

# Gravitational Waves from Axions

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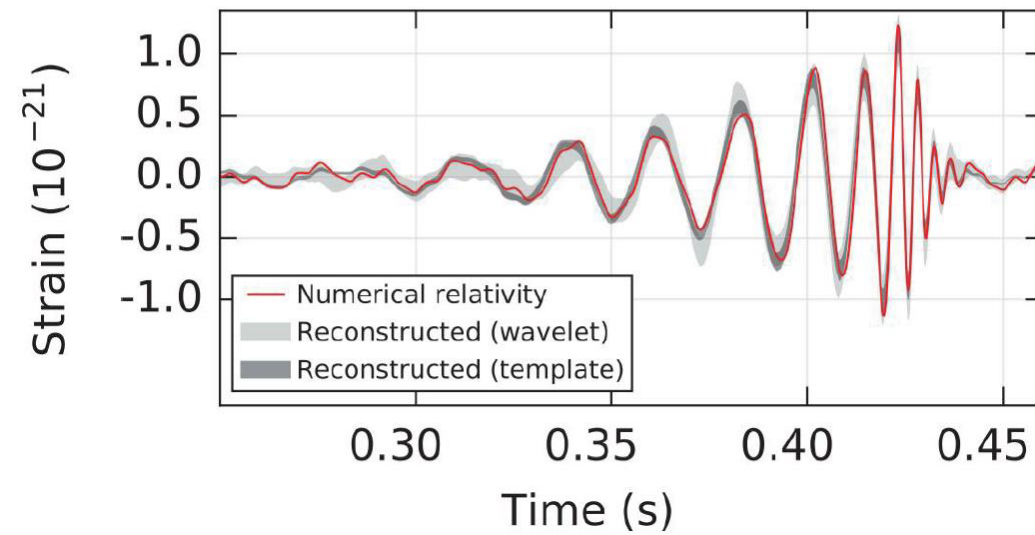
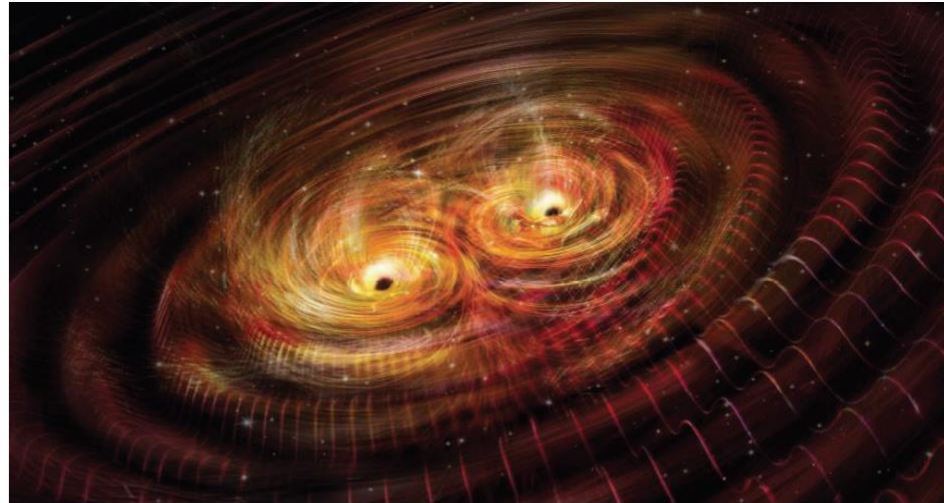
with **E.Hardy** and **H.Nicoleascu**

[2101.11007]

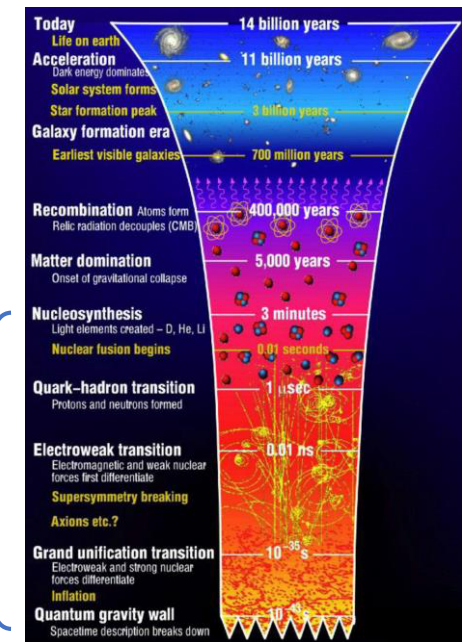
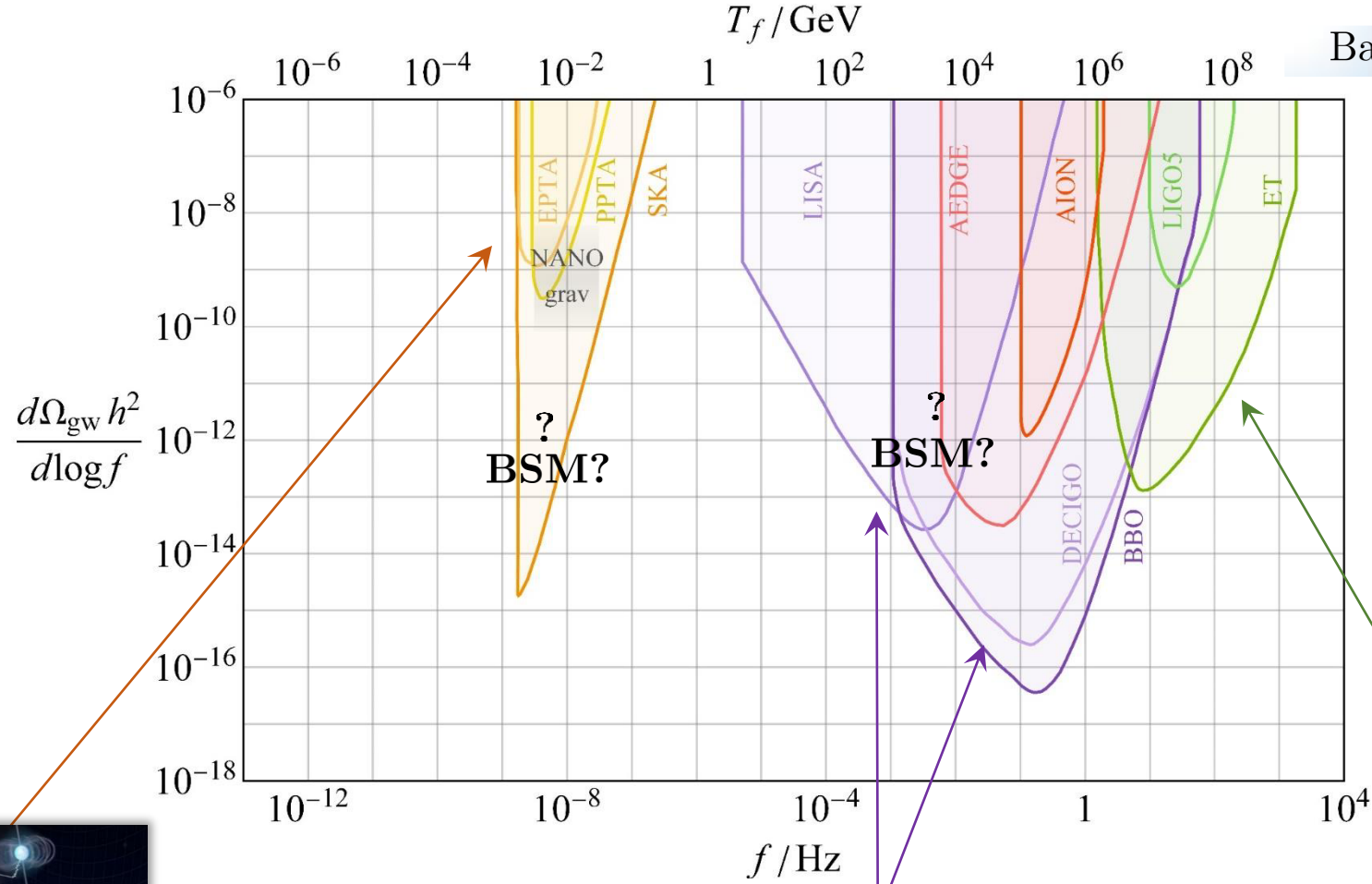
# Outline

