

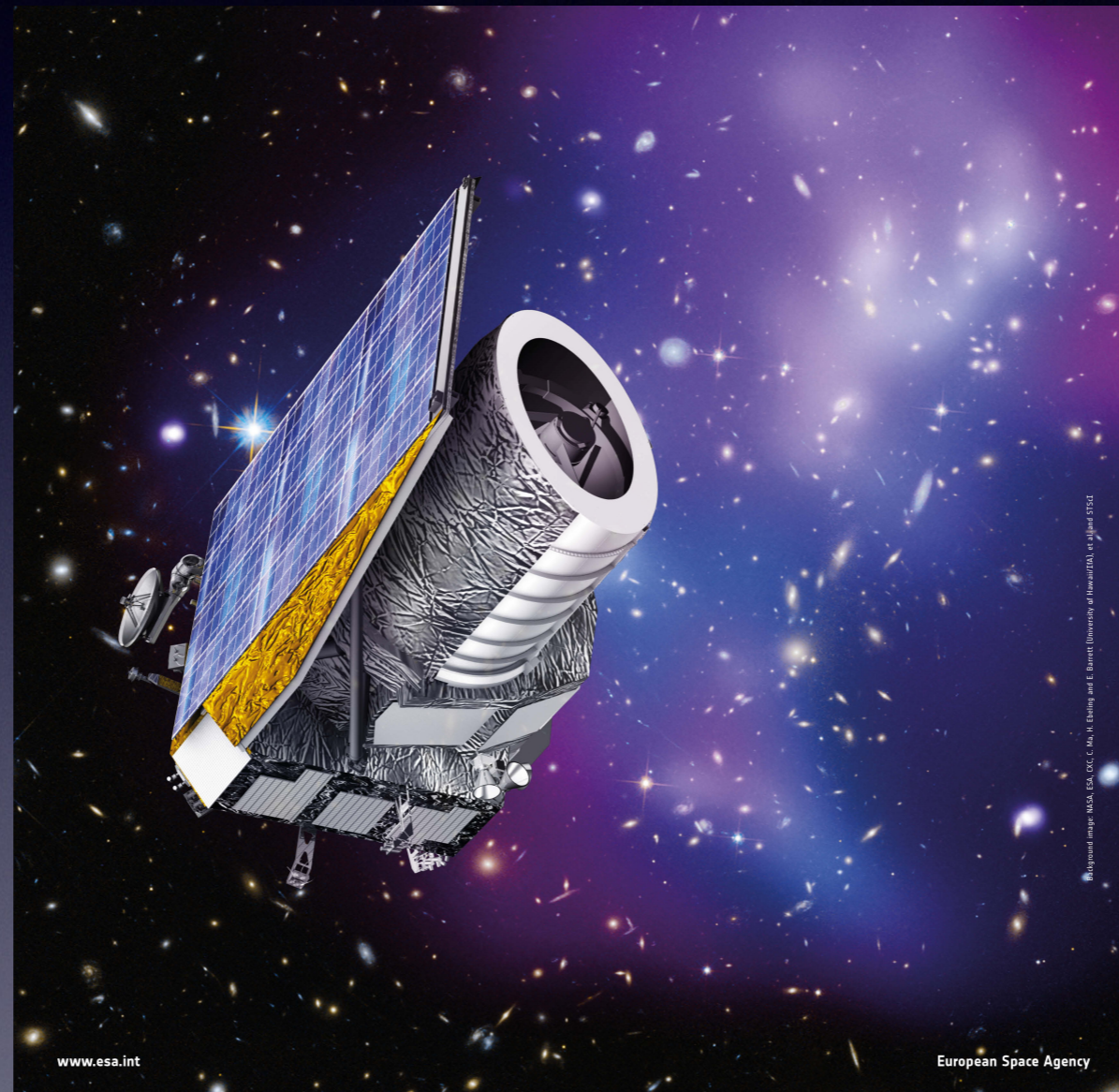
Multi-Scale Perturbation Theory

*A theoretical tool for late universe cosmology
with **nonlinear structures***

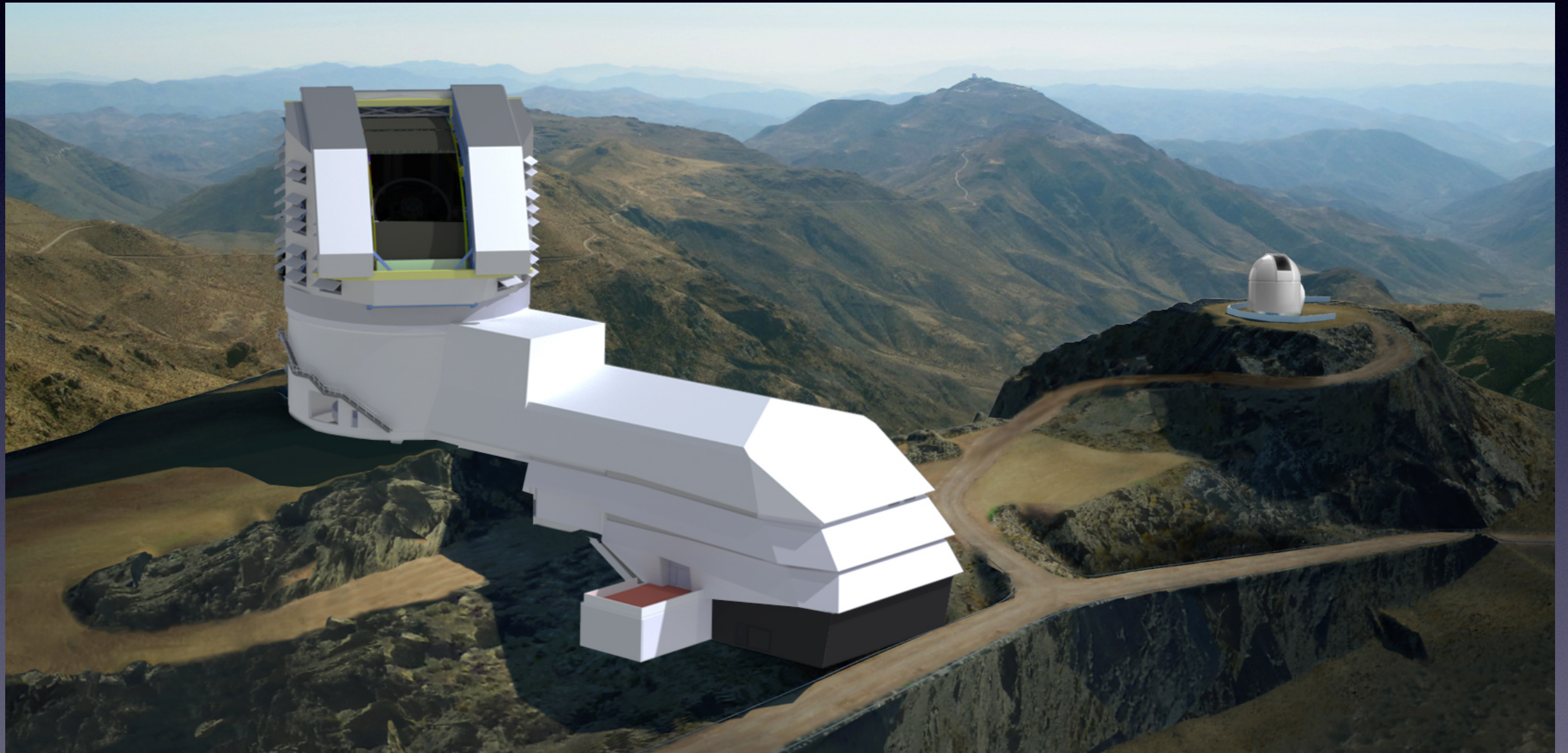
Kit Gallagher - PhD Student - QMUL

MOTIVATION

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- So our notion of “**realism**” becomes a question of approximation...
- What are the **most physical** simplifying approximations we can make?

