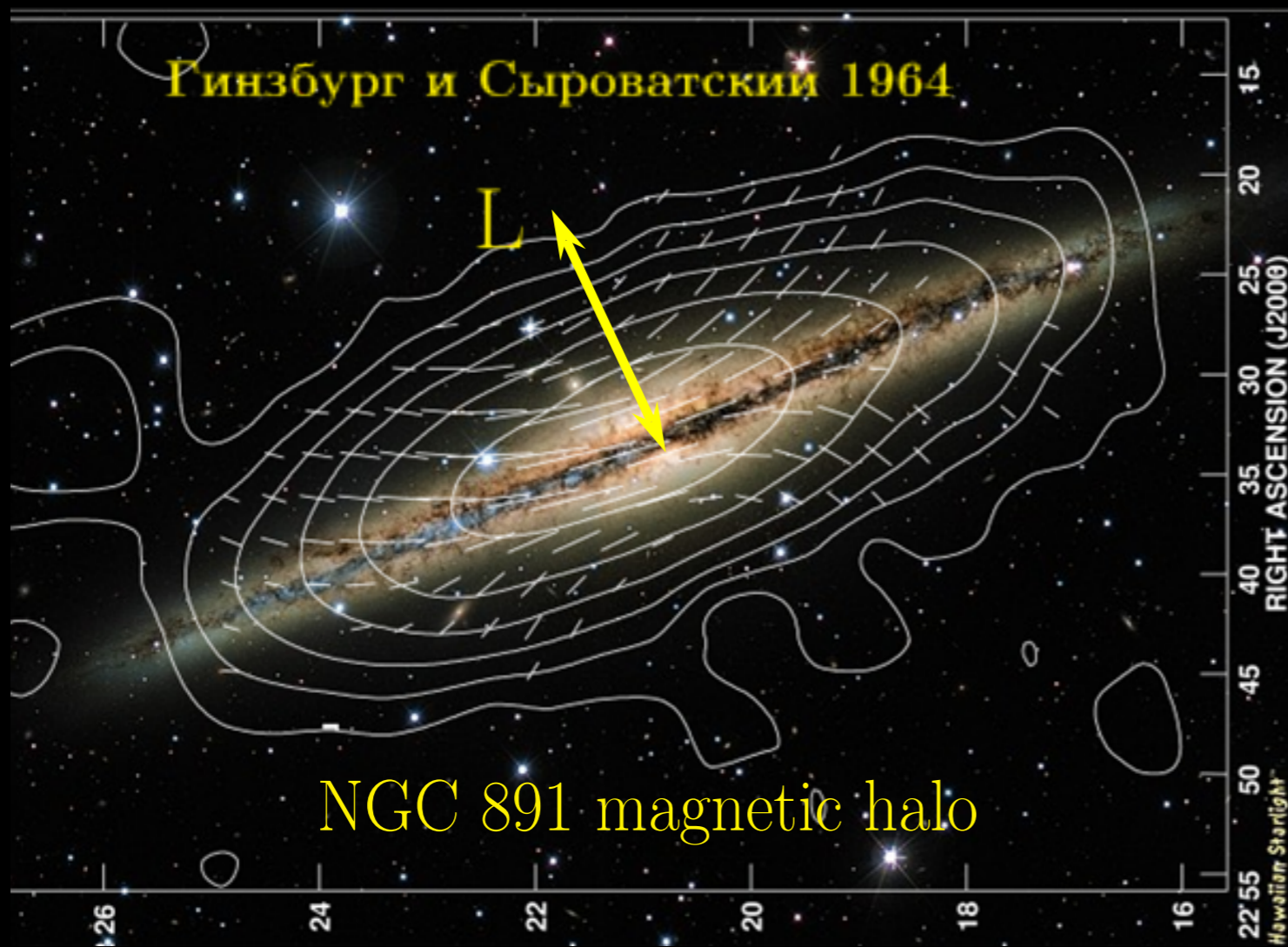
V. Poulin et al., Phys. Rev. **D99** (2019) 023016

M. Korsmeier et al., Phys. Rev. **D97** (2018) 103011

$$7.7 \times 10^{-7} \leq \frac{B_4}{\text{GeV}^6} \leq 3.9 \times 10^{-6}$$

Courtesy Pierre Salati

Cosmic-ray anti-nuclei Galactic propagation



Based on code by M. Boudaud
e.g. Genolini, Boudaud et al. 2019

$$\psi = \frac{dn}{dE} = \frac{d^4N}{d^3\mathbf{x}dE}$$

$$\Phi = \frac{1}{4\pi} v \psi$$

$$(\text{GeV/nuc})^{-1} \text{cm}^{-2} \text{s}^{-1} \text{sr}^{-1}$$

ISM spallation

$$\psi + \underbrace{\nabla \cdot \{-K \nabla \psi + \psi \mathbf{V}_C\}}_{\text{convection}} + \underbrace{\frac{\partial}{\partial E} \left\{ b \psi - D_{EE} \frac{\partial \psi}{\partial E} \right\}}_{\text{E losses}} = q - (\sigma v n_H) \psi$$

$q = q_{\text{acc}}, q_{\text{sec}}, q_{\text{DM}}$

\mathbf{x} diffusion

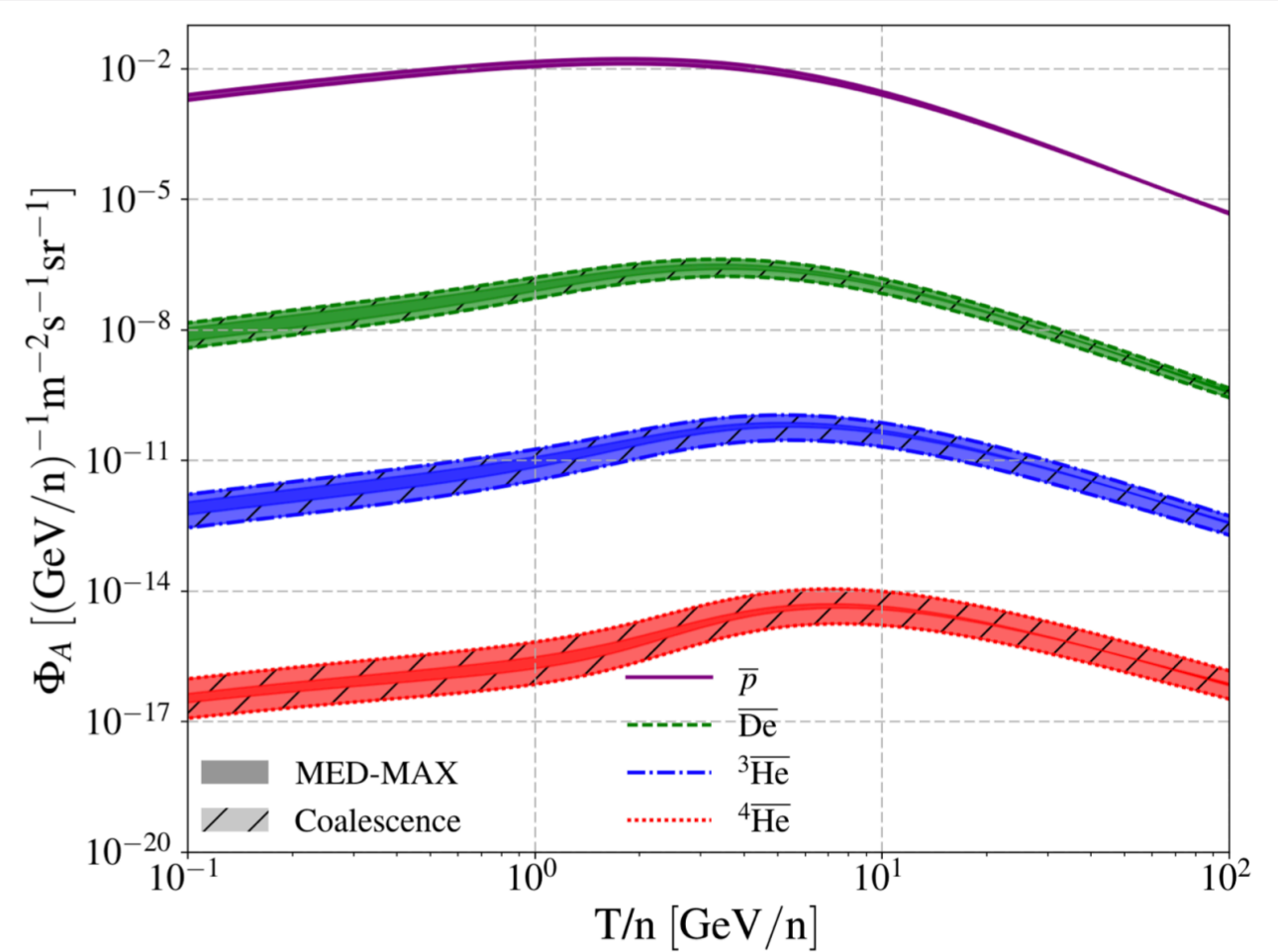
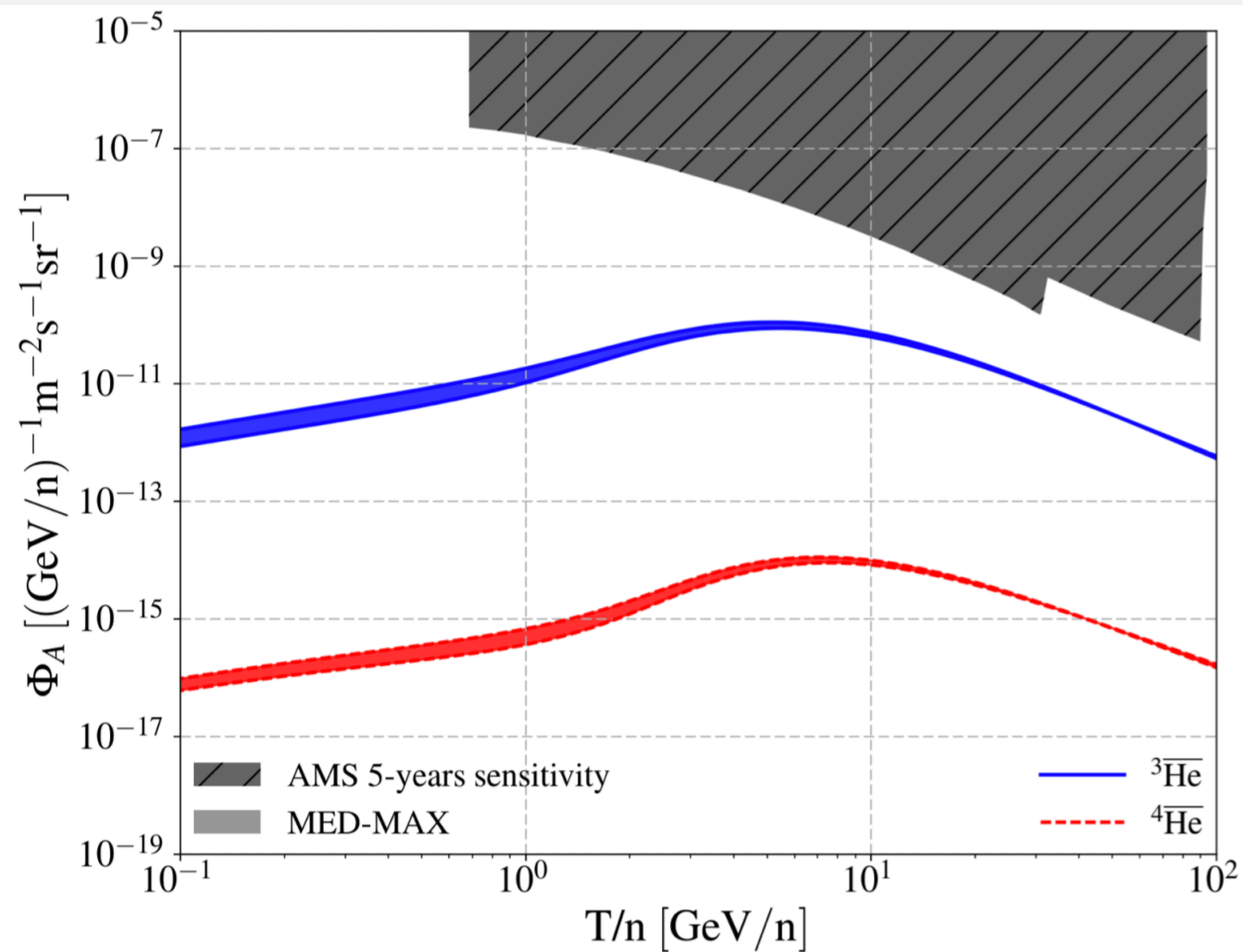
$$K = K_0 \beta \mathcal{R}^\delta \left\{ 1 + \left(\frac{\mathcal{R}}{\mathcal{R}_b} \right)^{\Delta\delta/s} \right\}^{-s}$$