- Prospects of gravitational wave detection
- Axions and axion strings
- String dynamics and generation of gravitational waves
  - Bounds on axion mass and decay constant
- Conclusion and Outlook

# Gravitational Waves



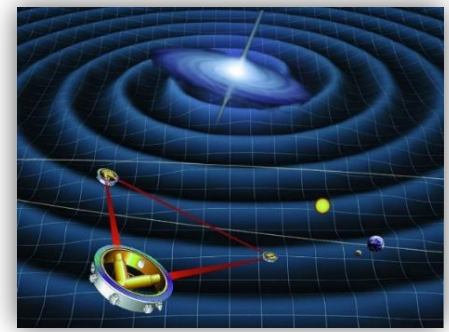
Back to the past



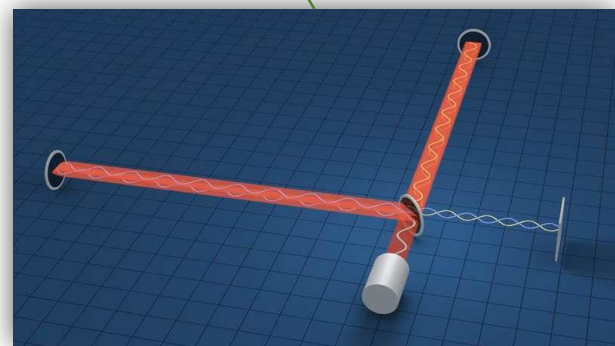
Pulsar



Pulsar Timing Arrays



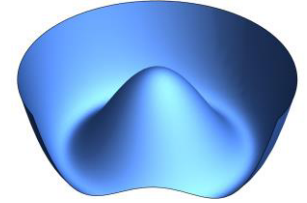
Space-based Interferometers



Ground-based Interferometers

# Axion $\equiv a$

- Goldstone boson of a (new) spontaneously broken U(1) symmetry (at the scale  $f_a$ )



Simplest realization:

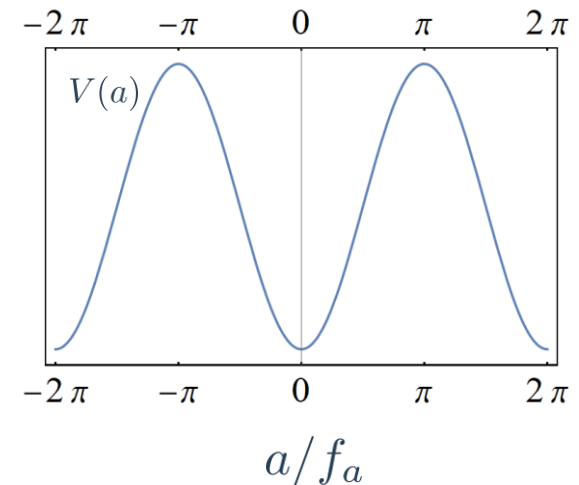
$$\phi = \frac{1}{\sqrt{2}}(r + f_a)e^{i\frac{a}{f_a}}$$

$$V_\phi = \frac{m_r^2}{2f_a^2} \left( |\phi|^2 - \frac{f_a^2}{2} \right)^2$$

- U(1) explicitly broken by the axion potential  $V(a)$

→  $V(a)$  invariant under  $a \rightarrow a + 2\pi f_a$

→ axion mass  $m_a$



## Example

QCD axion:

$$\mathcal{L} \supset \frac{a}{f_a} G_{\mu\nu} \overbrace{\tilde{G}^{\mu\nu}}^{\epsilon^{\mu\nu\rho\sigma} G_{\rho\sigma}}$$

$$\longrightarrow m_a \simeq m_\pi f_\pi / f_a$$

Strong CP:

$$\mathcal{L}_{SM} \supset \theta_{\text{QCD}} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

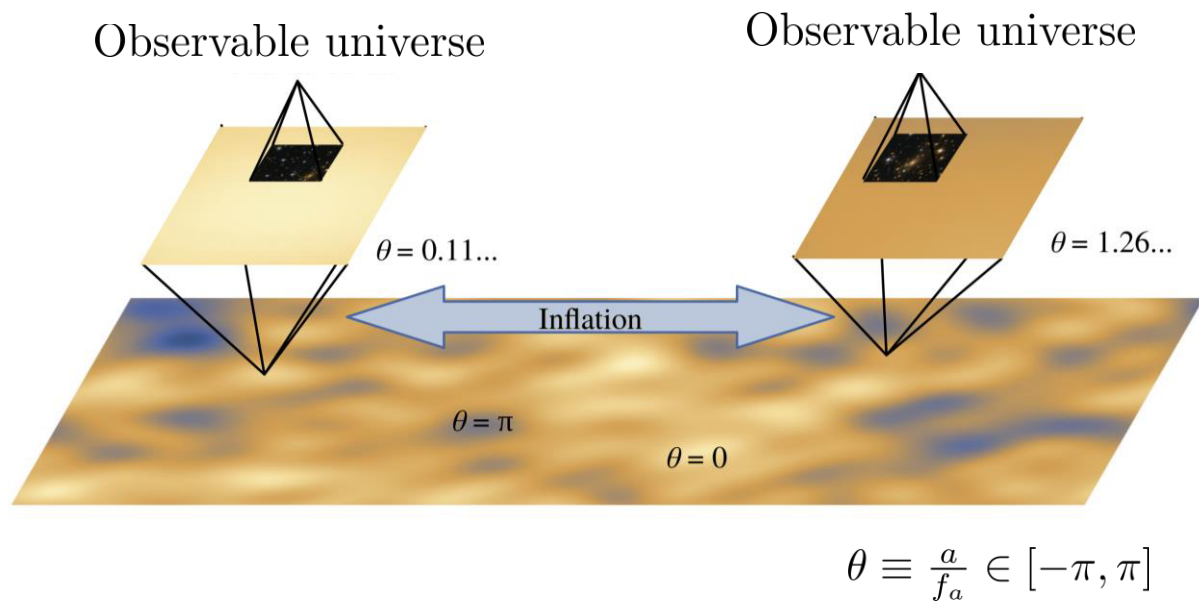
↓  
 $\lesssim 10^{-10}$

$$\longrightarrow \theta_{\text{QCD}} = \frac{\langle a \rangle}{f_a} \rightarrow 0$$

# Cosmological Initial Conditions

## Pre-inflationary

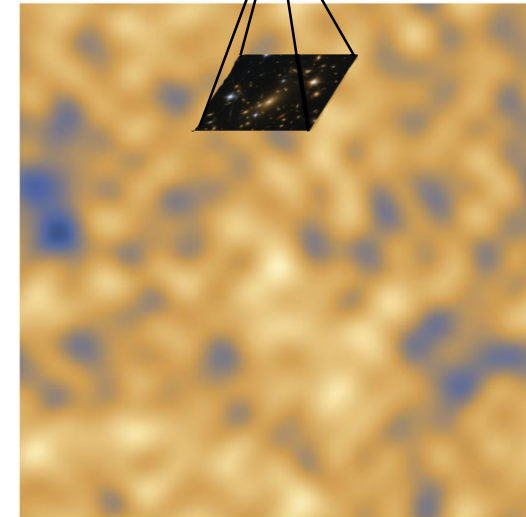
$$T_R \lesssim f_a \text{ and } H_I \lesssim f_a$$