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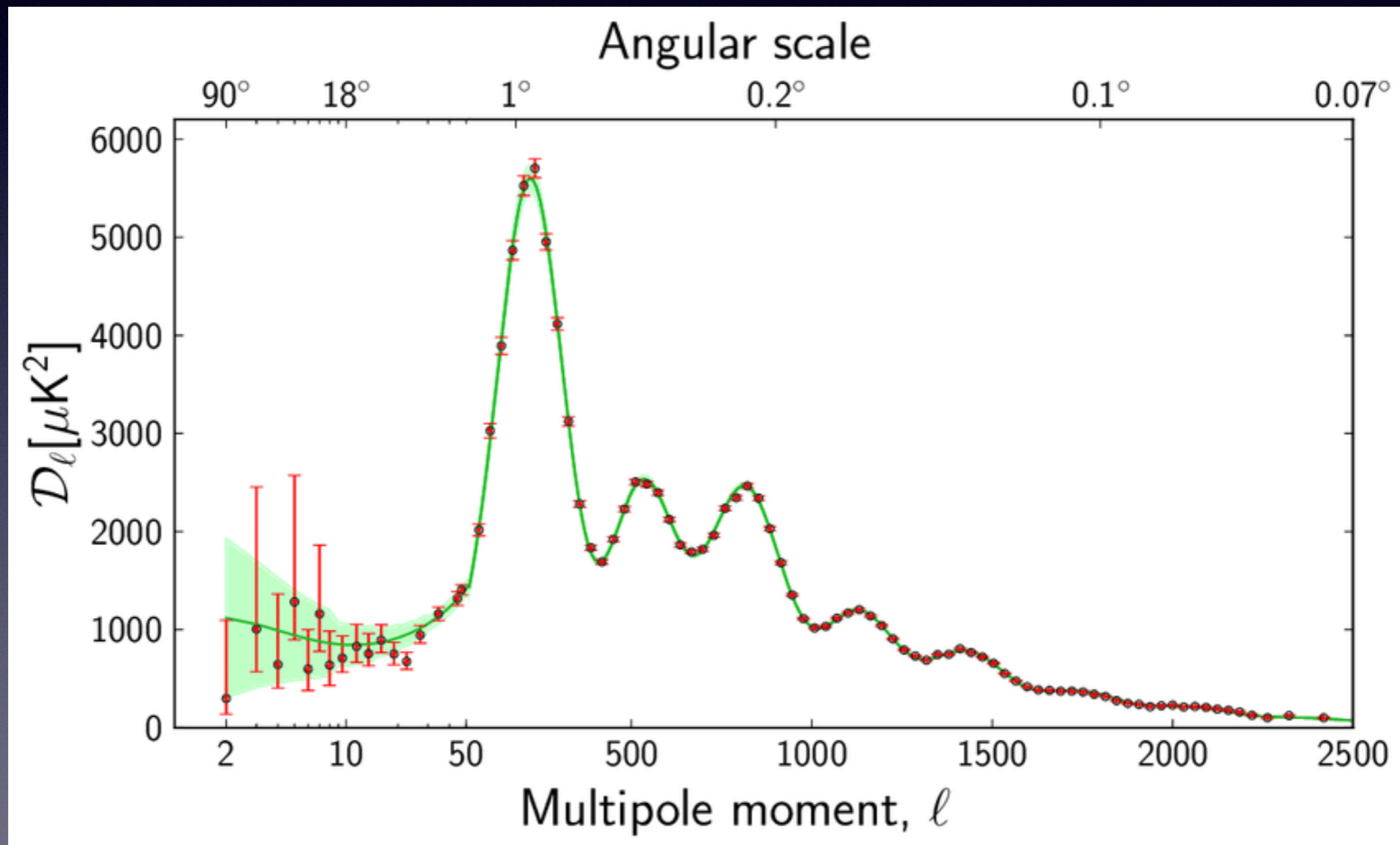
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- Equations linearise at **ALL ORDERS** in ϵ
- Smaller corrections can be calculated by considering inhomogeneous linear problem with **higher order** ($\epsilon^{n>1}$) **source terms**

SUCCESS!

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THIS APPROACH IS **EXTREMELY** SUCCESSFUL



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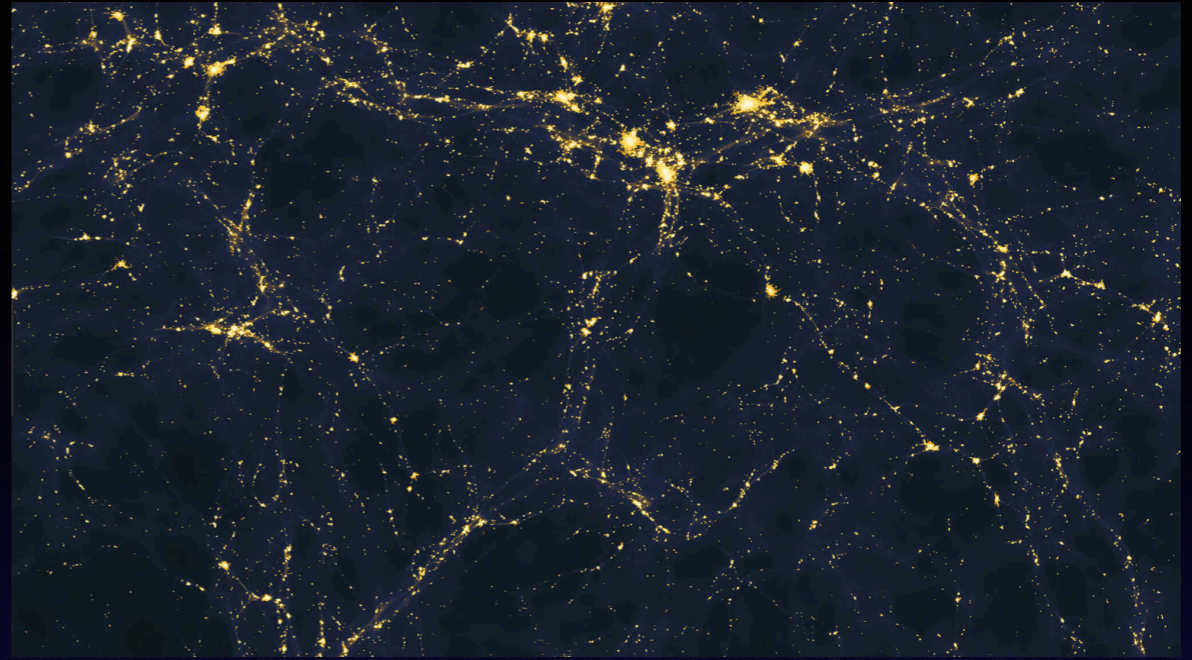
BUT...

Small fluctuation assumption **breaks down** in the late universe due to **GRAVITATIONAL COLLAPSE** on small scales!