E diffusion

$$D_{EE} = \frac{2}{9} \frac{V_A^2 \beta^4 E^2}{K}$$

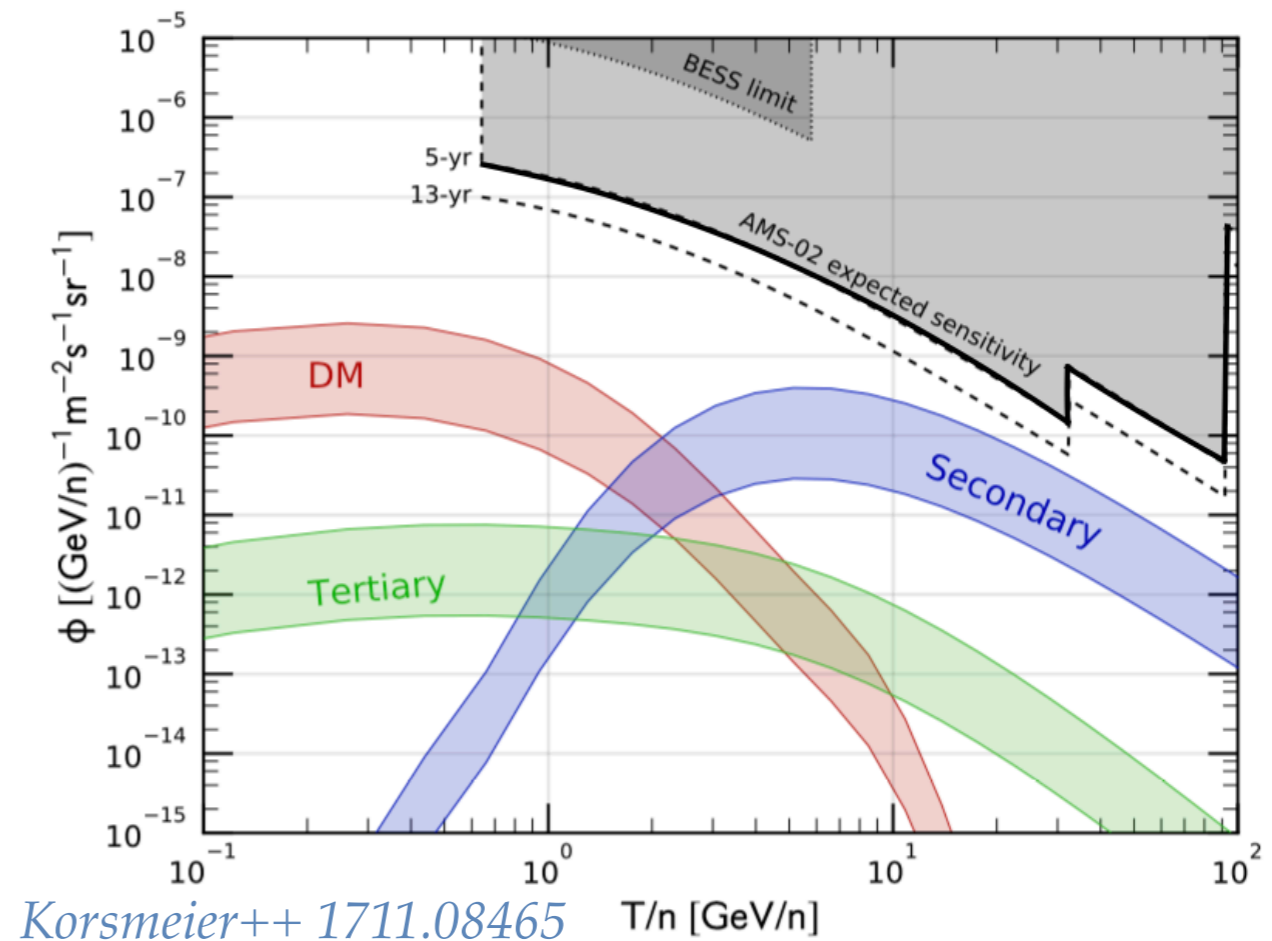
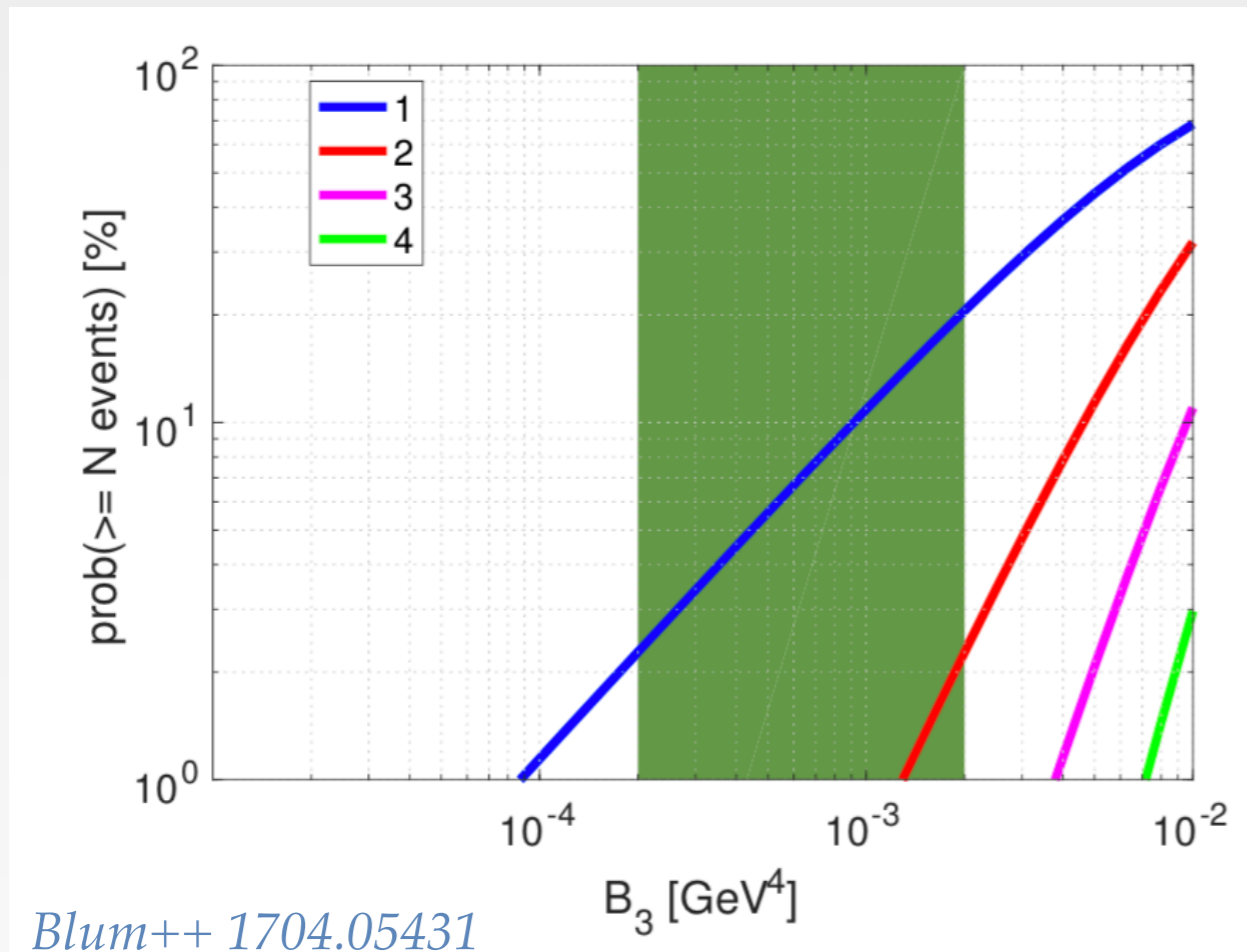
Secondaries cannot explain ${}^4\overline{\text{He}}$

- The coalescence scenario predicts a hierarchy in the flux of anti-nuclei $\phi_{A+1} \approx 10^{-3} - 10^{-4} \phi_A$
- AMS measurement is ~ 6 orders of magnitude above ${}^4\overline{\text{He}}$ “secondary” prediction
- Where is the anti-De???



VP, Salati++ PRD99 (2019)

All (recent) predictions agree!