## Post-inflationary

$$T_R \gtrsim f_a \text{ or } H_I \gtrsim f_a$$

Observable universe

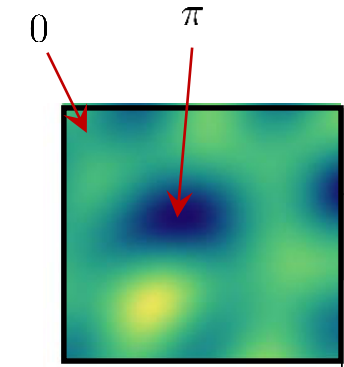
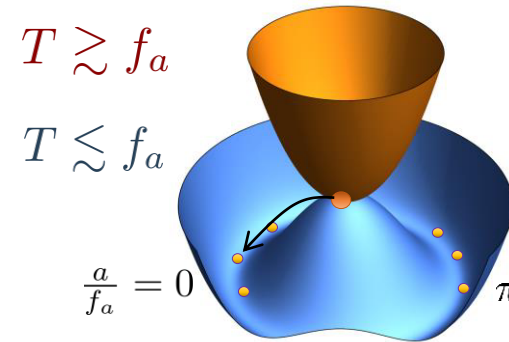


# Origin of the inhomogeneities in the post-inflationary scenario

$$T_R \gtrsim f_a \text{ or } H_I \gtrsim f_a$$

- Finite temperature corrections

$$T_{RH} \lesssim 10^{15} \text{ GeV}$$



spatial inhomogeneities

- Quantum inflationary fluctuations:

$$\langle \sigma_a^2 \rangle = \left( \frac{H_I}{2\pi} \right)^2 \rightarrow H_I \gtrsim \frac{f_a}{2\pi}$$

$$\frac{H_I}{2\pi} \lesssim 10^{13} \text{ GeV}$$

- Direct coupling of  $\phi$  to the inflation  $\varphi$ ?

[Kofman,Linde '87]

$$V_\phi \supset g\varphi^2|\phi|^2 \xrightarrow{\text{during inflation}} \underbrace{g\langle\varphi\rangle^2}_{\text{effective mass for } \phi} |\phi|^2 \quad g\langle\varphi\rangle^2 \gtrsim m_r^2$$

# Waves

$$\mathcal{L} = |\partial_\mu \phi|^2 - V_\phi \quad V_\phi = \frac{m_r^2}{2f_a^2} \left( |\phi|^2 - \frac{f_a^2}{2} \right)^2$$

Trivial inhomogeneous solutions:

$$\left\{ \begin{array}{l} r = 0 \text{ (radial mode on its VEV)} \\ \partial_\mu \partial^\mu a = 0 \end{array} \right.$$



**axion waves**

$$\phi = \frac{f_a}{\sqrt{2}} e^{ik(t-x)}$$

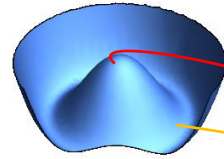


# Strings



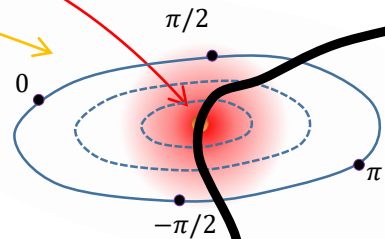
axion waves

$$\phi = \frac{f_a}{\sqrt{2}} e^{ik(t-x)}$$



string core

$$m_r^{-1} \sim f_a^{-1}$$



$$\phi = \frac{f_a}{\sqrt{2}} g(|x|) e^{i\theta}$$

$$\begin{cases} g(0) = 0 \\ g(\infty) = 1 \end{cases}$$

$$d \sim H^{-1}$$

**Nonlinear dynamics:**

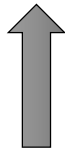
- ~~Analytical approach~~ 😞

**Large ratio of scales:**

- ~~Numerical approach~~ 😞

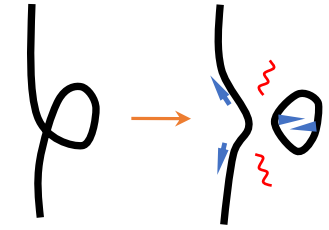
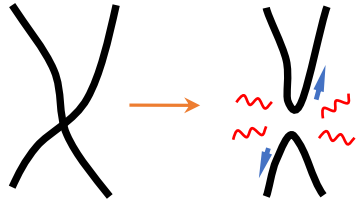
string tension

$$\mu = \frac{E}{L} \sim \underbrace{\pi f_a^2}_{\text{core}} \underbrace{\log \frac{d}{m_r^{-1}}}_{\text{axion gradient}} \sim \pi f_a^2 \log \frac{m_r}{H}$$

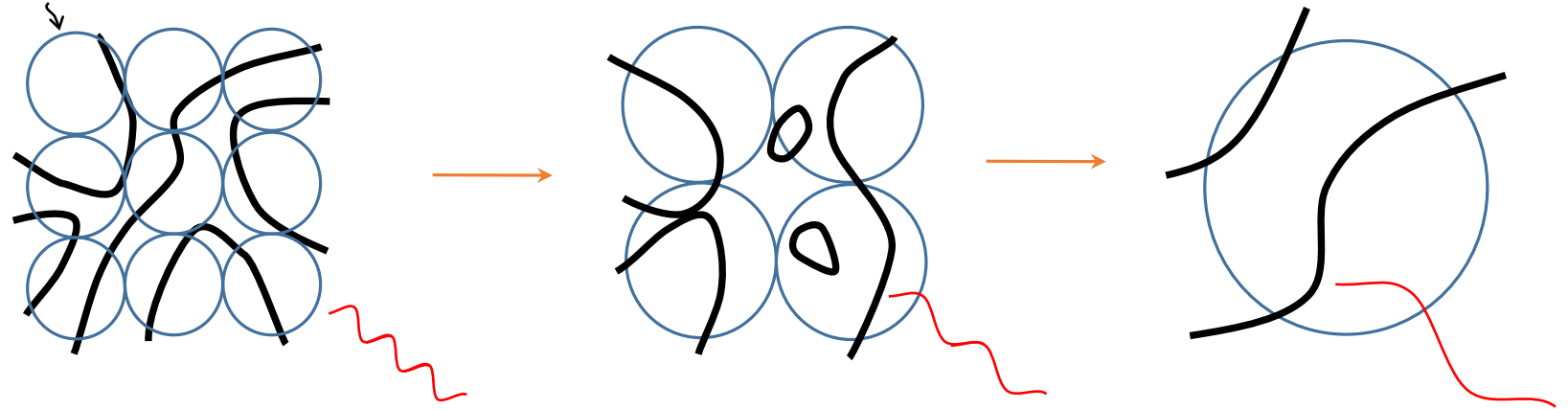


grows logarithmically in time

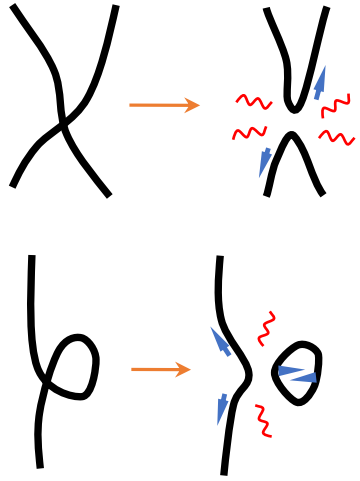
# The Scaling Regime



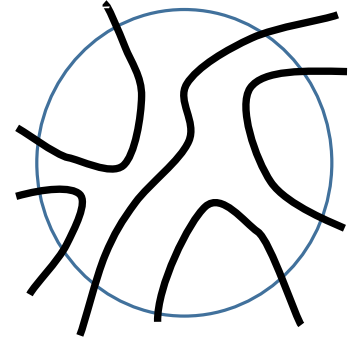
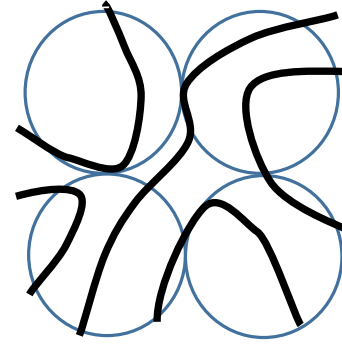
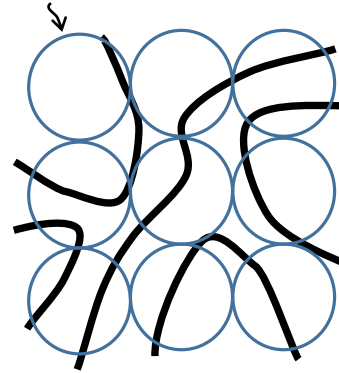
causal patch  $\propto 1/H = 2t$