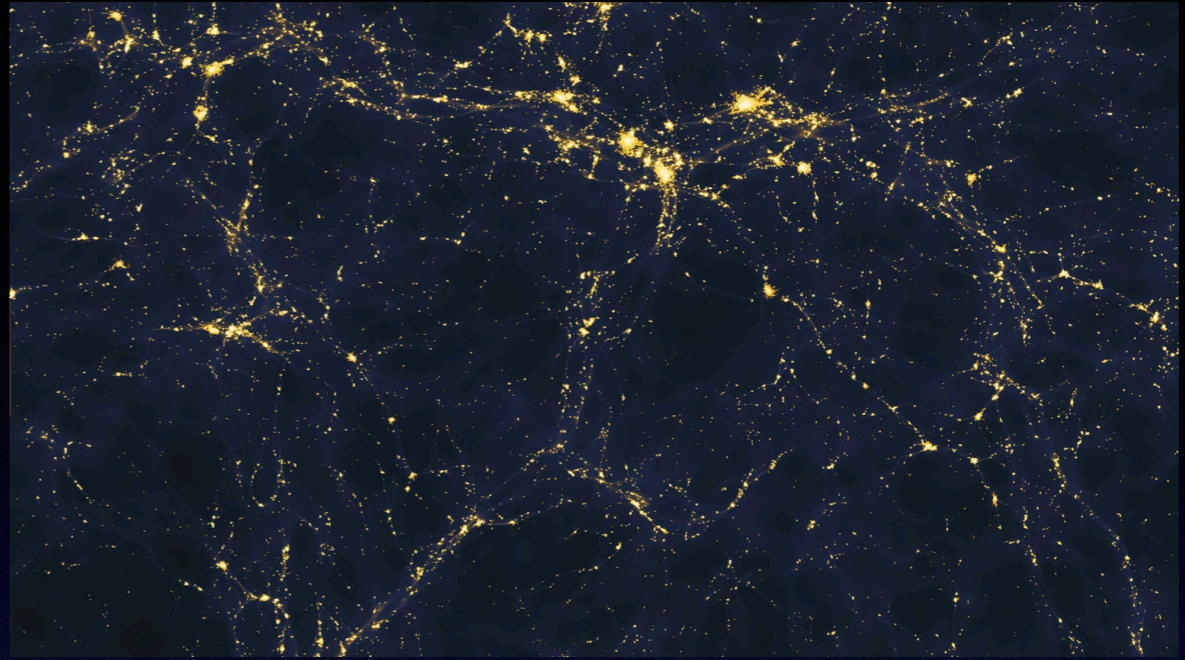
This



This



This



is not very well modelled by

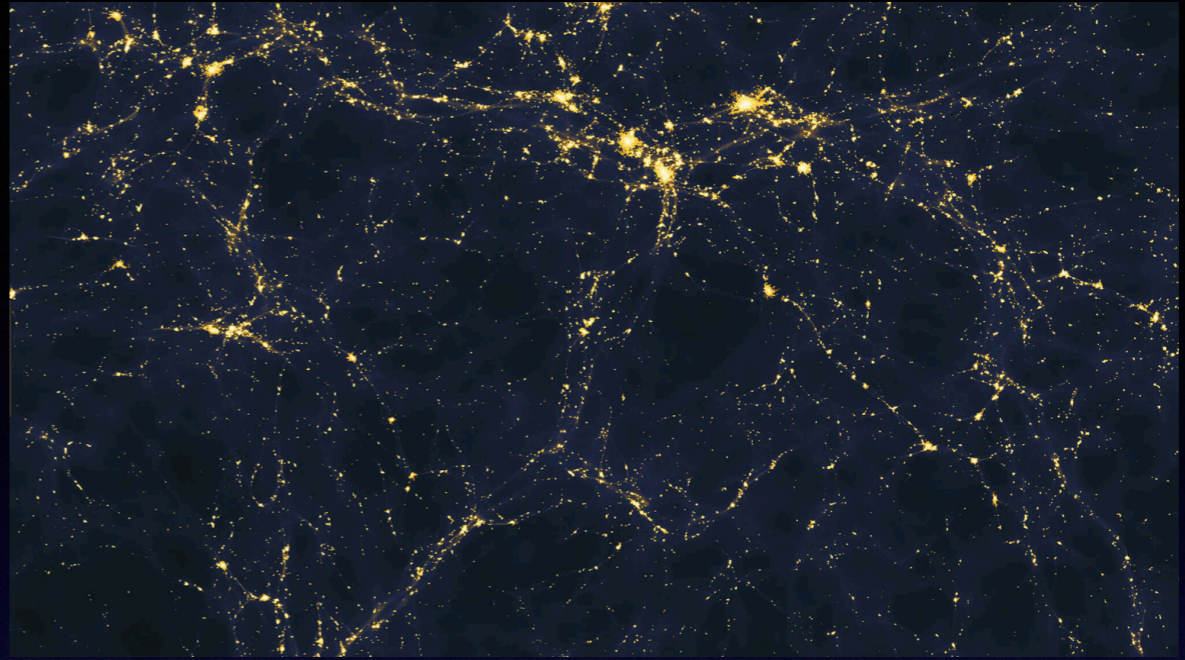
$\bar{\rho}$

+

$\delta\rho \sim \epsilon\bar{\rho}$



This



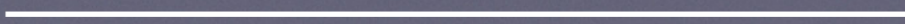
is not very well modelled by

$$\bar{\rho} + \delta\rho \sim \epsilon\bar{\rho}$$



...more like...

$$\bar{\rho} + \delta\rho \sim \bar{\rho}$$



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BUT NONLINEARITY COUPLES SCALES!

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$$\phi_2(k) \sim \int dq \phi_1(q - k) \delta_1(q)$$

Conceivable that large δ at small scales could have an effect on quantities at larger scales!

Could try to use **renormalisation** or **EFT**...