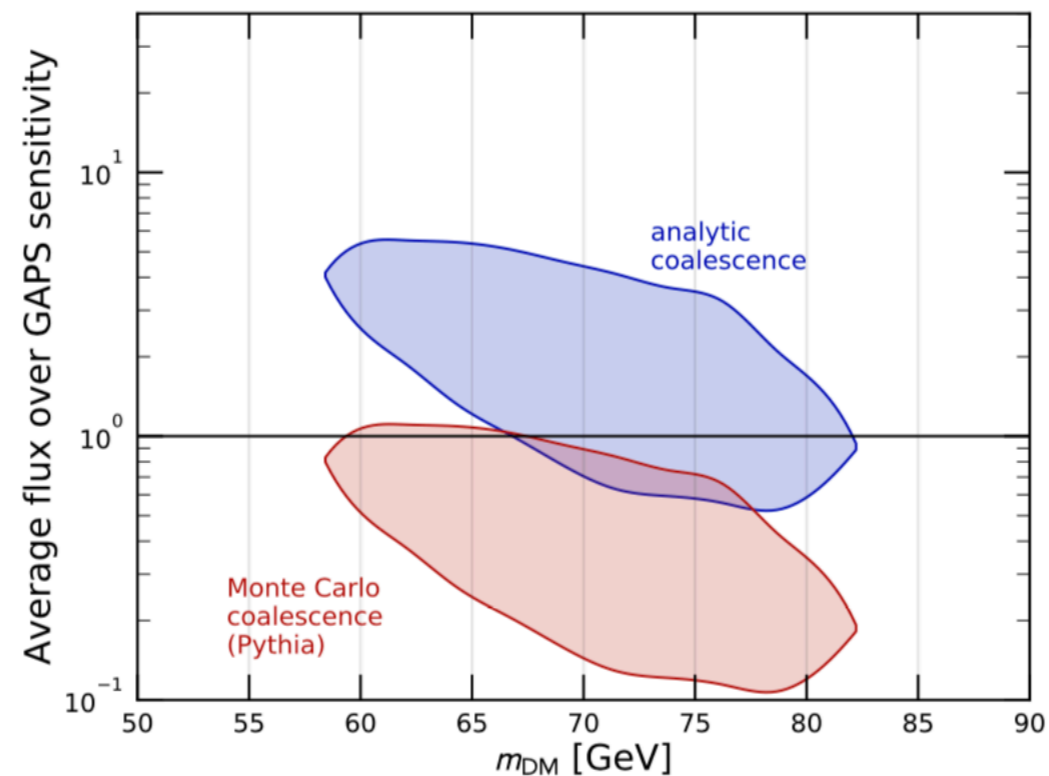
- Blum++ 2017: AMS (5yrs) could detect ~1 or 2 events **if $B_3 = 10 \times B_3$ from Alice!** AMS has detected ~6 events. probability $\rightarrow 0$.
- Korsmeier++ 2017: ~1-2 orders of magnitude **below measurement**.
- Same conclusions in Cirelli++ 2014, Herms++2016 etc...

What about Dark Matter?

- The Dark Matter explanation suffers from very similar issues! Anti-He produced via coalescence of anti-proton and anti-neutron.

$$q_{\text{DM}}(E_{\bar{D}}, \vec{x}) = \frac{1}{2} \left(\frac{\rho(\vec{x})}{m_{\text{DM}}} \right)^2 \langle \sigma v \rangle_{b\bar{b}} \frac{dN_{\bar{D}}^{b\bar{b}}}{dE_{\bar{D}}}$$

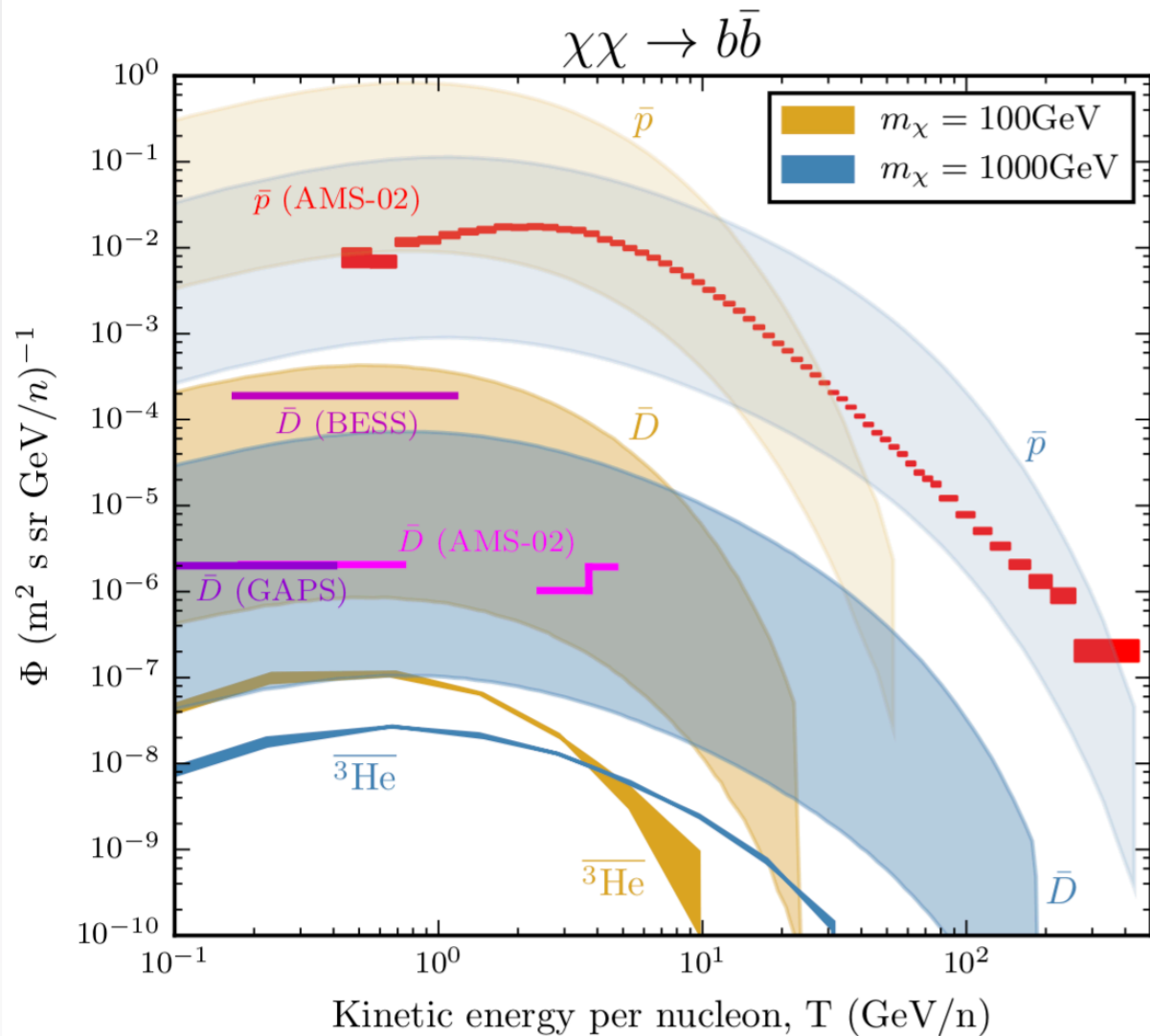
- Coalescence factor can change: very different kinematic + non-nuclear material. It leads to typically **smaller values of B_A** .



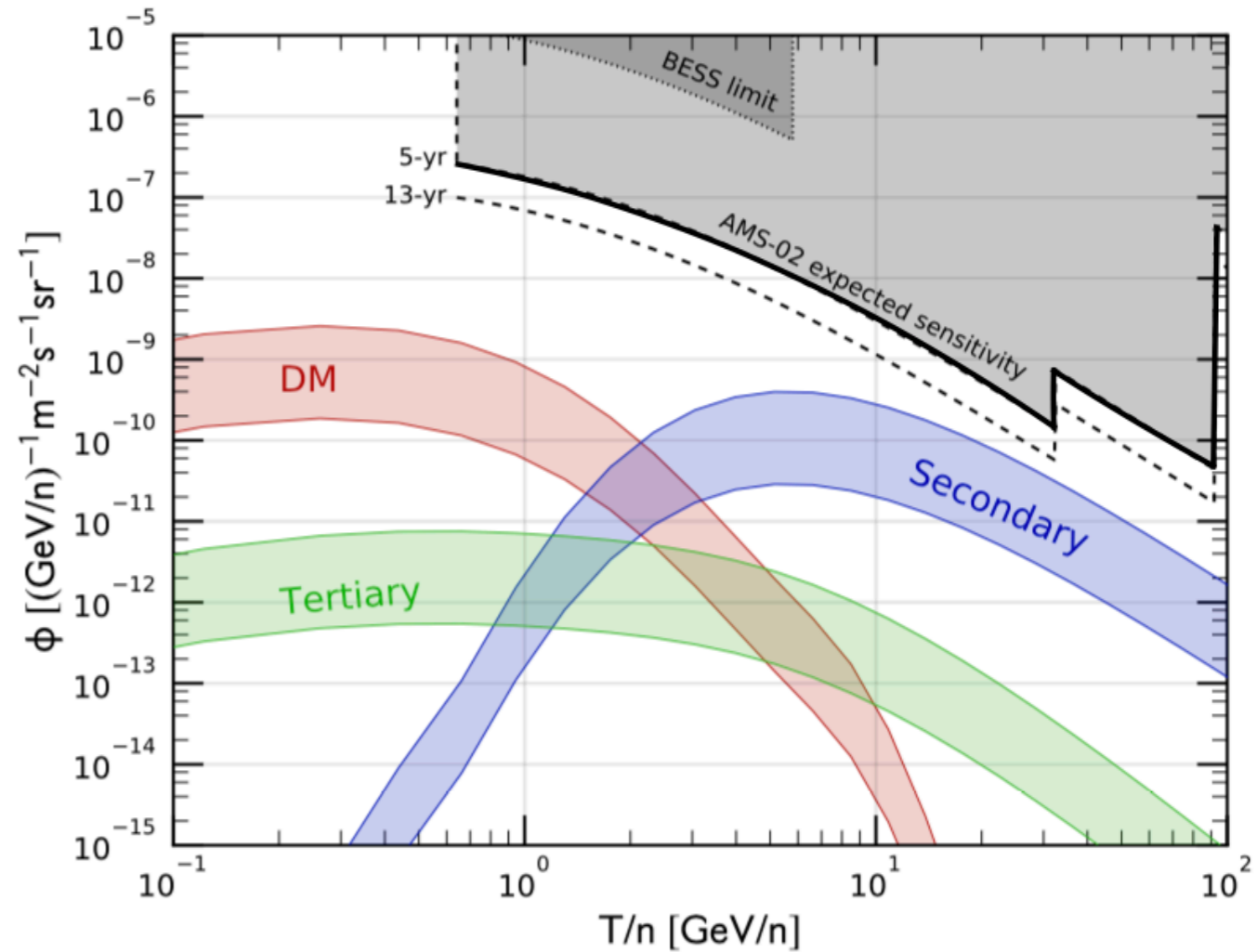
(a) Coalescence model

Korsmeier++ 1711.08465

Dark Matter is at odds with AMS-02 events



Coogan&Profumo, PRD96 (2017)



Korsmeier++ 1711.08465

- The Dark Matter flux **peaks at low kinetic energy** compared to the background.
- AMS should see associated $\bar{D}e$ and \bar{p} : Most of the **parameter space is ruled out** by \bar{p} .
- Dark Matter models cannot produce ${}^4\bar{\text{He}}$ via coalescence.

II/ Basics of baryogenesis: how to produce an anti-world

Three types of cosmological baryon asymmetry

e.g. Bambi&Dolgov 2007

- η is homogeneous, the universe is 100% matter dominated;
- average η is 0 but there are very large domains of matter and anti-matter;
- η is not spatially constant: there are lumps of antimatter in a matter dominated universe.

AMS-02 can typically probe scenario iii)

Sakharov Conditions



See Kolb&Turner for pedagogic discussion

A successful baryogenesis requires

Sakharov 1967

- **Baryon number violation:** if $B = 0$ at an initial time, and there are not B -violating processes, non-zero B cannot be generated.

nb: inflation makes sure that $B = 0$ initially.