# The Scaling Regime

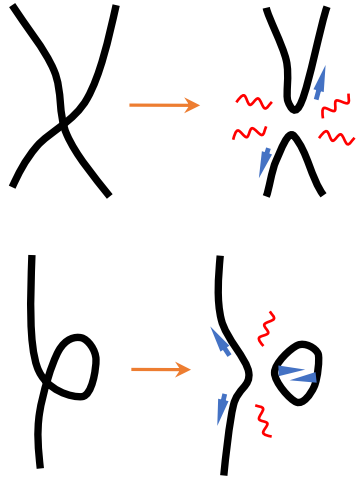


causal patch  $\propto 1/H = 2t$

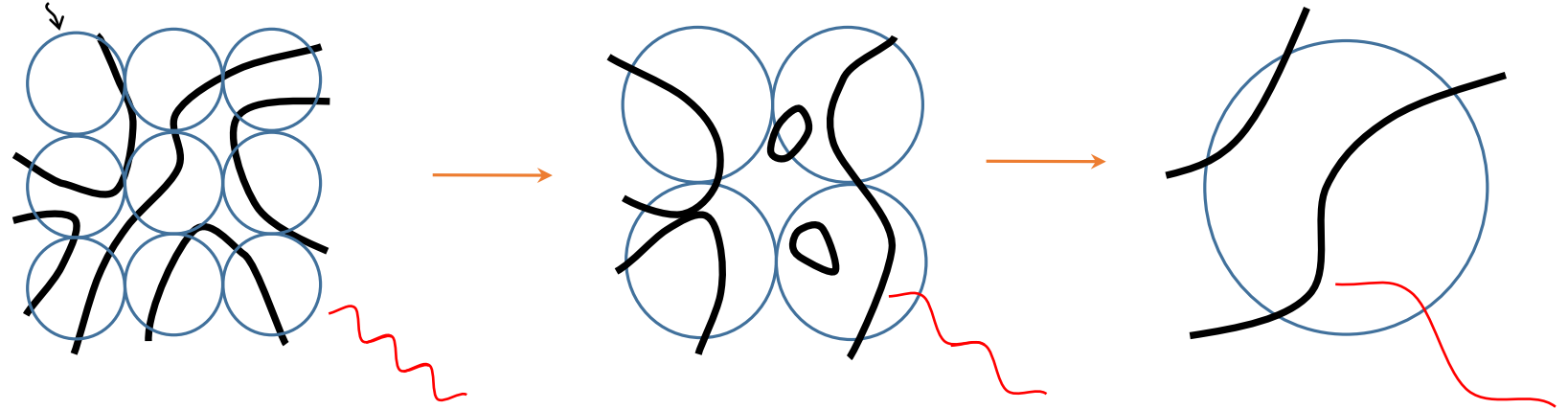


free strings:  $\rho^{\text{free}} \propto \frac{1}{R^2} \propto \frac{1}{t}$

# The Scaling Regime



causal patch  $\propto 1/H = 2t$



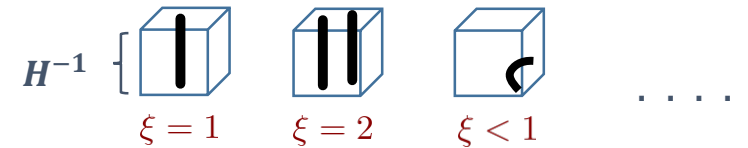
rate of energy loss:

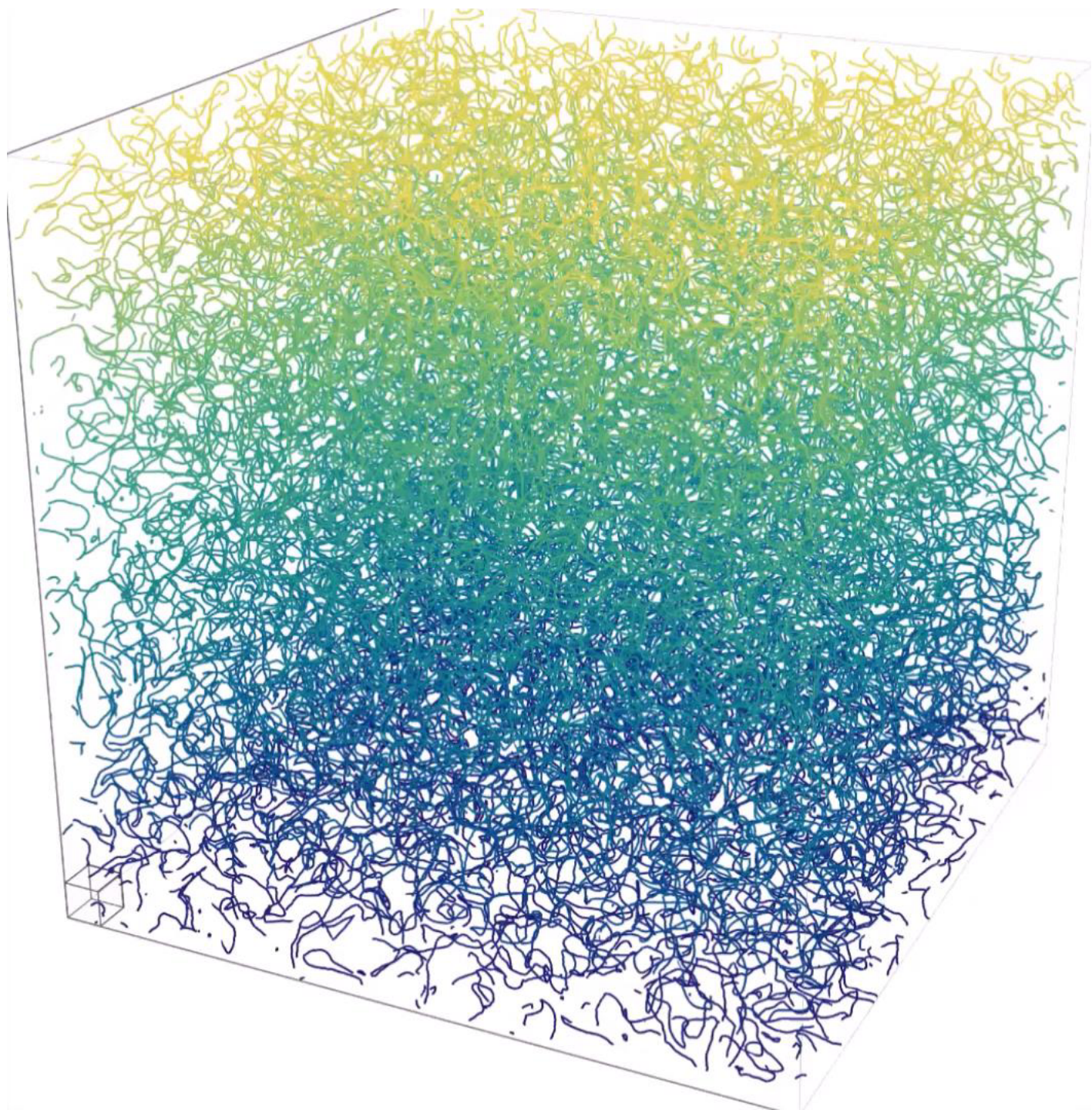
$$\Gamma \equiv \frac{d}{dt} [\rho^{\text{free}} - \rho^{\text{scal}}] \approx \frac{\xi \mu}{t^3}$$