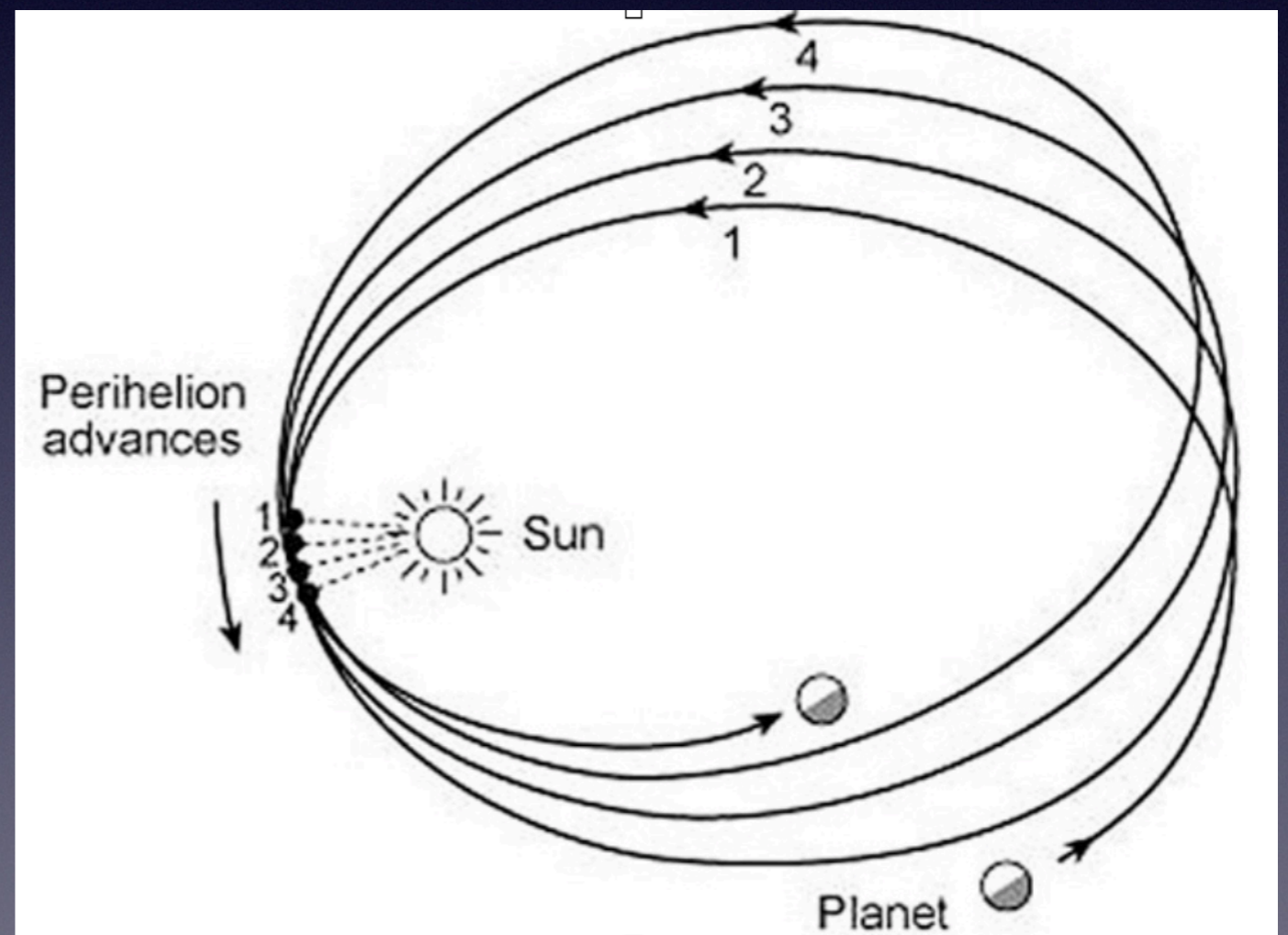
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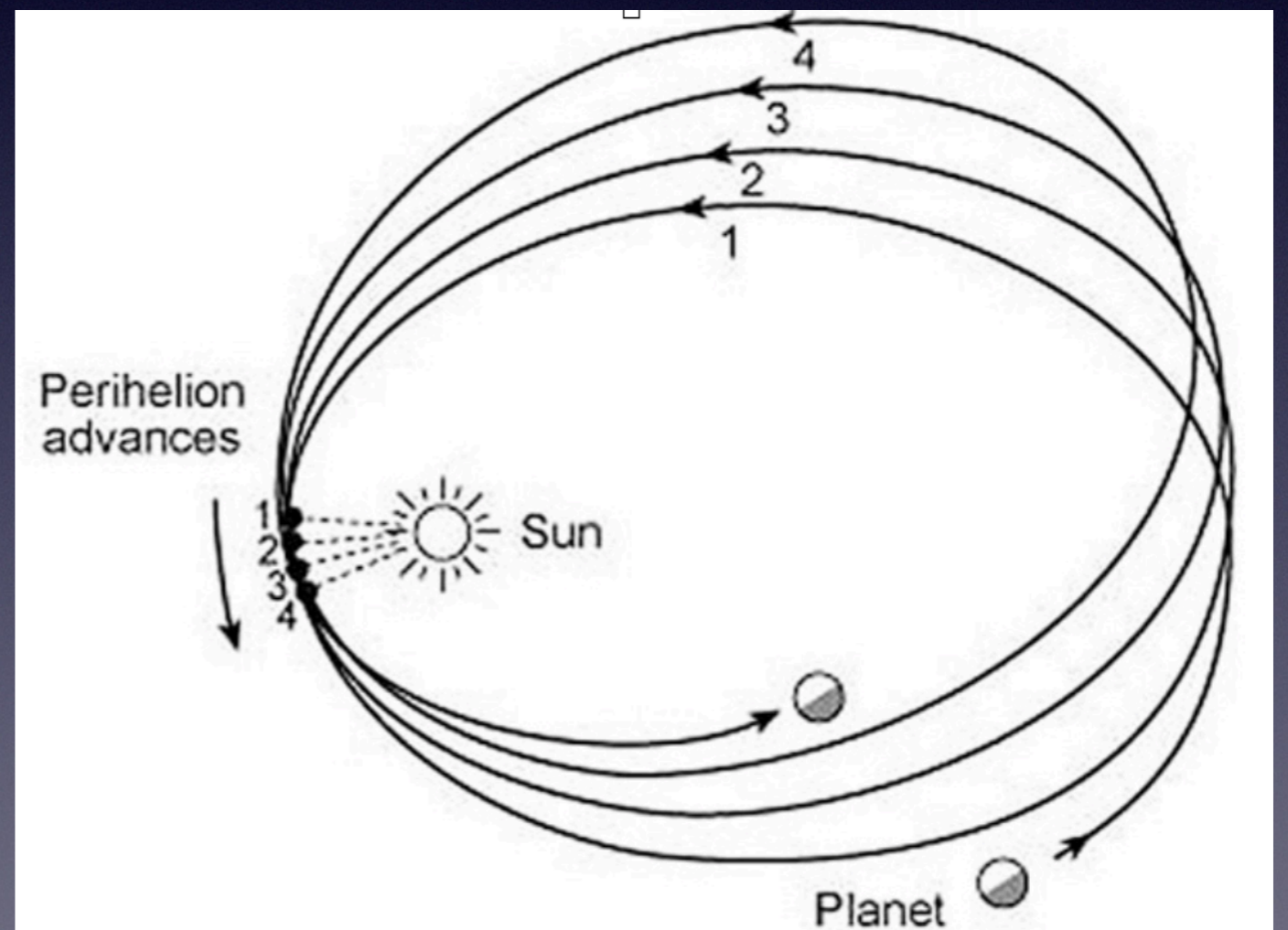
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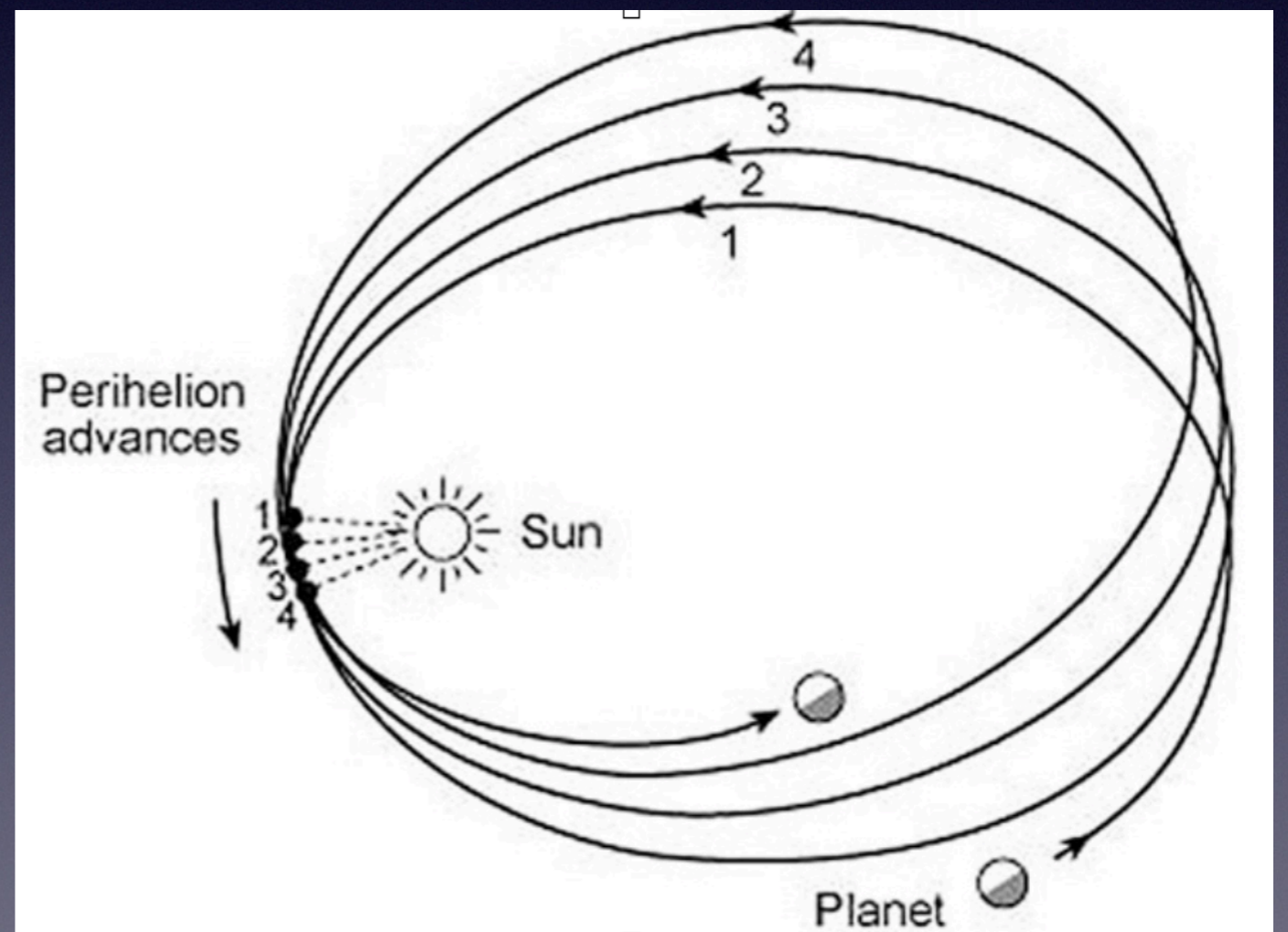
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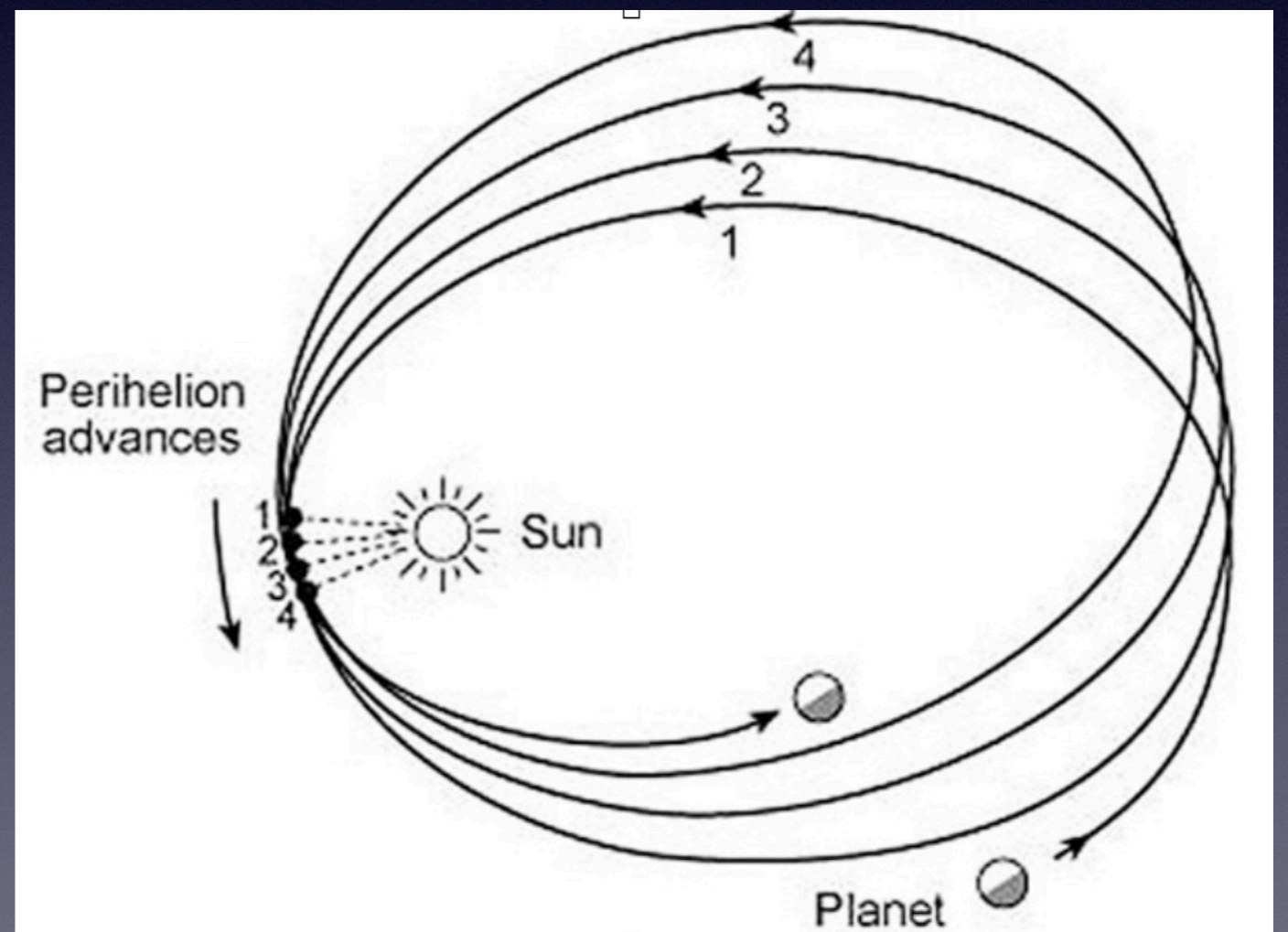
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3. The gravitational field is weak everywhere:

$$U \sim \eta^2$$

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Small velocities and weak fields lead to...

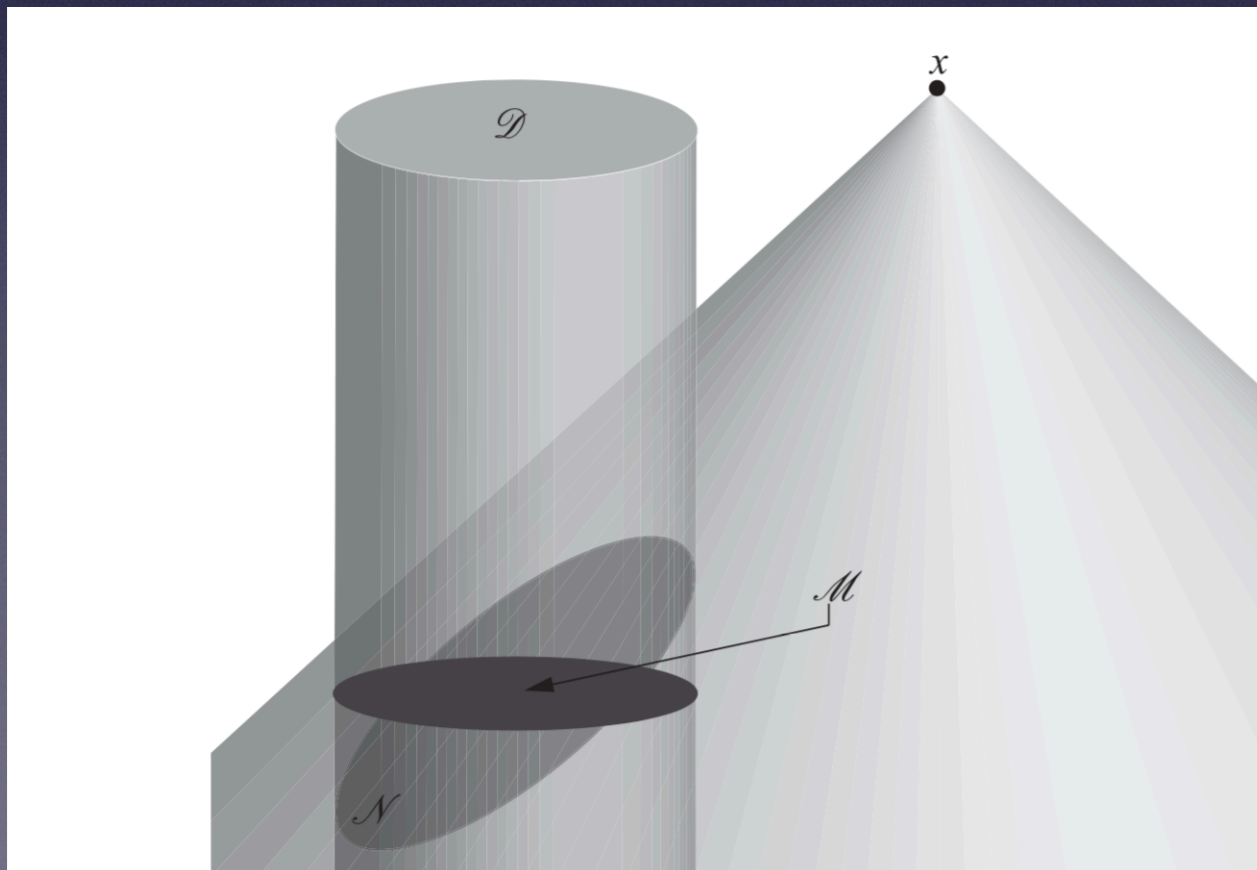
$$\frac{v}{c} = \frac{1}{c} \frac{dx}{dt} \sim \eta \quad \Longrightarrow \quad \frac{d}{dt} \sim \eta \frac{d}{dx} \quad \begin{array}{l} v \sim \eta \\ \rho \sim U \sim \eta^2 \end{array}$$

SMALL TIME DERIVATIVES:

Equations change structure from wave equations to Poisson equations

$$B_i \sim \eta^3$$

$$h_{ij} \sim \eta^4$$



Can think of approximating $\mathcal{N}(x, t)$ by $\mathcal{M}(x)$

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Can we use **normal perturbation theory** on **large scales**...

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...and see how they might **interact**?

ANSWER...YES!