- **P and CP violation:** ensure that opposite B -violating processes do not take place with an equal rate, resulting in no net Baryon asymmetry.

$$\Gamma(i \rightarrow f; p_i, s_i; p_f, s_f) \neq \Gamma(f \rightarrow i; -p_f, -s_f; -p_i, -s_i).$$

- **Departure from equilibrium (or CPT violation):** at equilibrium distribution are Fermi-Dirac or Bose-Einstein at temperature T . B -violating processes leads to $\mu = 0$. Because CPT ensure that $m_B = m_{\bar{B}}$, there cannot be any net Baryon asymmetry.

$$f_{\pm} = \frac{1}{e^{(E-\mu)/T} \pm 1}$$

$$n(t) = \frac{g}{(2\pi)^3} \int d^3p f_s(t, p)$$

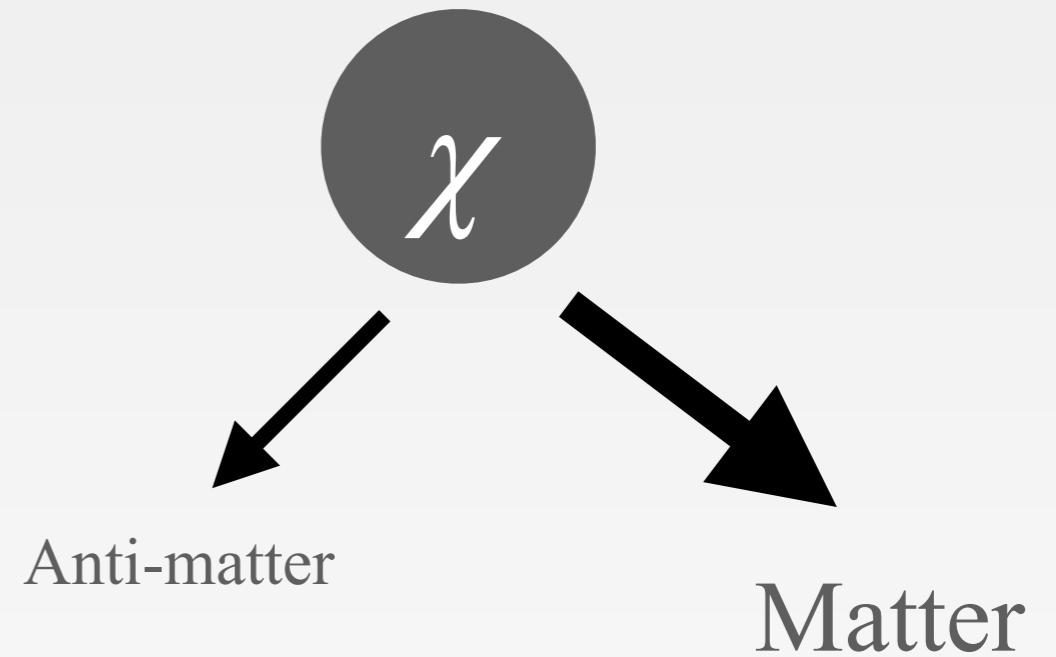
Two types of scenarios

- Cosmic phase transition



Electroweak baryogenesis,...

- Decay of new particles



GUT baryogenesis, leptogenesis...

Standard scenario predicts homogeneous baryon number, no anti-objects

Baryogenesis in the standard model

- **Maximal P violation:** SM is a chiral theory! Right-handed fermions (Left-handed anti-fermions) are gauge singlets w/r to $SU(2)_L$.
- **CP violation:**
 - first detected in observing neutral K meson decay, now also established in B meson decays.
 - *All* CP-breaking effects in the quark sector can be understood in terms of the phase δ which appears in the CKM matrix.
 - What about the neutrino sector?
- **B-violation:** B - and L -number global symmetry are *accidental* in the SM. Only violated via non-perturbative effects: ‘instantons’ and ‘sphalerons’. Rate is negligible today huge before EWPT.

Christensen et al, 1964

$$\emptyset \leftrightarrow 2u + d + 2c + s + t + 2b + e^- + \nu_\mu + \tau^-$$

't Hooft 1976

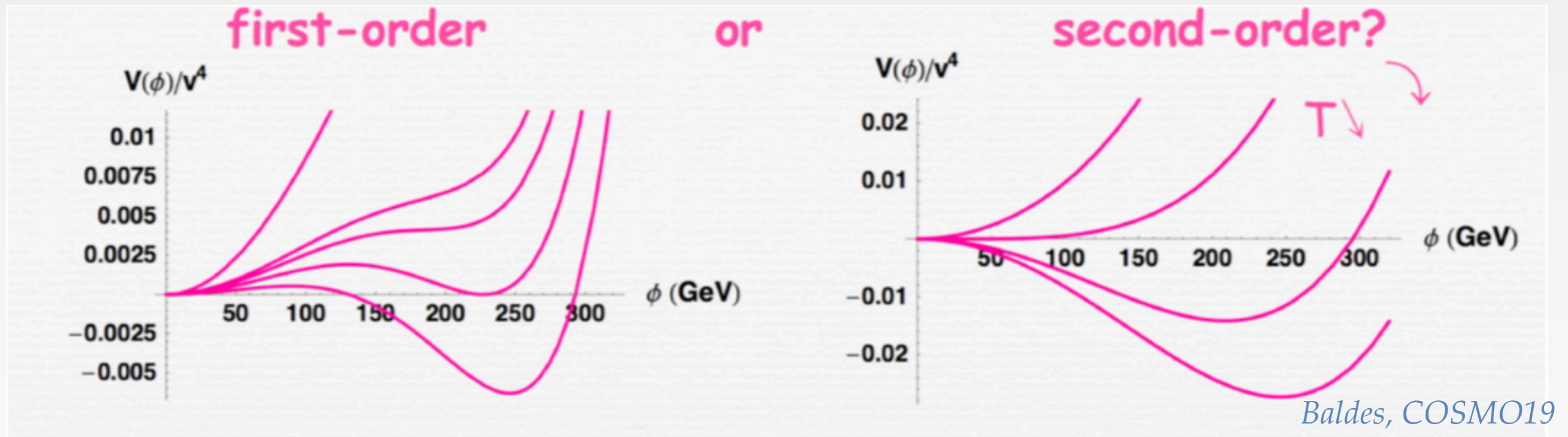
nb: $B + L$ violated, but $B - L$ still a symmetry

- **Out-of-equilibrium:** Provided by the expansion of the universe, when $H > \Gamma$.

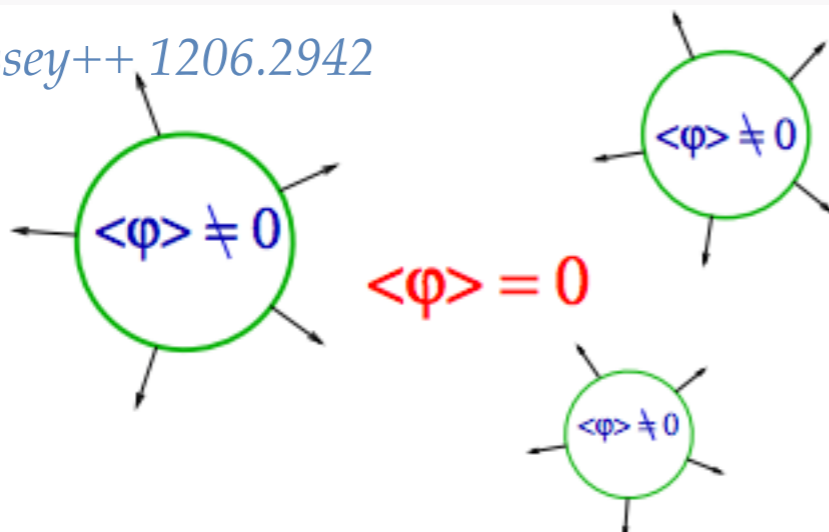
Electro-Weak baryogenesis

- Occurs during the EW phase-transition, i.e., when $SU(2)_L \times U(1)_Y$ is broken to $U(1)_{em}$ at $T \sim 100$ GeV.

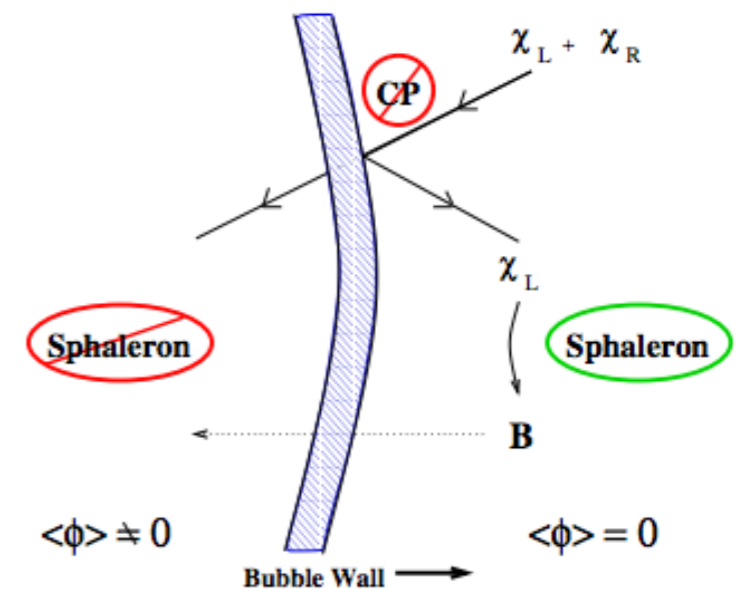
$$V_{eff}(\phi, T) \simeq D(T^2 - T_0^2)\phi^2 - ET\phi^3 + \frac{\bar{\lambda}}{4}\phi^4$$



Morrissey++ 1206.2942



v. POUIN - LUPM (CNRS)