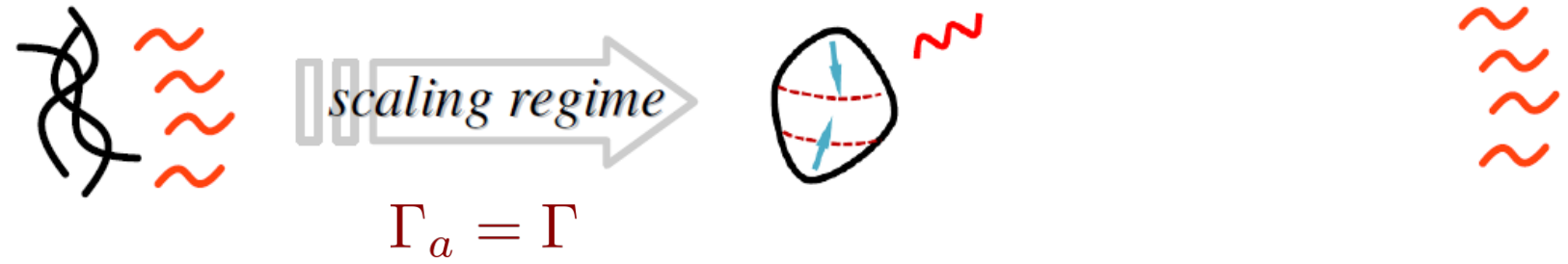
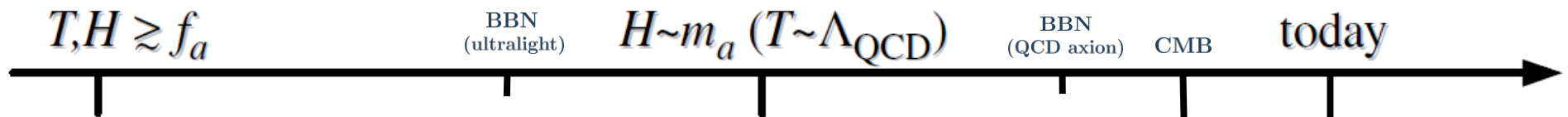
$$\propto \frac{1}{R^2} \propto \frac{1}{t}$$

$$\frac{\xi \mu}{t^2}$$

number of strings  
per Hubble patch







$$\Gamma_a = \Gamma$$

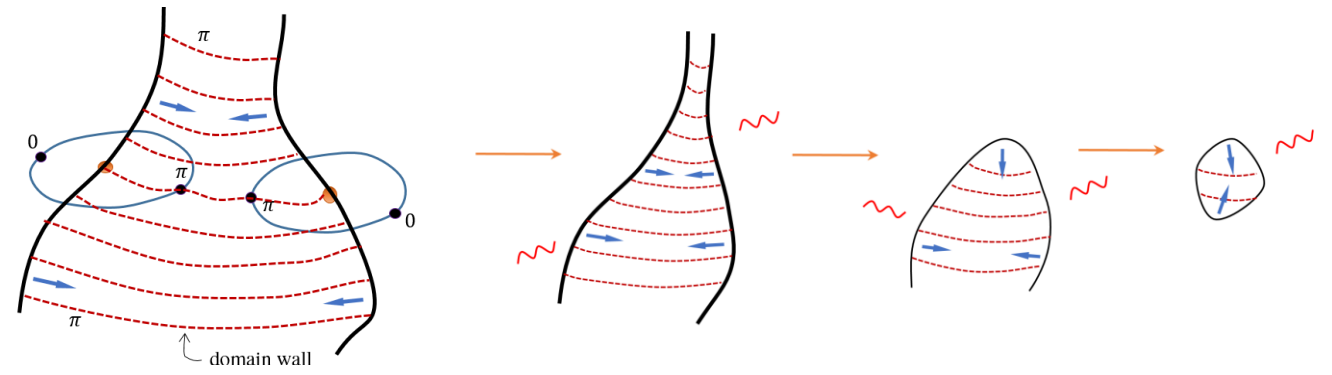
$$\log(m_r/H) \quad \sim 1 \div 15$$

$$\sim 70 \div 100$$

strings form

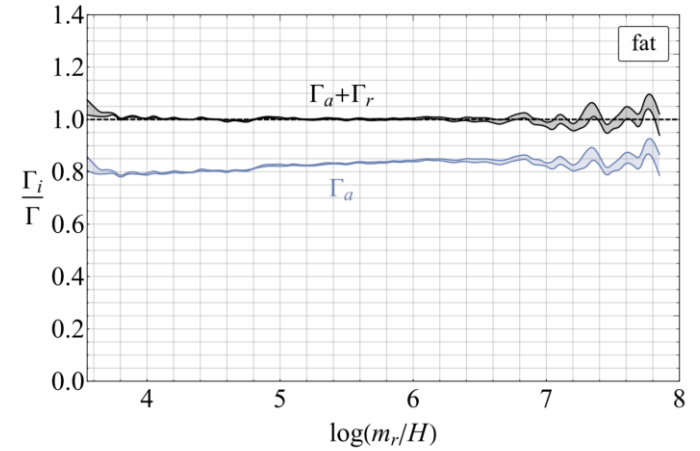
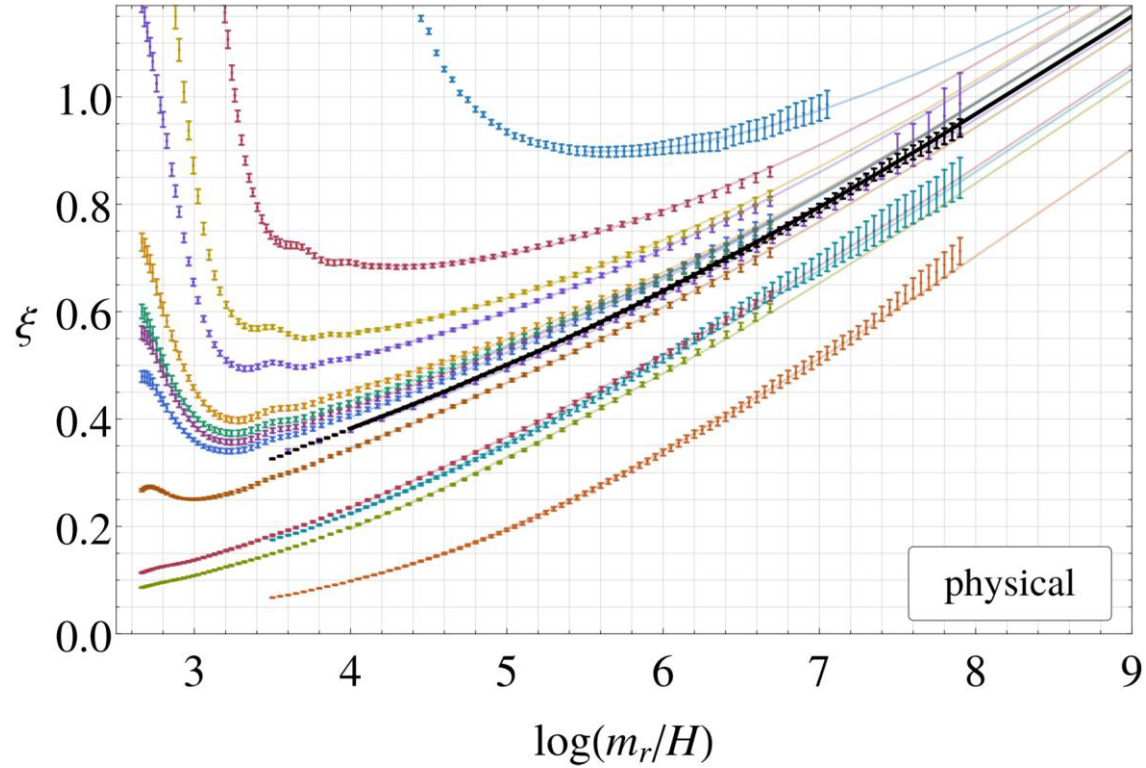
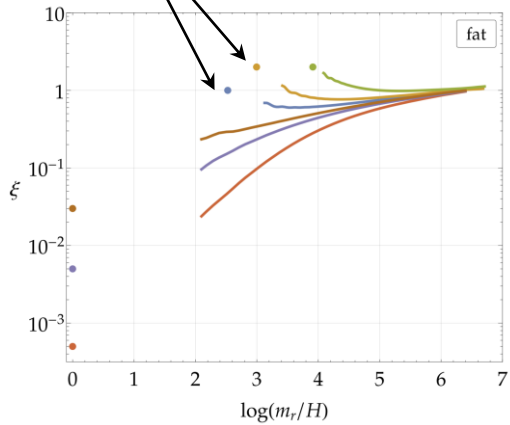
domain walls form and annihilate