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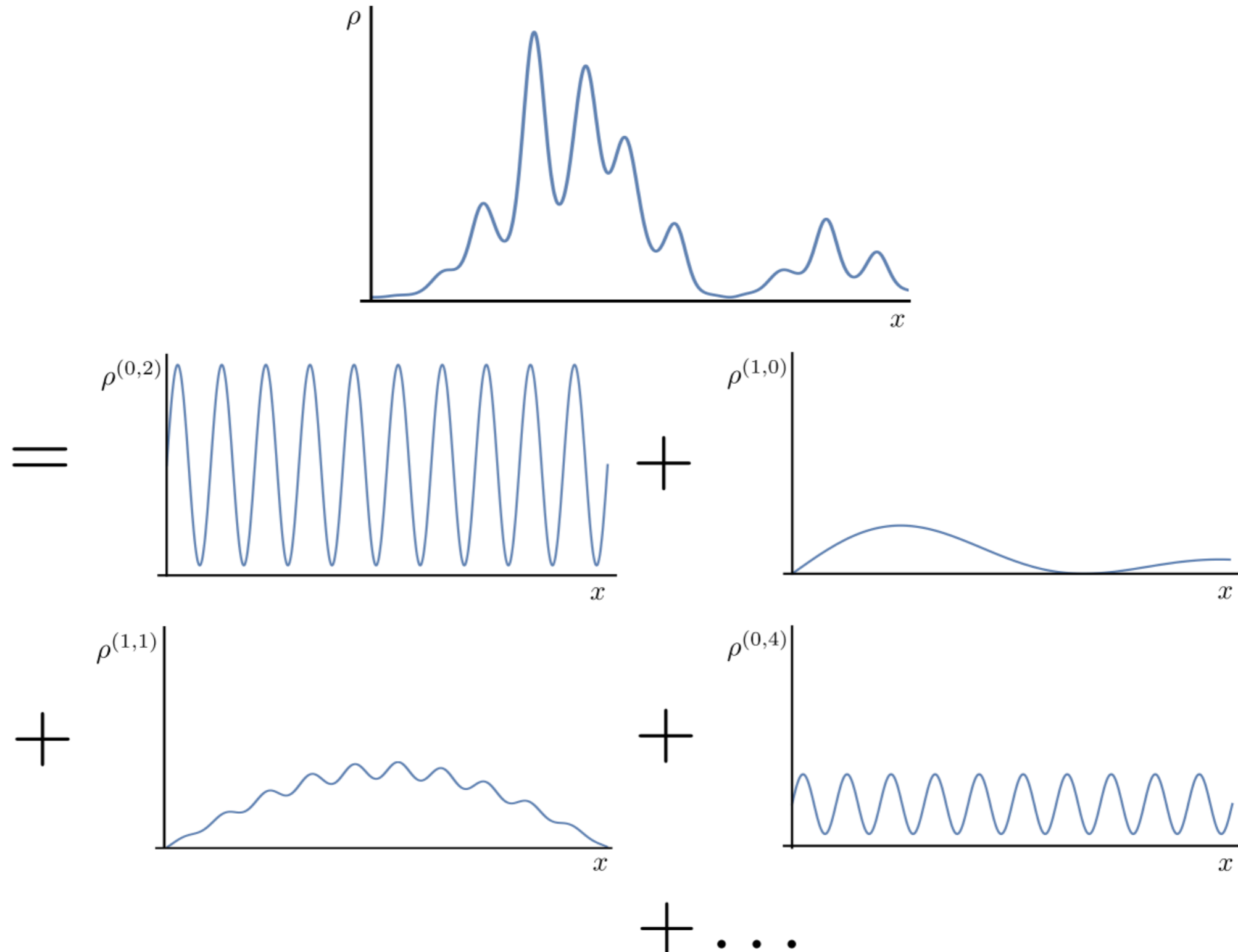
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This choice is **well motivated by physical observations!**

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...where $\{U, \rho_N, v_{Ni}\}$ are treated like the leading order parts of a **post-Newtonian expansion...**

...and $\{\phi, \psi, \rho, v_i\}$ are treated like **standard 1st order cosmological perturbations.**

At **LEADING ORDER:**

Obtain **homogeneous
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$$\rho^{(0,2)} = \rho_N = \bar{\rho} + \delta\rho_N \sim \frac{\eta^2}{L_N^2} \implies \delta_N = \frac{\delta\rho_N}{\bar{\rho}} \sim 1$$

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- **AND spatially varying coefficients in linear operator!?!**

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- Or...make **further approximations** and try to proceed **analytically**...

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- Could take solutions for $\{U, \rho_N, v_{Ni}\}$ from **N-body** - then try to **solve numerically**?
- Or...make **further approximations** and try to proceed **analytically**...

We take **approximate solutions from Newtonian perturbation theory** for $\{U, \rho_N, v_{Ni}\}$
- see what happens?

Compare to normal perturbation theory?

Since the big nasty equations are still **LINEAR**...

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...then if we can split the
Newtonian quantities as...

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...we can do the same thing for the **cosmological ones!**

$$\psi = \psi^{(1)} + \psi^{(2)} + \dots$$

$$\delta = \delta^{(1)} + \delta^{(2)} + \dots$$

General Relativity

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$\mathcal{O}(\varphi)$



**First order
cosmological
perturbation
theory**

$\mathcal{O}(\varphi^2)$



**Second order
cosmological
perturbation
theory**

$\mathcal{O}(\varphi^3) \dots \mathcal{O}(\varphi^n)$



General Relativity
 $G_{\mu\nu} + \Lambda g_{\mu\nu} = \kappa T_{\mu\nu}$

$$\begin{array}{c} \mathcal{O}\left(\frac{\eta^2}{L_N^2}\right) \\ \longrightarrow \\ \mathcal{O}\left(\frac{\epsilon\eta^2}{L_N^2}\right) \end{array}$$

**Two-parameter
perturbation
theory (2PPT)**

$$\begin{array}{c} \mathcal{O}\left(\frac{\eta^2}{L_N^2}\right) \\ \longrightarrow \\ \epsilon \rightarrow 0 \end{array}$$

**Newtonian
gravity**

$$\mathcal{O}(\varphi)$$



**First order
cosmological
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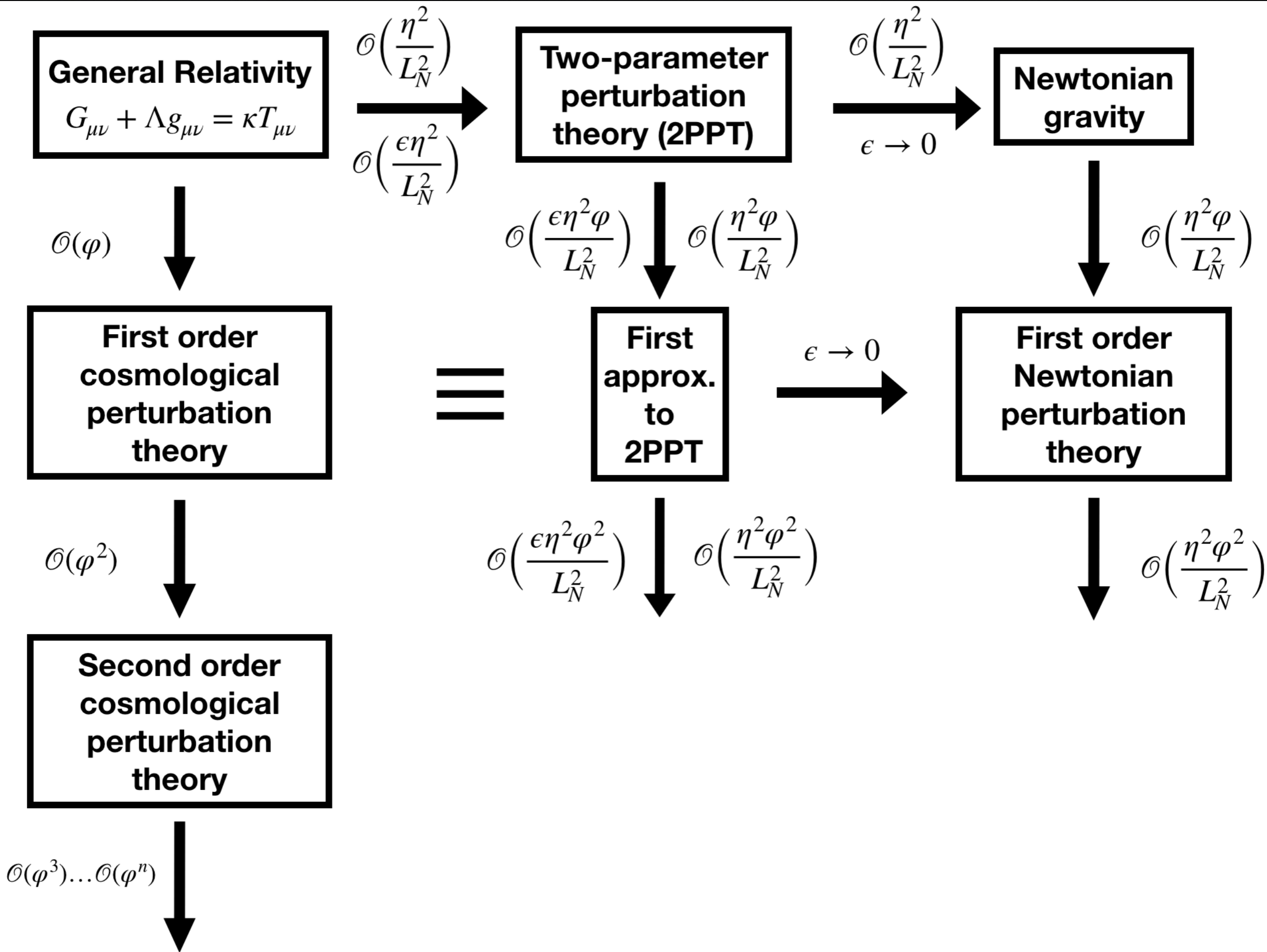
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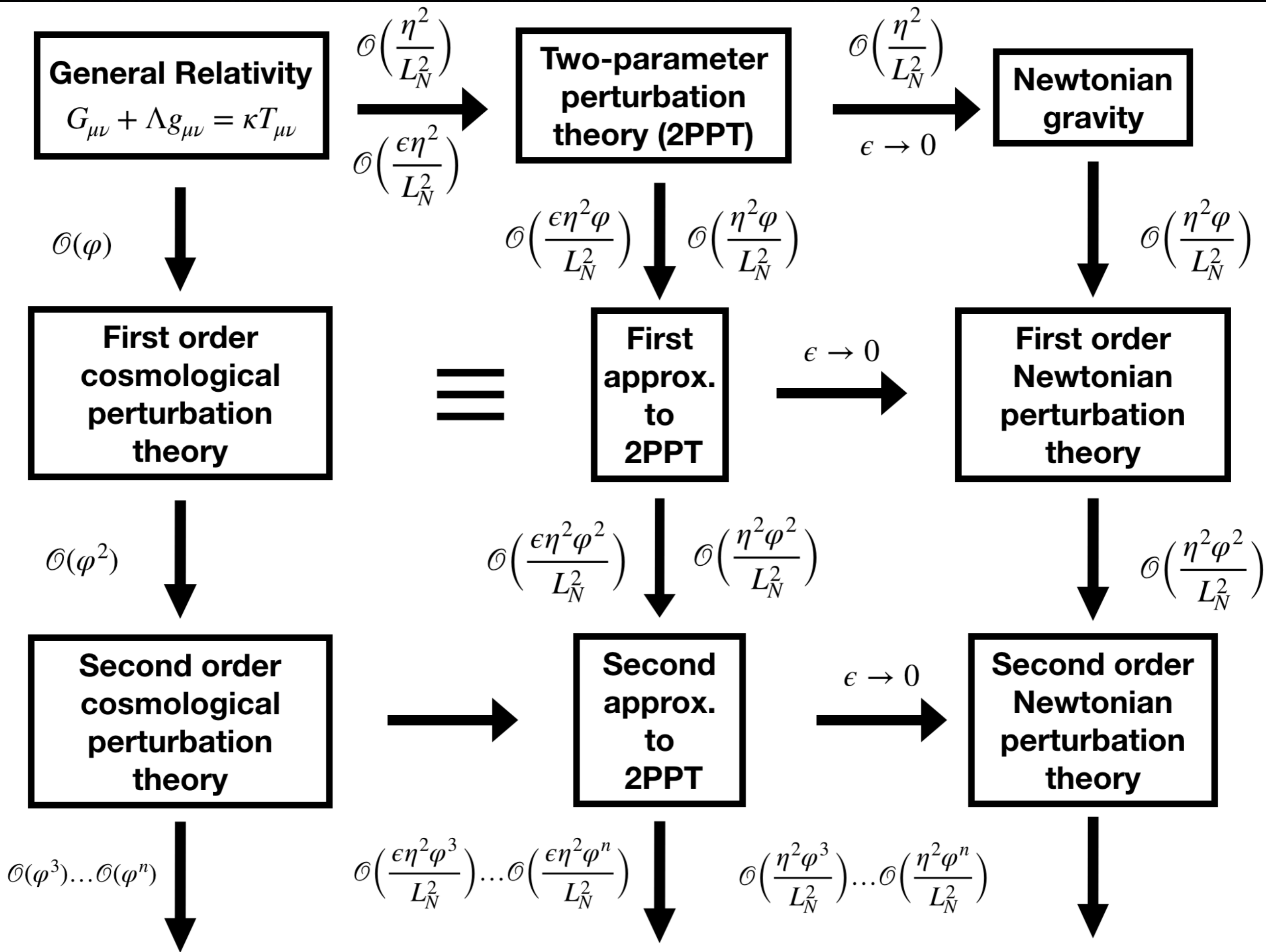