Problems with SM EW Baryogenesis

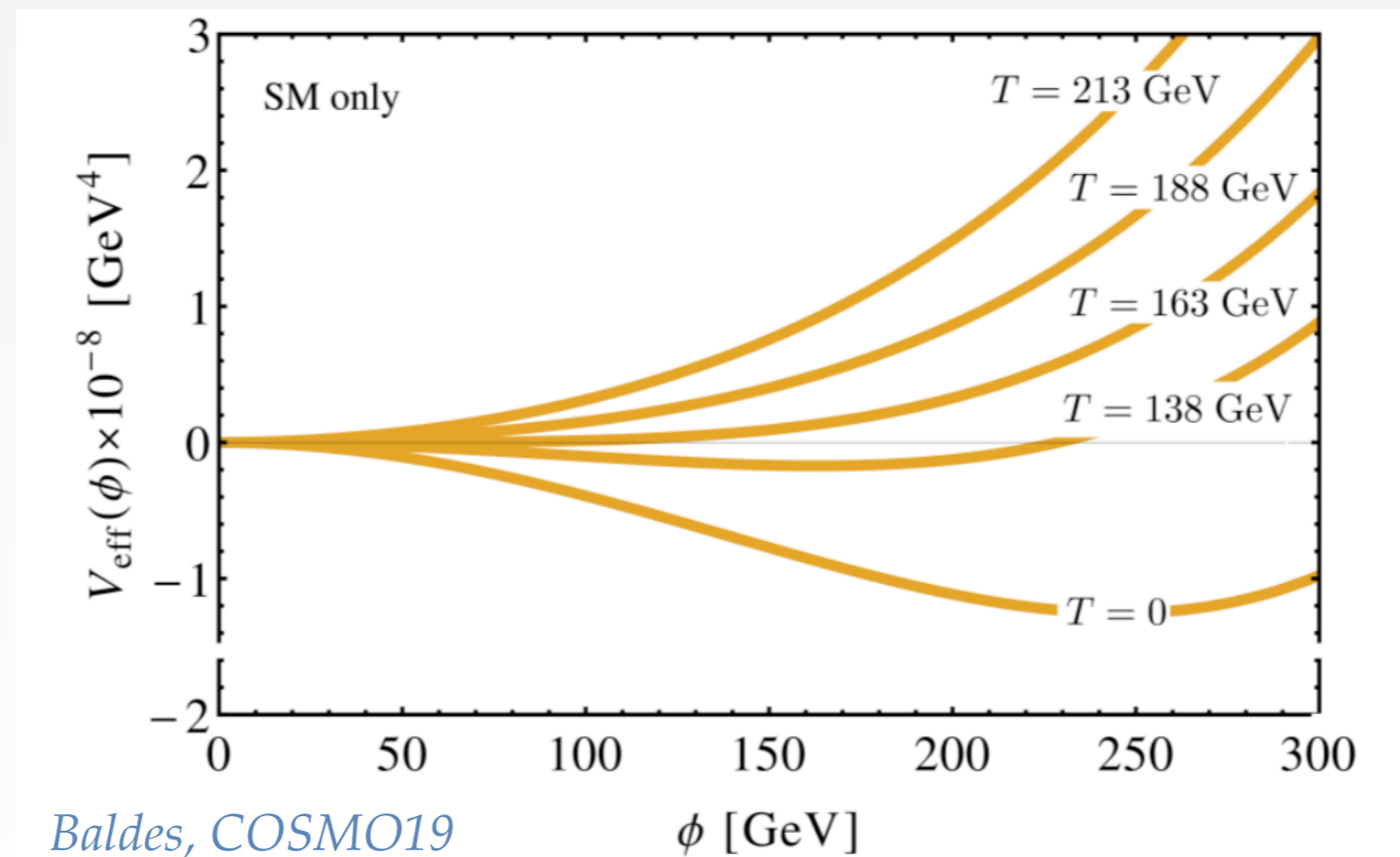
- Too little CP violation: $\epsilon_{\text{CP}} \sim 10^{-20} \ll \eta$

$$\epsilon_{\text{CP}} \sim J \frac{(m_t^2 - m_c^2)(m_t^2 - m_u^2)(m_c^2 - m_u^2)(m_b^2 - m_s^2)(m_b^2 - m_d^2)(m_s^2 - m_d^2)}{E^{12}},$$

$$J = -\text{Im} [V_{us} V_{cd} V_{cs}^* V_{ud}^*] = c_{12} c_{23} c_{13}^2 s_{12} s_{23} s_{13} s_{\delta} \sim 3 \times 10^{-5}.$$

$$E_{\text{EW}} \sim v \sim 10^2 \text{ GeV}$$

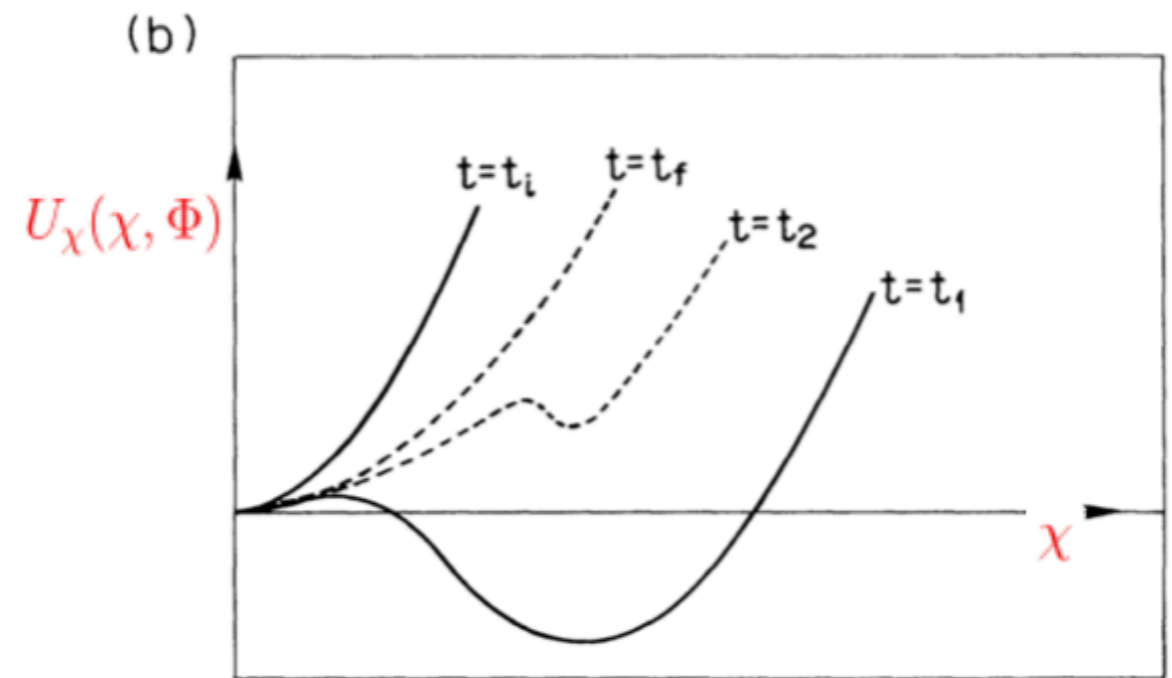
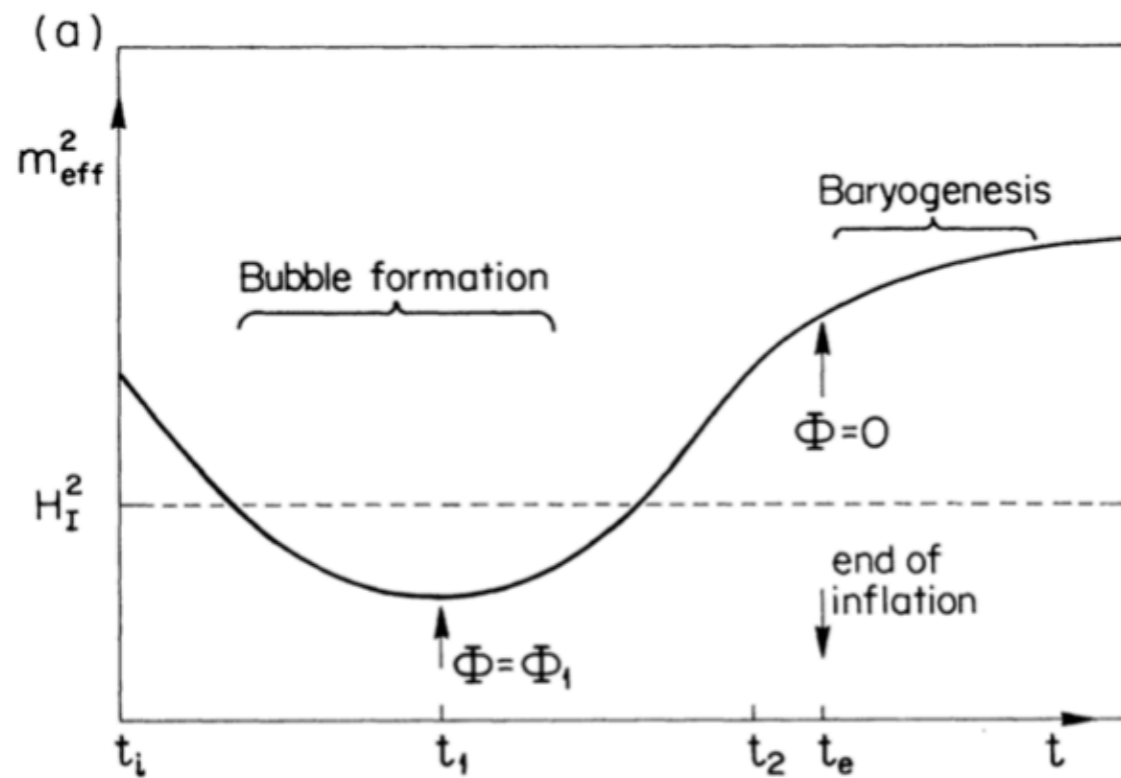
- Too strong Sphaleron rate in the broken phase: EWPT is not strongly first order (crossover), Higgs is too heavy.



Inhomogeneous baryon number

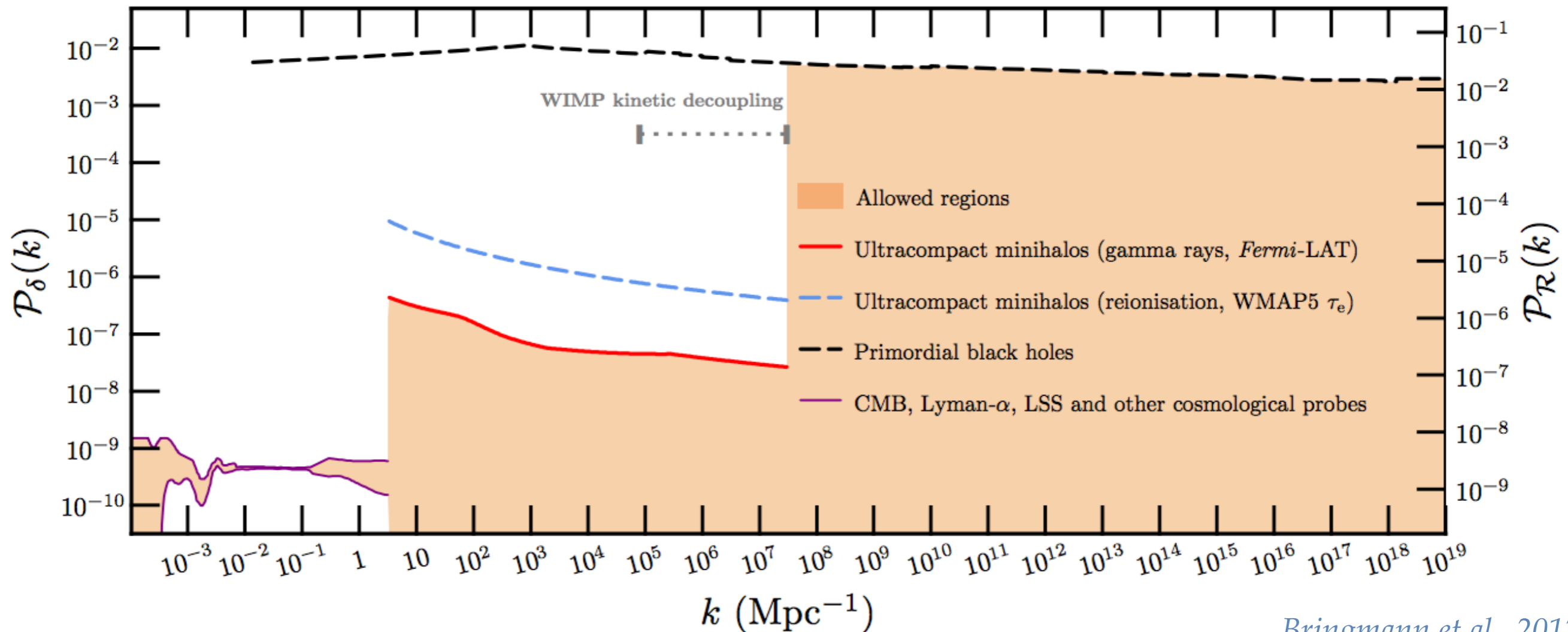
Dolgov & Silk 1994

$$U_\chi(\chi, \Phi) = \lambda_1(\Phi - \Phi_1)^2|\chi|^2 + \lambda_2|\chi|^4 \ln \frac{|\chi|^2}{\sigma^2} + m_0^2|\chi|^2 + m_1^2\chi^2 + m_1^{*2}\chi^{*2}.$$



- Modified “Affleck-Dine” baryogenesis: a complex scalar field carrying a non-zero baryon number coupled to the inflaton.
- Time-dependent effective mass allows for formation of ‘bubbles’ with high baryon number in some regions of space + overall homogeneous baryon number.
- After inflation, at horizon re-entry bubble collapse into compact objects (black holes, stars...)