relic axions and gravitational waves



# Attractor and Logarithmic Running

different initial conditions



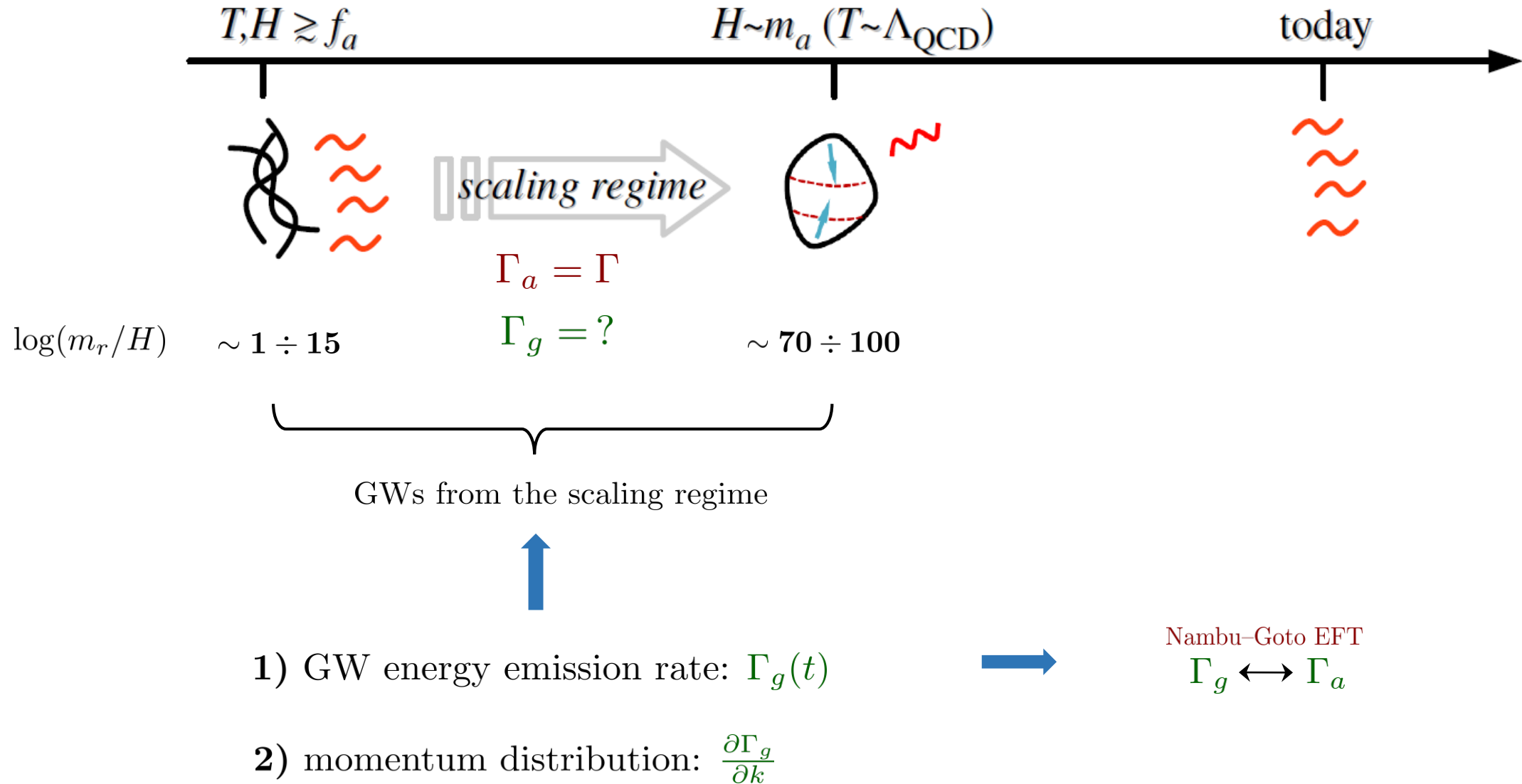
Scaling Violation

$$\xi = c_1 \log + c_0 + \frac{c_{-1}}{\log} + \frac{c_{-2}}{\log^2}$$



$$\Gamma_a = \frac{\xi \mu}{t^3} \propto \frac{f_a^2 \log^2}{t^3}$$

# Gravitational Waves



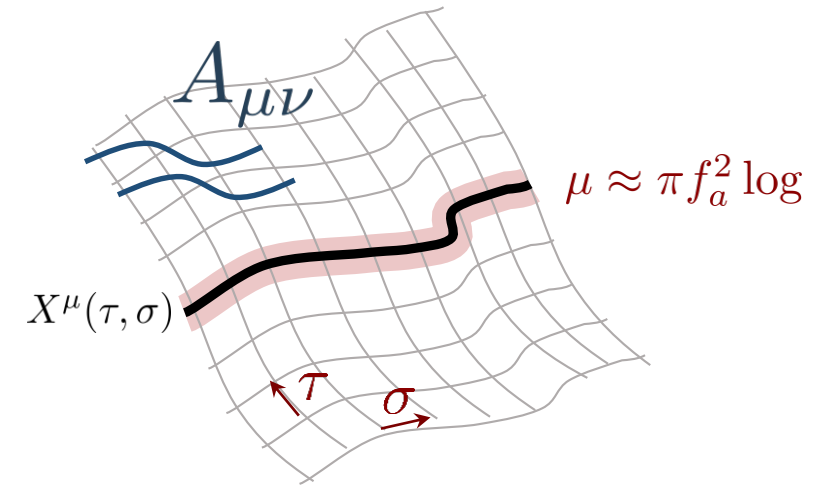


# String Effective Theory

- 1) GW energy emission rate:  $\Gamma_g(t)$
- 2) momentum distribution:  $\frac{\partial \Gamma_g}{\partial k}$

Degrees of freedom:

- $a \longleftrightarrow A_{\mu\nu}$
  - $X^\mu(\tau, \sigma)$
- $$\partial A \sim F^{\mu\nu\rho} = \epsilon^{\mu\nu\rho\sigma} \partial_\sigma a$$