**Second order
cosmological
perturbation
theory**

$$\mathcal{O}(\varphi^3) \dots \mathcal{O}(\varphi^n)$$







WHAT DO YOU GET?

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- At **first approximation** - recover **1st order Poisson gauge cosmological perturbation theory!**

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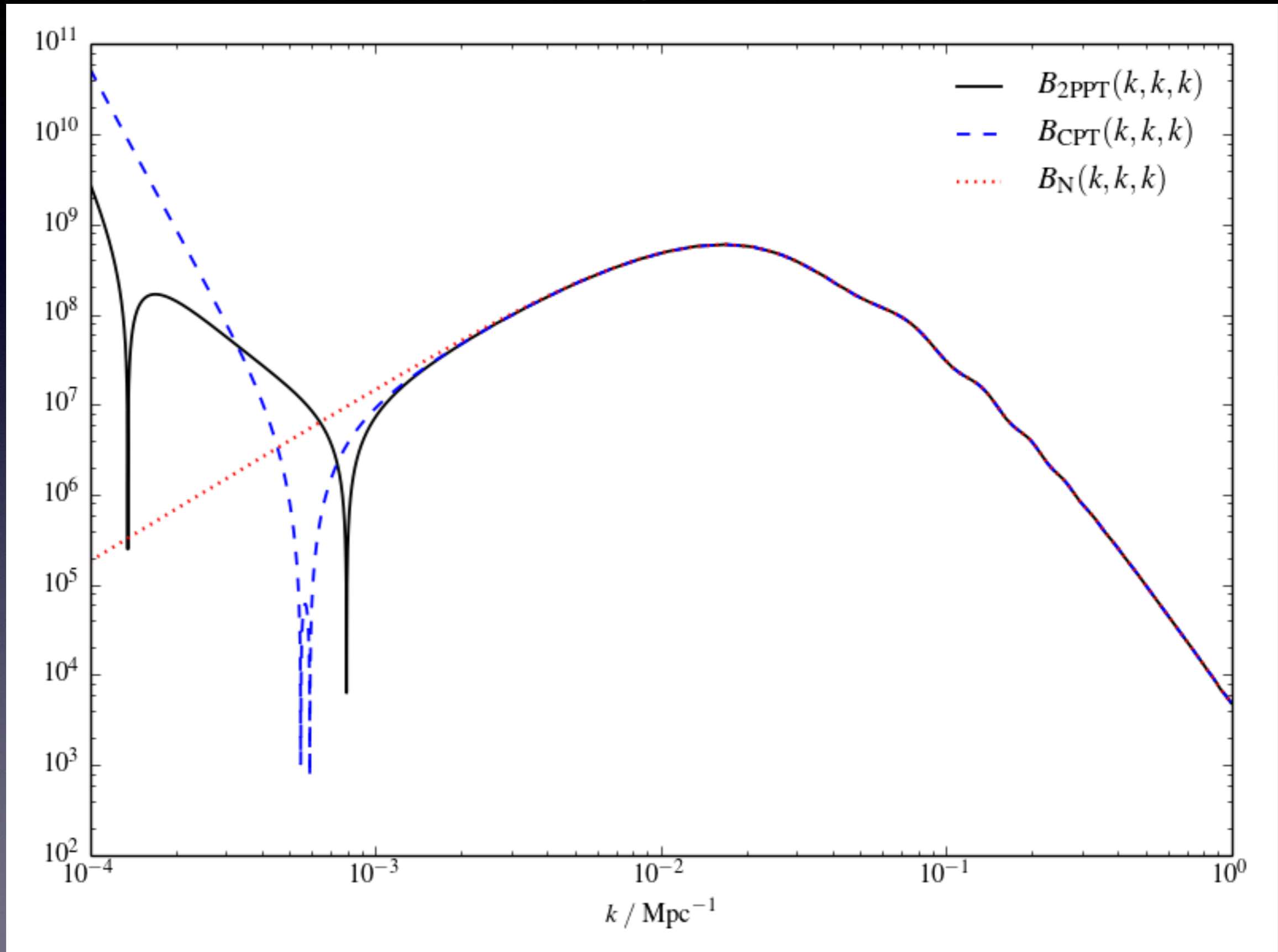
- At **first approximation** - recover **1st order Poisson gauge cosmological perturbation theory!**
- At **second approximation** - recover an approximation to 2nd order perturbation theory (**missing some terms**)

WHAT DO YOU GET?