We know almost nothing about small scales

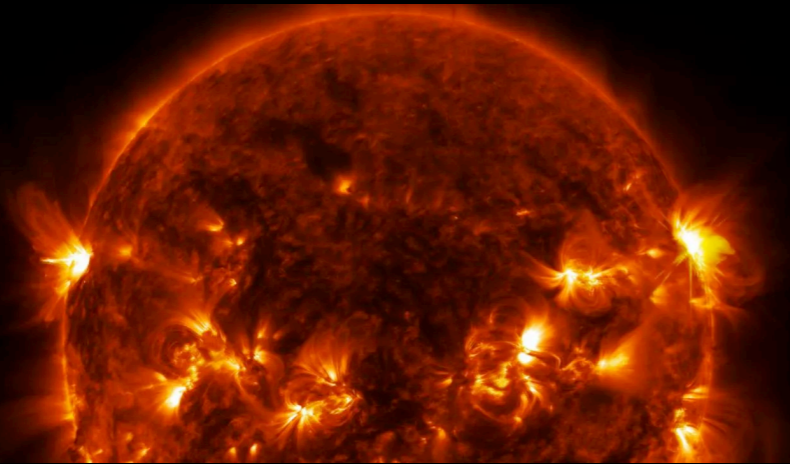


Bringmann et al., 2012

- Could there be pocket of antimatter at small scales?
- AMS probes antimatter at sub-galactic scales (i.e, $k > 10^2 \text{ Mpc}^{-1}$).

III/ Constraining the population of anti-objects

What can we learn from current data?



- AMS-02 might have identified $6 \text{ } ^3\overline{\text{He}}$ and $2 \text{ } ^4\overline{\text{He}}$. The event rate is $\sim 1\overline{\text{He}}$ for 10^8He .
- Questions: i) what population do we need to explain the measurements?
ii) Can such objects survive over cosmological timescale?
iii) How can such objects accelerate CRs?

Clouds of anti-matter in our Galaxy?

- How many of them? What are their densities? What volume would they occupy?
- AMS-02 measurements can help us answer these questions.

Assumption: acceleration and propagation of Cosmic Rays are identical for matter and anti-matter.

$$\frac{\phi_{\overline{\text{He}}}}{\phi_{\text{He}}} \simeq \frac{N_{\overline{\text{He}}}}{N_{\text{He}}} = \left(\frac{n_{\overline{\text{He}}} V_{\overline{\text{He}}}}{n_{\text{He}} V_{\text{He}}} \right) \Rightarrow n_{\overline{\text{He}}} V_{\overline{\text{He}}} \simeq 10^{-8} (n_{\text{He}} V_{\text{He}})$$

• Measured by AMS-02: 10^{-8}

• what we want to learn

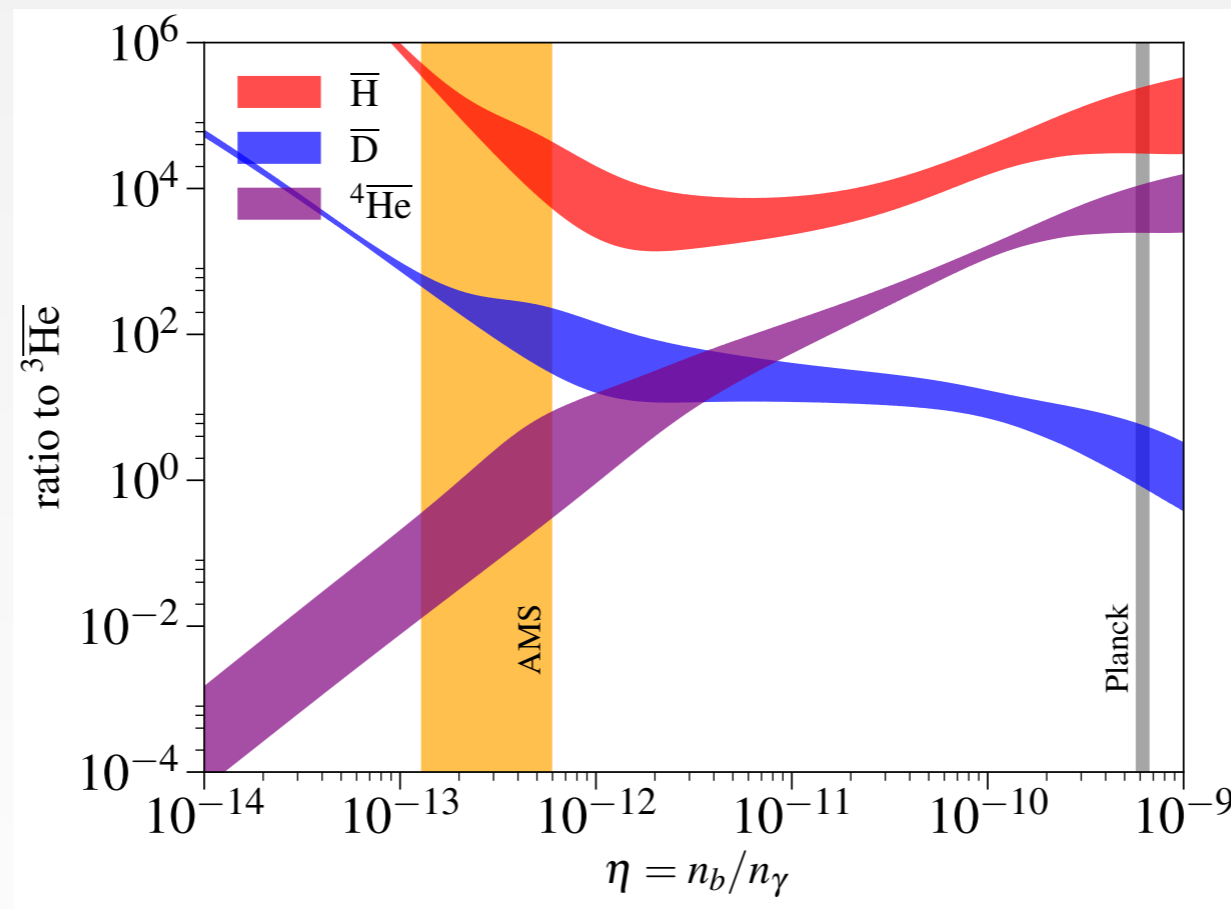
- Are there small, very dense objects or large, very dilute anti-domains?

Anisotropic BBN and the isotopic ratio

- Standard BBN predicts in the ISM: ${}^4\text{He}/{}^3\text{He} \sim 10^4$. Within CRs, spallation leads to ${}^4\text{He}/{}^3\text{He} \sim 5$.

Problem: observed isotopic ratio is 0.3.

- Solution: anisotropic BBN! if η is not homogeneous, there could be pockets dominated by antimatter with very low density.



Correct isotopic ratio if anti- $\eta = 10^{-3} \eta$

produced with AlterBBN
Arbey 1106.1363
Checked with PRIMAT
Pitrou++ 1909.12046

Some implications of the BBN calculation

- This immediately predicts density ratio: $\frac{N(^4\overline{\text{He}})}{N(^3\overline{\text{He}})} \simeq 0.3 \Rightarrow \frac{N(\overline{p})}{N(^3\overline{\text{He}})} \simeq 10^5$
- We predict $\sim 10^4$ primary anti-proton and ~ 0.1 De event.

This is potentially detectable with AMS-02!

- Moreover, we know in the ISM: $n_p = 10n_{\text{He}}$. AMS-02 therefore implies:

$$\frac{\phi_{\overline{\text{He}}}}{\phi_{\text{He}}} \simeq \frac{n_{\overline{\text{He}}} V_{\overline{\text{He}}}}{n_{\text{He}} V_{\text{He}}} \simeq 10^{-8} \Rightarrow \left(\frac{n_{\overline{p}}}{n_p}\right) \left(\frac{V_{\overline{M}}}{V_M}\right) \simeq 10^{-4}$$

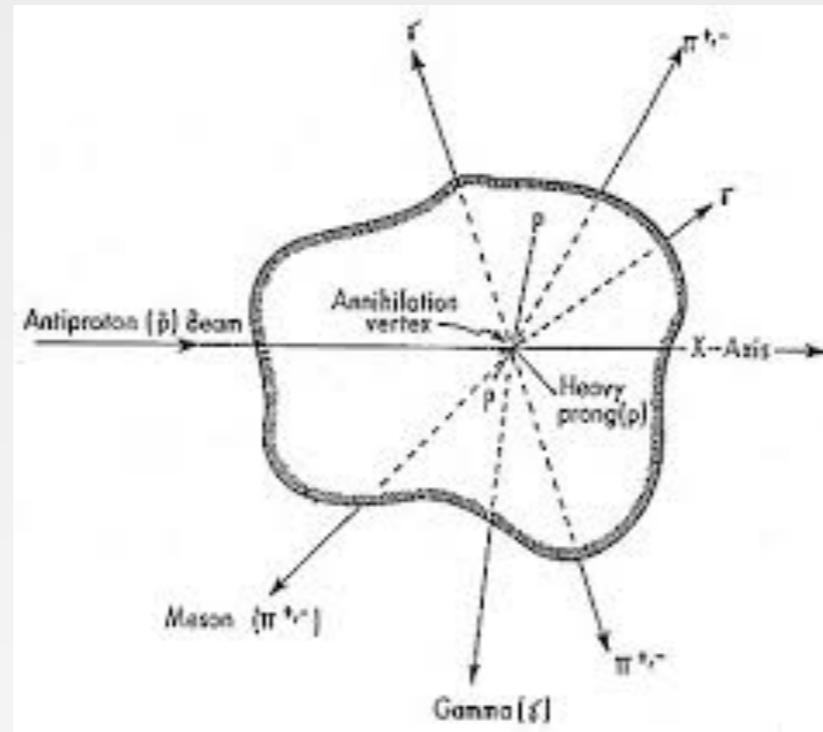
- If we assume anti-clouds are spherical with radius 1 parsec (arbitrary)

$$n_{\overline{p}} \simeq 10^5 - 10^{6.5} N_{\overline{c}}^{-1} \left(\frac{n_p}{1 \text{ cm}^{-3}}\right) \left(\frac{r_{\overline{c}}}{1 \text{ pc}}\right)^{-3} \text{ cm}^{-3}.$$

A few, very dense anti-clouds could explain AMS events!

- Question: can such objects survive in our galaxy? can we see them in γ -rays?