$S_\phi[\phi]$

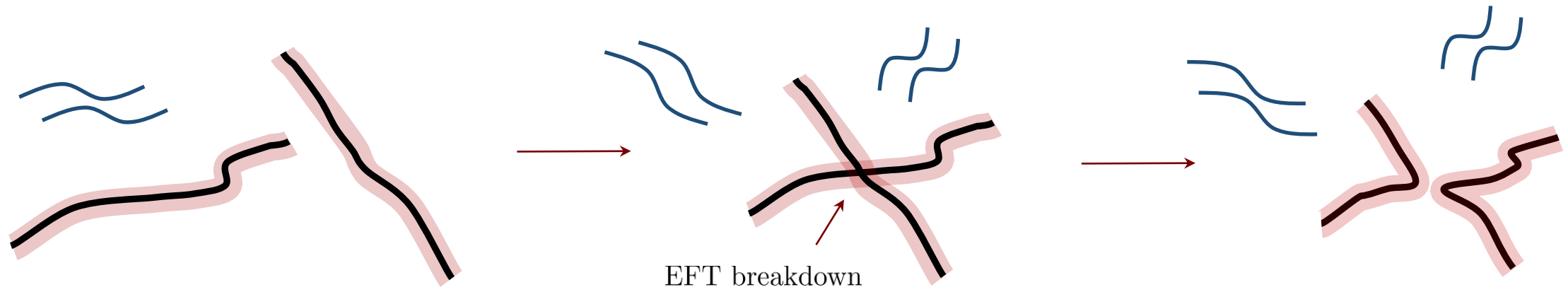


$$S_{\text{EFT}}[X, A] = \underbrace{-\mu \int d\tau d\sigma \sqrt{-\gamma}}_{\text{Nambu-Goto action}} \underbrace{-\frac{1}{6} \int d^4x (\partial A)^2}_{\text{Axion kinetic term}} + \underbrace{2\pi f_a \int d\tau d\sigma \partial_\tau X^\mu \partial_\sigma X^\nu A_{\mu\nu}}_{\text{Axion-string interaction (Kalb-Ramond action)}}$$

axion-to-string coupling

$\gamma_{ab} = \partial_a X^\mu \partial_b X_\mu$

- 1) GW energy emission rate:  $\Gamma_g(t)$
- 2) momentum distribution:  $\frac{\partial \Gamma_g}{\partial k}$



# Gravitational Wave Emission

- 1) GW energy emission rate:  $\Gamma_g(t)$
- 2) momentum distribution:  $\frac{\partial \Gamma_g}{\partial k}$

EoM for  $A^{\mu\nu}$ :  $\square_x A^{\mu\nu} = 2\pi \boxed{f_a} \int d\sigma \dot{X}^{[\mu} X'^{\nu]} \delta^3(\vec{x} - \vec{X})$

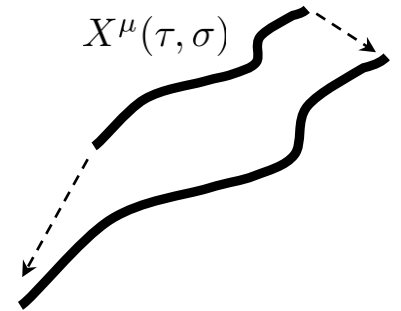
Einstein Eq:  $\square_x h^{\mu\nu} = 16\pi G (T_s^{\mu\nu} - \frac{1}{2}\eta^{\mu\nu} T_s^\lambda{}_\lambda)$

$$T_s^{\mu\nu} = \boxed{\mu} \int d\sigma (\dot{X}^\mu \dot{X}^\nu - X'^\mu X'^\nu) \delta^3(\vec{x} - \vec{X})$$

$$\frac{dE_a}{dt} = \underbrace{r_a[X]}_{\uparrow} \boxed{f_a^2}$$

$$\frac{dE_g}{dt} = \underbrace{r_g[X]}_{\uparrow} G \boxed{\mu^2}$$

dimensionless functionals of the shape of the string trajectory  $X^\mu$



$$\frac{\Gamma_g}{\Gamma_a} = \underbrace{\frac{r_g[X]}{r_a[X]}}_{\equiv r} \frac{G\mu^2}{f_a^2}$$

$$\equiv r = \text{const}$$

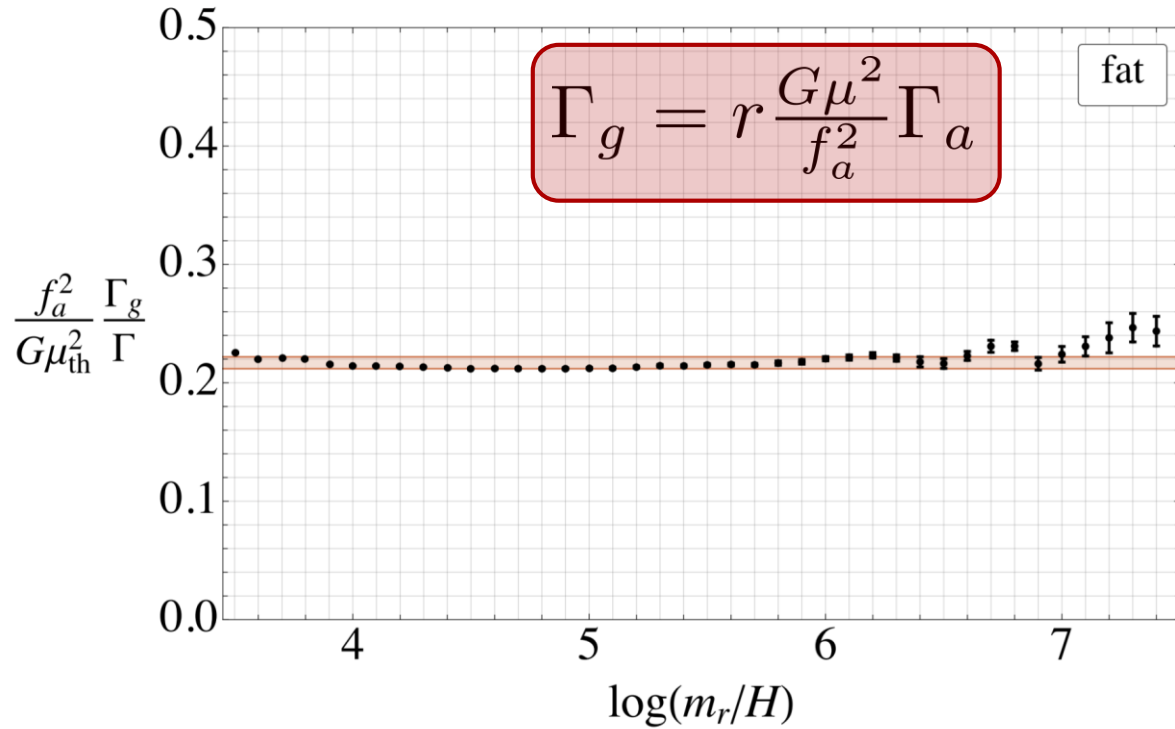
$$\Gamma_g = r \frac{G\mu^2}{f_a^2} \Gamma_a$$

$\frac{\xi_\mu}{t^3}$

$$\propto \frac{\log^4}{t^3}$$

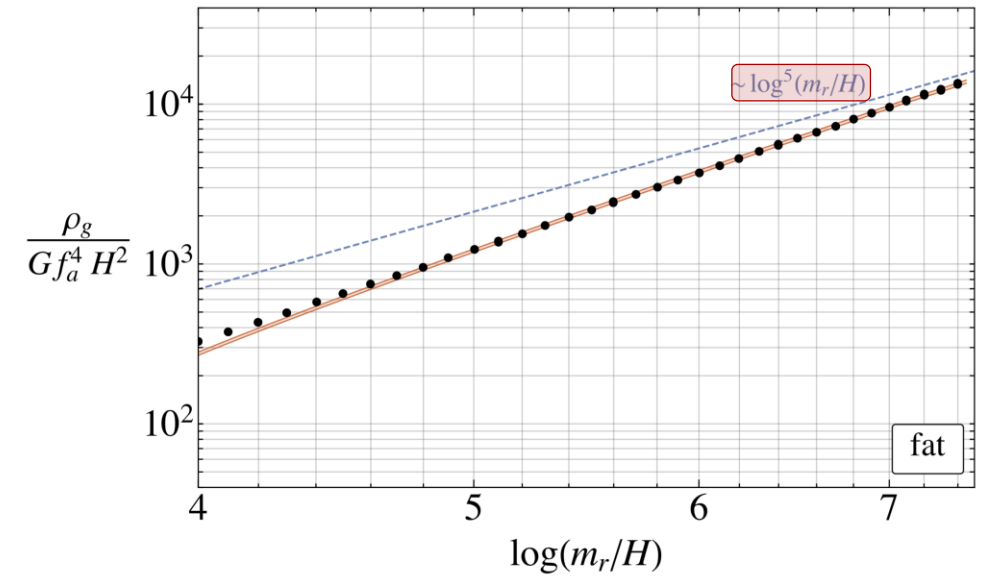
# Comparison with the Field Theory Evolution