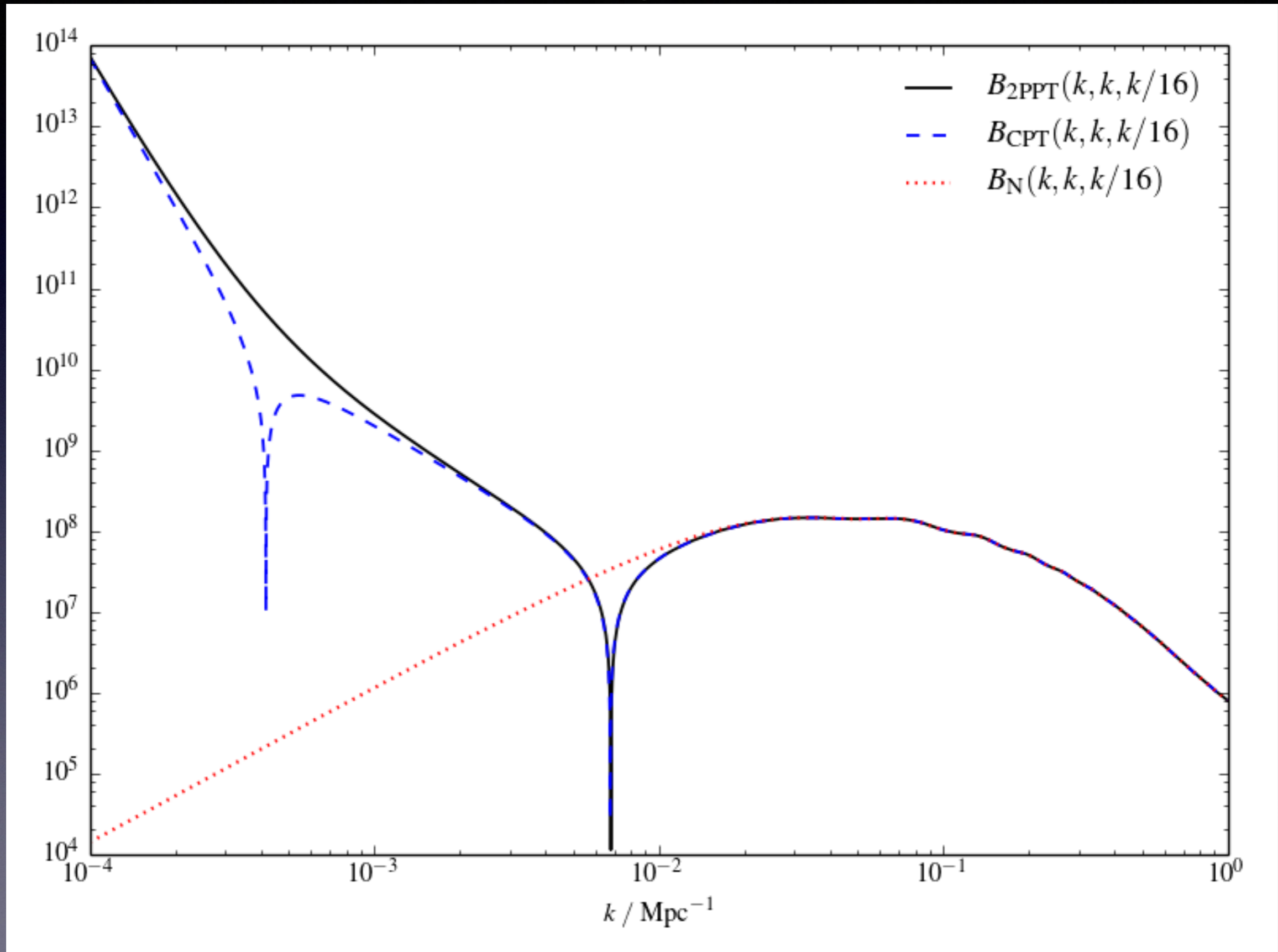
- At **first approximation** - recover **1st order Poisson gauge cosmological perturbation theory!**
- At **second approximation** - recover an approximation to 2nd order perturbation theory (**missing some terms**)
- At **third approximation** - recover an approximation to 3rd order perturbation theory (missing **A LOT** of terms)

We plotted **tree-level matter bispectrum** to compare behaviour to full cosmological perturbation theory:

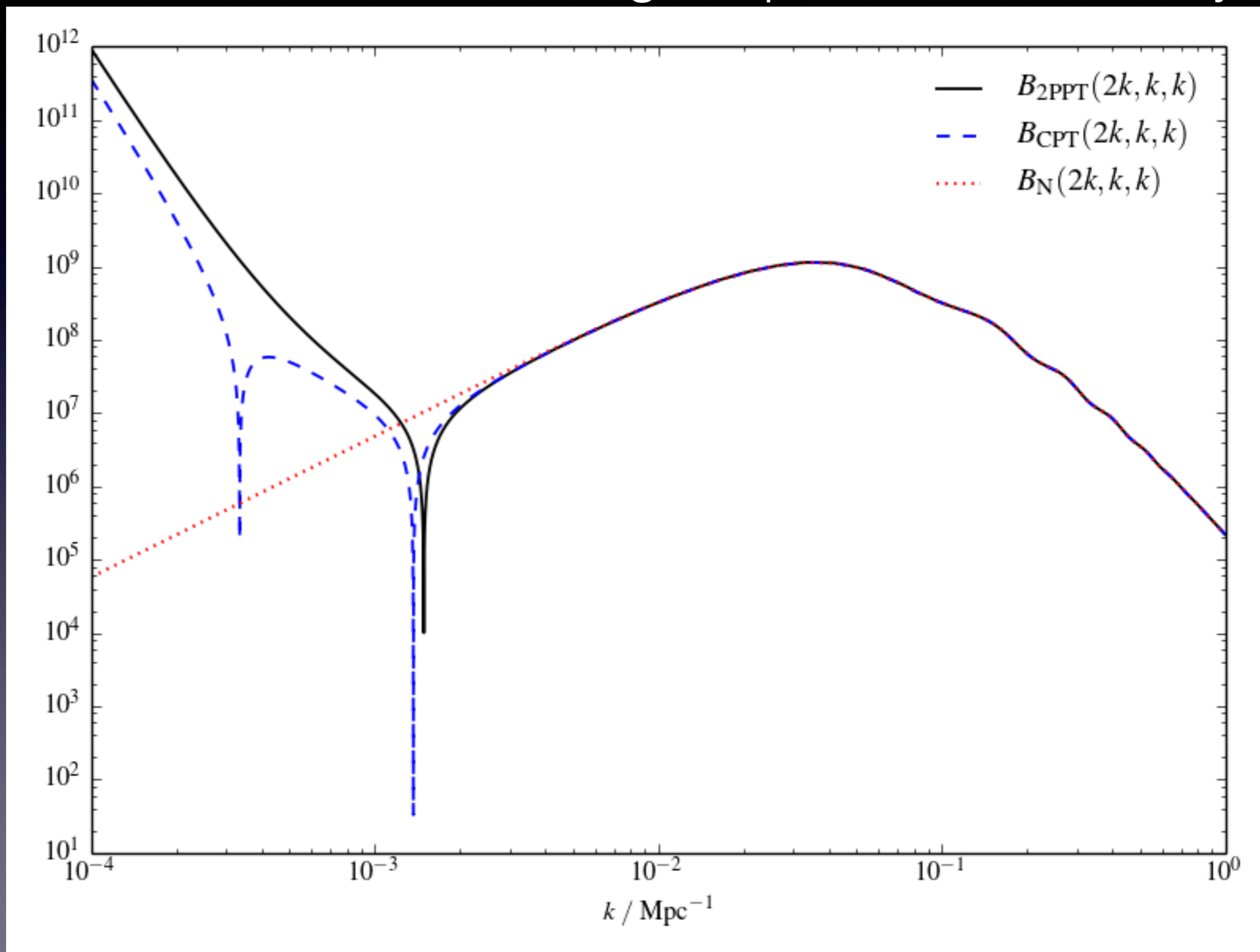
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Seem to have **good agreement**

down to

$$k > 10^{-3} \text{ Mpc}^{-1}$$

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The **2PPT correctly identifies** all the Poisson gauge
terms scaling as $\sim \frac{\mathcal{H}^2}{k^2}$

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...**BUT** then would have to consider **even higher order products of the nonlinear terms**

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4. $k_{\min} \sim 0.001 \text{ Mpc}^{-1}$ is expected for surveys

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4. “Upgrade” the Newtonian perturbation theory sector (**Renormalised PT, EFTofLSS**)
5. Check the **soft theorems!**

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THANK YOU!