Anti-cloud cannot survive in our Galaxy



- \bar{p} can annihilate with p in the ISM at a rate: $\tau_{\text{ann}}^{-1} = (n_p \langle \sigma_{p\bar{p}} v \rangle)$

$$\langle \sigma_{p\bar{p}} v \rangle \simeq \begin{cases} 1.5 \times 10^{-15} \text{ cm}^3/\text{s} & T > 10^{10} \text{ K}, \\ 10^{-10} \left(\frac{\text{K}}{T}\right)^{1/2} \text{ cm}^3/\text{s} & 10^{10} \text{ K} > T > 10^4 \text{ K}, \\ 10^{-10} \text{ cm}^3/\text{s} & 10^4 \text{ K} > T. \end{cases}$$

Steigman 1976

- Our Galaxy exists since roughly $t_{\text{gal}} \simeq 2.8 \times 10^{17} \text{ s}$ and $n_p^{\text{ism}} = 1 \text{ cm}^{-3}$
- Requiring $t_{\text{ann}} > t_{\text{gal}}$ leads to $n_p^{\text{cold}} < 3.5 \times 10^{-8} \text{ cm}^{-3}$ $n_p^{\text{hot}} < 6.1 \times 10^{-5} \text{ cm}^{-3}$.

Anti-clouds cannot survive unless there is a segregation between matter and anti-matter

Survival rate in the Early Universe

- In the early universe, larger densities lead to larger annihilation rate and stronger constraints.
- The hubble time before matter-radiation equality ($z_{\text{eq}} > 3500$) is $t_H \simeq 5 \times 10^{19} (1+z)^{-2}$ s
- Before BBN ($z \sim 10^9$), annihilation happens in the relativistic regime. The constraint on the local proton density from requiring $t_{\text{ann}} > t_H$ is:

$$n_p^{\text{local}}(z_{\text{BBN}}) < 1.9 \times 10^{-8} n_p^{\text{cosmo}}(z_{\text{BBN}})$$

- Below z_{eq} , the constraint relaxes to

$$\frac{n_p^{\text{local}}}{n_p^{\text{cosmo}}}(z < z_{\text{eq}}) < \frac{6.3 \times 10^{-2}}{(1+z)^{3/2}}$$

If anti-domains were formed before BBN,
there must be less than 1 baryon per 10^8 anti-baryons within them!

γ -Ray constraints

- Annihilations lead to γ -rays that can be detected.
- There are three types of searches that can provide strong constraints:
 - i) searches for distinctive spectral features such as **a gamma-ray line**;
 - ii) searches for morphological features localized on the sky, either from **extended or point sources**;
 - iii) searches for a continuous spectrum of gamma-rays extending over **large area on the sky** (e.g. extragalactic γ -ray background).
- Type i) and iii) can provide very strong constraints on the overlap of matter/anti-matter region. Type ii) could explain some unassociated sources in the 3FGL catalog.
- Line search in FermiLAT allows to set (for a cold cloud) $n_p^{\text{local}} \lesssim 10^{-10} - 2 \times 10^{-9} \text{ cm}^{-3}$.

FermiLAT can be used to improve constraints by 2 orders of magnitude!

Anti-stars in the galaxy?

- Alternatively, anti-domains **could have formed compact objects**: naturally free of normal matter! Annihilations only occur at the surface of these objects.
- A one solar-mass would survive if formed at $z < 10^{16}$

Anti-stars cannot form from a anti-cloud because it would not survive in the early universe:
they have to be primordial!

- The Dolgov & Silk scenario could produce such objects. How many of them? What mass & composition? What is the acceleration mechanism?
- **Massive stars are short-lived** compared to t_{gal} : they would require anti-stars to form again from a cloud. This is excluded!

High-energy cosmic rays from anti-stars

- Even if such objects were created in the early universe, it is **unclear how they can lead to high-energy cosmic rays**.
- Do they lead to **supernovae explosion** that accelerate the surrounding medium? Do they experience **solar flares**? Could there be thermo-nuclear explosions from annihilations at the surface?
- Parametrically we can estimate that from a single event occurring at a given time:

$$\Phi_{\overline{\text{He}}} = \left(\frac{c}{V_{\text{gal}}} \right) \left(\frac{f_{\overline{\text{He}}} M_{\overline{*}}}{m_{\overline{\text{He}}}} \right) f_{\text{acc}} = 10^{-9} \left(\frac{(4\pi/3)(10 \text{ kpc})^3}{V_{\text{gal}}} \right) \left(\frac{M_{\overline{*}}}{M_{\odot}} \right) \left(\frac{f_{\text{acc}}}{10^{-8}} \right) \left(\frac{f_{\overline{\text{He}}}}{1} \right) \overline{\text{He}} \text{ cm}^{-2} \text{ s}^{-1}$$

If 10^{-8} of the mass of a *single* anti-helium star with $M = M_{\odot}$ is ejected in the galaxy, it can explain AMS-02 events!

- Helium would escape the galaxy in 10^8 yrs $\sim 10^{-3} t_{\text{gal}}$: there might be a population of stars!

A coherent scenario for AMS-02 anti-stars

- One possible scenario: White dwarf anti-stars were formed in the early universe in clusters.
- Binary of (long-lived) white dwarfs can lead to **type Ia supernovae!** Measurements of such events indicate a rate:

$$1.4 \times 10^{-13} \text{yr}^{-1} M_{\odot}^{-1}$$

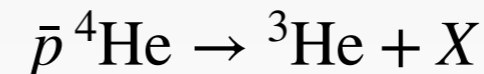
Badenes&maoz 1202.5472

- Requiring **one such event over one CR diffusion time scale** leads to a total anti-star mass of

$$\sum M_{\bar{*}} = 10^{-5} - 10^{-6} \sum M_{*}$$

- If anti-stars are heavier than 0.6Msun, producing the correct isotopic ratio requires **spallation around the anti-star.**

- We can compute the grammage required to inverse the isotopic ratio from the result of the LEAR collaboration measuring



Balestra++ 1985

- We find that it requires 20g/cm². For comparison: **this represents 1/50th of our atmosphere.**

How to see an anti-star

- Normal matter falling onto the anti-star could lead to **characteristic annihilation spectra** (line and continuum below the proton mass).
- Within 150 pc from the Sun, non-observation of such event from Bondi accretion leads to $N_{\bar{*}} < 4 \times 10^{-5} N_{*}$. *Von Ballmoos, 1401.7258*
- We can **check the 3FGL catalog for un-associated sources**: the brightest source can be used to estimate the closest distance at which an anti-star could be.
- Luminosity from annihilations to pions and subsequent decay

$$L_{\bar{*}} = 8\pi R_{*}^2 v n_p \simeq 10^{31} \left(\frac{R_{\bar{*}}}{10^{11} \text{ cm}} \right)^2 \left(\frac{v}{300 \text{ km s}^{-1}} \right) \left(\frac{n_p}{1 \text{ cm}^{-3}} \right) \# \gamma \text{ s}^{-1}$$

- Assuming isotropic emission, the 3FGL constrains: $\frac{L_{\bar{*}}}{4\pi d_{*}^2} \leq 2 \times 10^{-8} \# \gamma \text{ cm}^{-2} \text{ s}^{-1}$

- And therefore: $d_{*} \geq 6 \times 10^{18} \sqrt{\left(\frac{R_{\bar{*}}}{10^{11} \text{ cm}} \right) \left(\frac{v}{300 \text{ km s}^{-1}} \right) \left(\frac{n_p}{1 \text{ cm}^{-3}} \right)} \text{ cm}$

There could be an anti-star at ~ 1 pc from us!

Conclusions

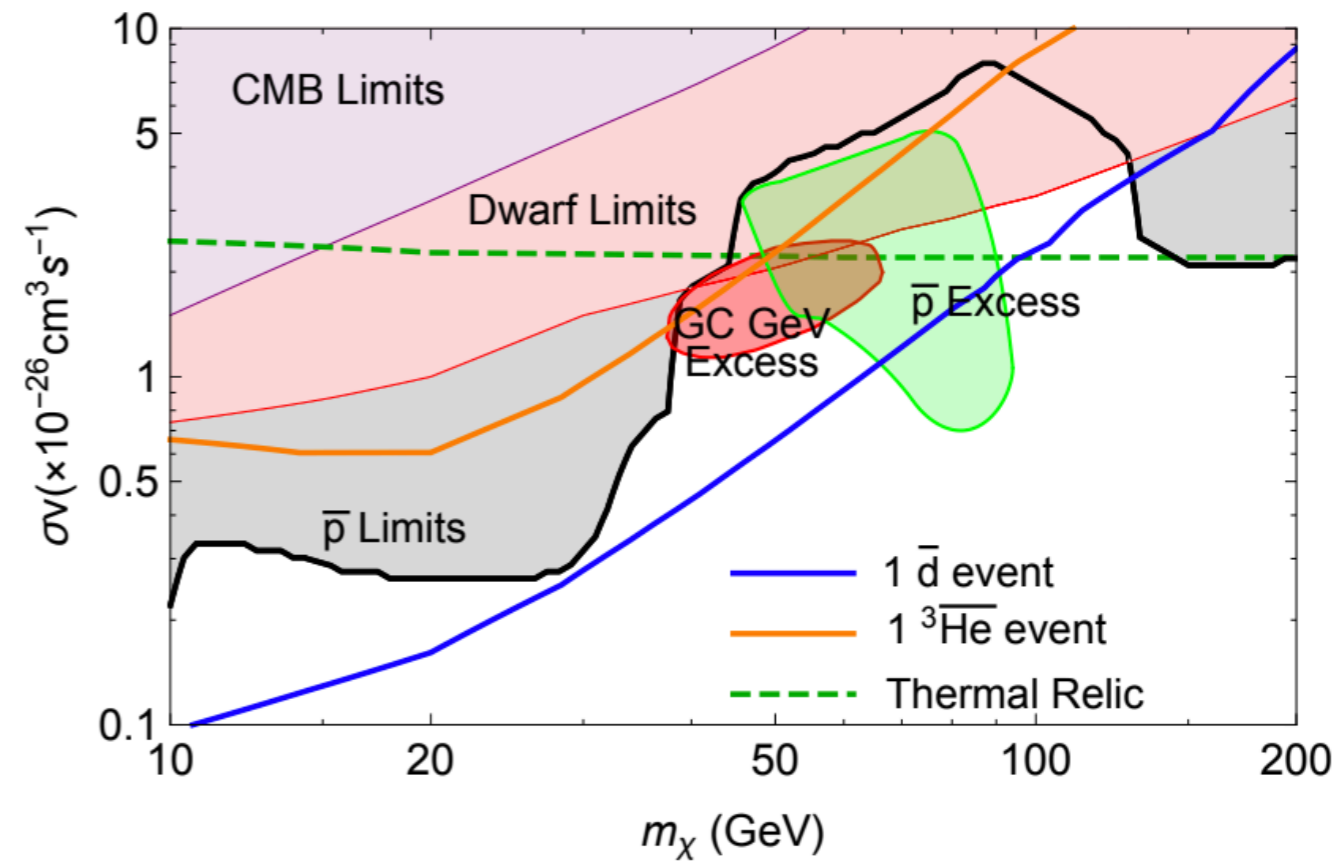
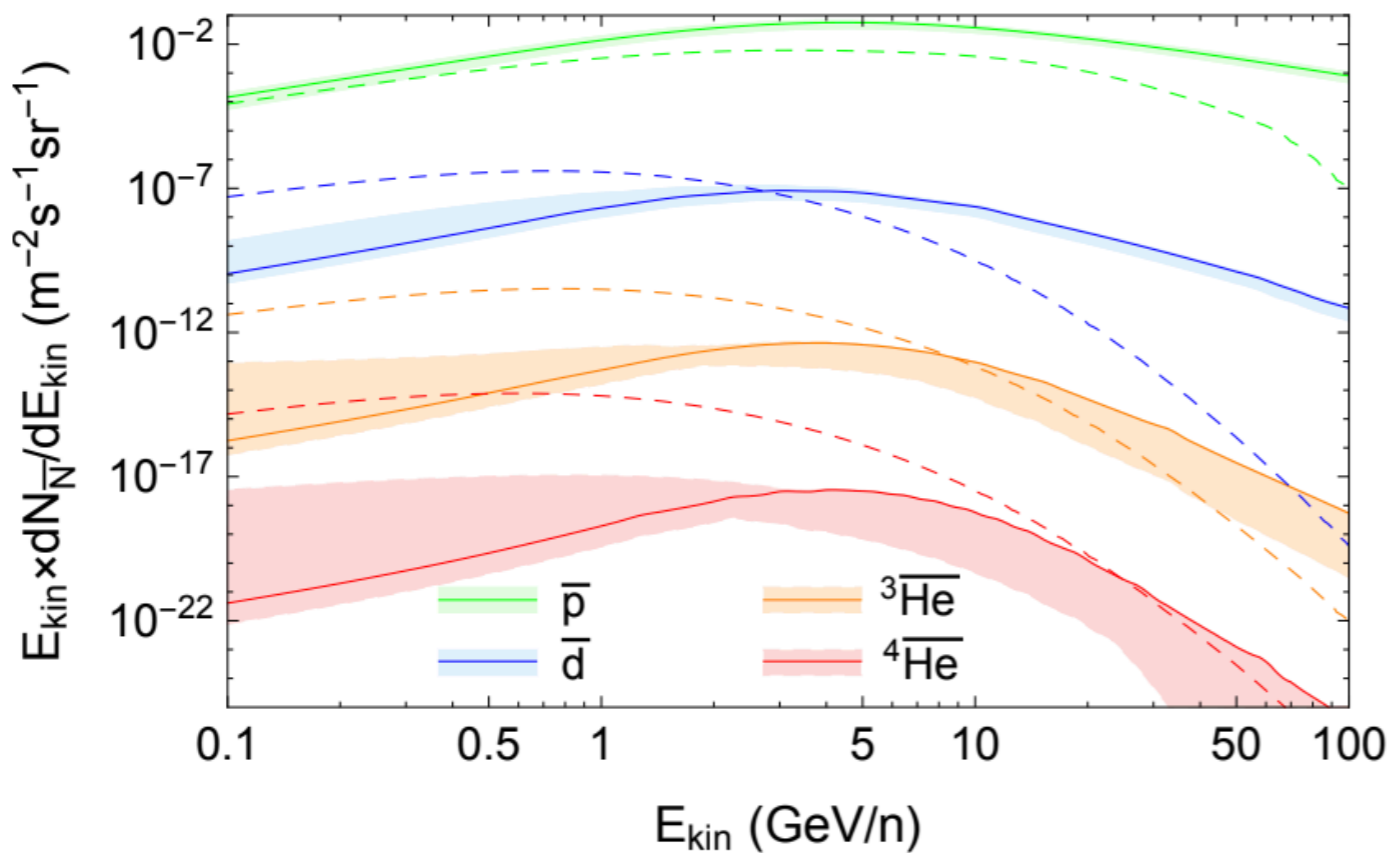
- AMS-02 has tentatively measured 6 anti- ^3He and 2 anti- ^4He : These events cannot be explained by the standard spallation and coalescence scenario. Dark Matter faces similar difficulty.

AMS-02 (tentative) discovery has major consequences for our understanding of baryogenesis in the early universe: it is far from trivial to explain these events.

- Anti-clouds might explain AMS but cannot survive unless they are almost free of normal matter along cosmic history: segregation mechanism?
- Alternatively, primordial anti-stars could be formed in the early universe from strong iso-curvature perturbations at small scales.
- Depending on the (unknown) acceleration mechanism, it is conceivable that a single near-by anti-star contributes to the AMS-02 observation.

Back-up

A common explanation to CR/ γ -rays anomalies?

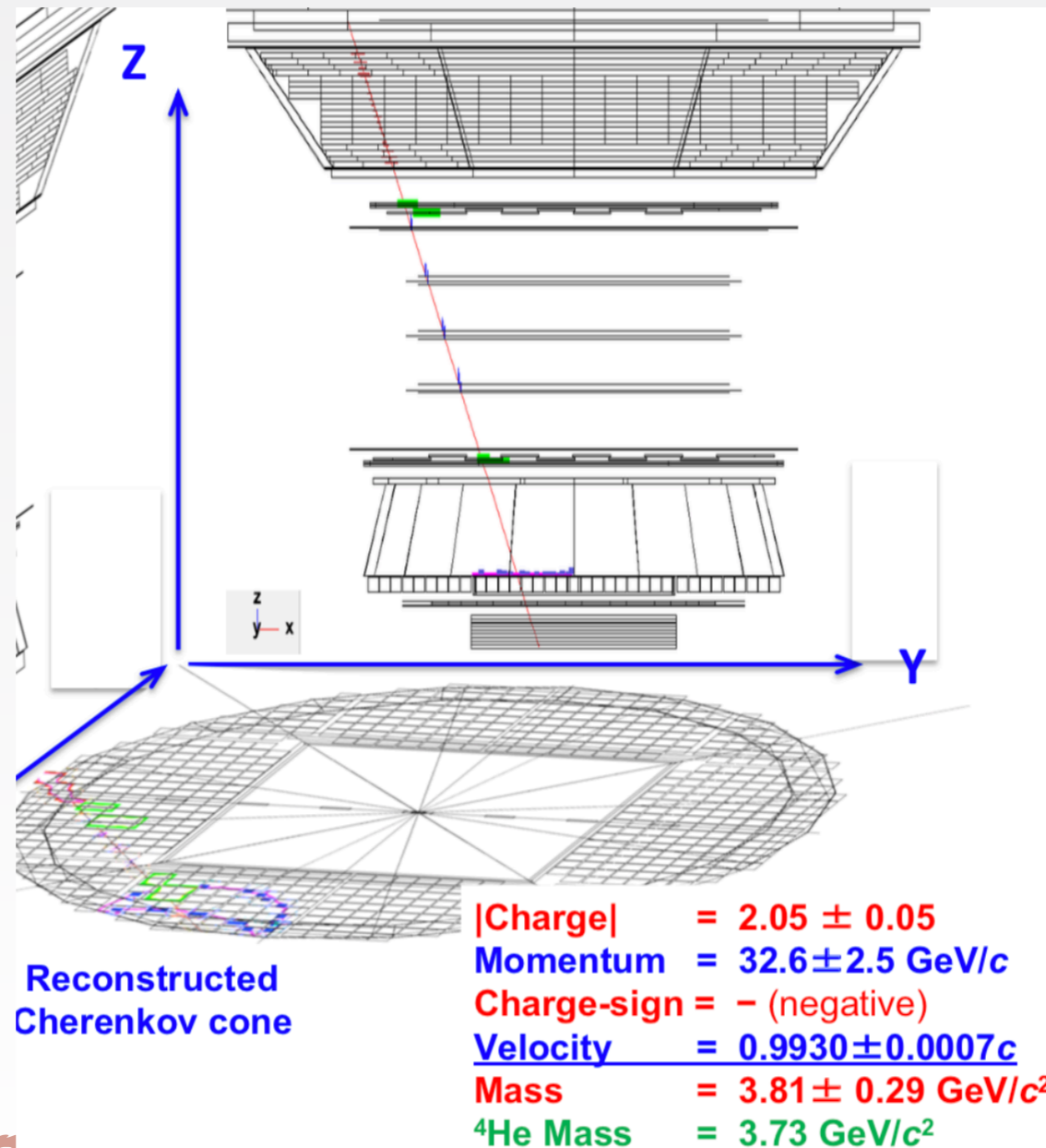


Cholis, Linden, Hooper 2001.08749

- Where is the \bar{d} ??
- How to produce ${}^4\bar{\text{He}}$??

What could be wrong?

- The measurements could be problematic:
 - Sensitivity to anti-De is **much worse** than that to anti-³He: did we miss them?
 - The mass of the anti-⁴He could have been **mis-reconstructed**.
 - Of course, the sign could be wrong...



Constraints from a γ -ray line

- γ -ray constraints can be **much stronger than the survival rate**. Let's see for instance the case of a line from $p\bar{p} \rightarrow \pi^0\gamma, \eta\gamma, \omega\gamma, \eta'\gamma, \phi\gamma, \gamma\gamma$.

- These processes produce **line with energy between 0.66 GeV and 0.933 GeV**. Decay of mesons will lead to continuum below the proton mass. We ignore this for simplicity.

- Using the FermiLAT data and the largest region “R180”, we calculate

$$\Phi_{\pi^0\gamma}^{m_p} = \frac{\int^{R180} d\ell d\Omega \rho_{\pi^0\gamma}^{\text{MW}}}{\int^{R180} d\Omega} < 6.8 \times 10^{-7} \text{cm}^{-2}\text{s}^{-1}$$

Ackermann++ 1506.00013

- We assume clouds homogeneously distributed in the disk, with a small thickness of 0.1 kpc perpendicular to the disk.

- FermiLAT allows to set (in the case of a cold cloud) $n_p^{\text{local}} \lesssim 10^{-10} - 2 \times 10^{-9} \text{cm}^{-3}$.

FermiLAT can be used to improve constraints by 2 orders of magnitude!

CMB constraints

- From Planck data we have: $\left. \frac{d^2 E}{dV dt} \right|_{\text{ann}} < 8.1 \times 10^{-31} (1+z)^6 \text{ J m}^{-3} \text{ s}^{-1}.$
- The annihilation rate is: $\left. \frac{d^2 E}{dV dt} \right|_{b\bar{b}\text{-ann}} = \langle \sigma_{p\bar{p}} v \rangle n_p n_{\bar{p}} 2m_p c^2$
- This leads to $n_{\bar{p}}^0 < 1.35 \times 10^{-10} \text{ cm}^{-3}$ on cosmological scales: ok for AMS02.
- Similarly, for anti-stars we find (assuming annihilation to pion injects energy).

$$\left. \frac{d^2 E}{dV dt} \right|_{\bar{\star}} = 8\pi R_{\bar{\star}}^2 v n_p m_p c^2 n_{\bar{\star}} \simeq 10^{13} n_{\bar{\star}} \text{ J s}^{-1} \times \left(\frac{R_{\bar{\star}}}{10^{11} \text{ cm}} \right) \left(\frac{v}{30 \text{ km s}^{-1}} \right) \left(\frac{n_p^0}{2 \times 10^{-7} \text{ cm}^{-3}} \right).$$

- And therefore $n_{\bar{\star}} \lesssim 10^{24} (1+z)^3 \text{ Mpc}^{-3}$