- 1) GW energy emission rate:  $\Gamma_g(t)$
- 2) momentum distribution:  $\frac{\partial \Gamma_g}{\partial k}$



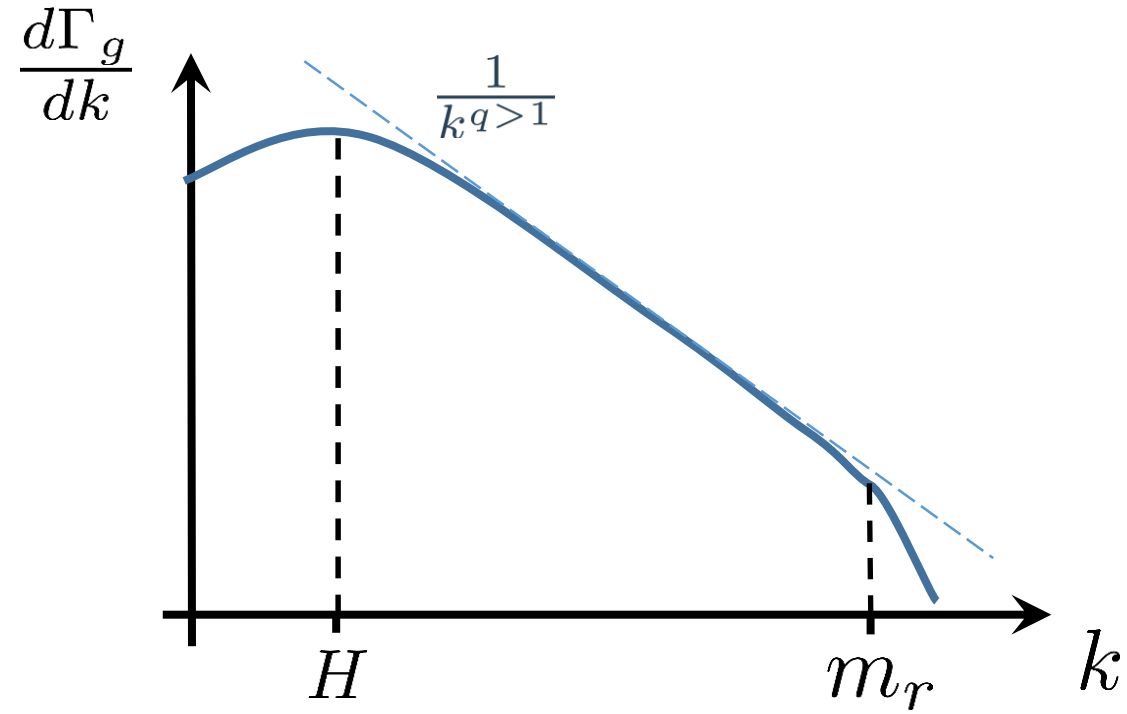
total GW energy:

$$\rho_g = \int dt' \left(\frac{R'}{R}\right)^4 \Gamma'_g \propto \frac{\log^5}{t^2}$$



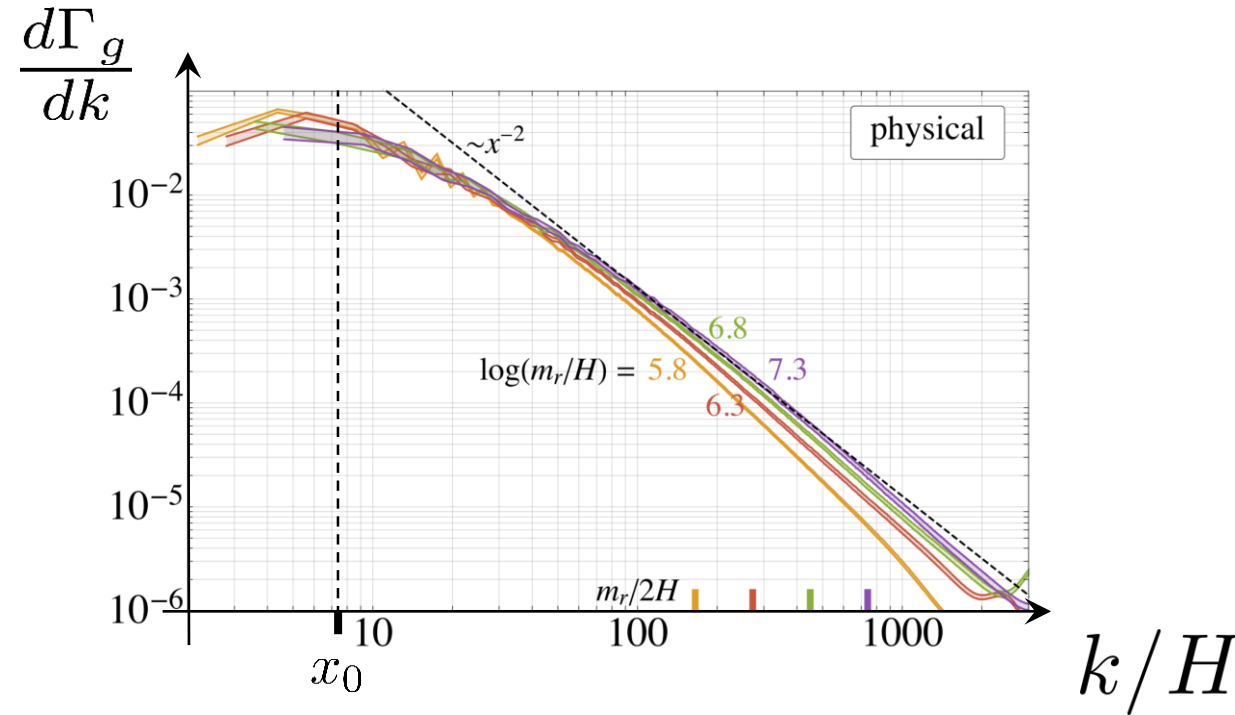
# The Gravitational Wave Spectrum

- 1) GW energy emission rate:  $\Gamma_g(t)$
- 2) momentum distribution:  $\frac{\partial \Gamma_g}{\partial k}$



# The Gravitational Wave Spectrum

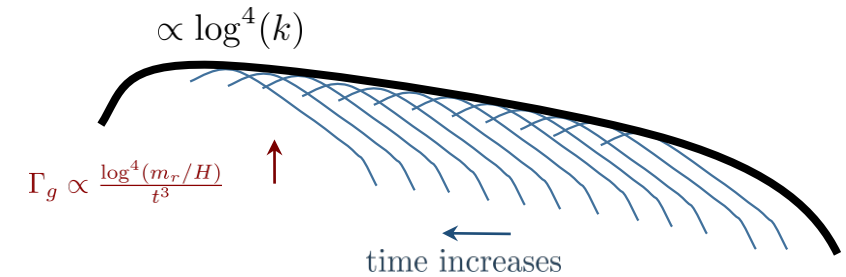
- 1) GW energy emission rate:  $\Gamma_g(t)$
- 2) momentum distribution:  $\frac{\partial \Gamma_g}{\partial k}$



$$\frac{\partial \rho_g}{\partial \log k} \equiv \int dt' \frac{d\Gamma'_g}{d \log k} \left(\frac{R'}{R}\right)^4 =$$

$$\simeq 8\pi^3 r G f_a^4 H^2 \log^4 \left[ \frac{m_r}{H} \left(\frac{x_0 H}{k}\right)^2 \right]$$

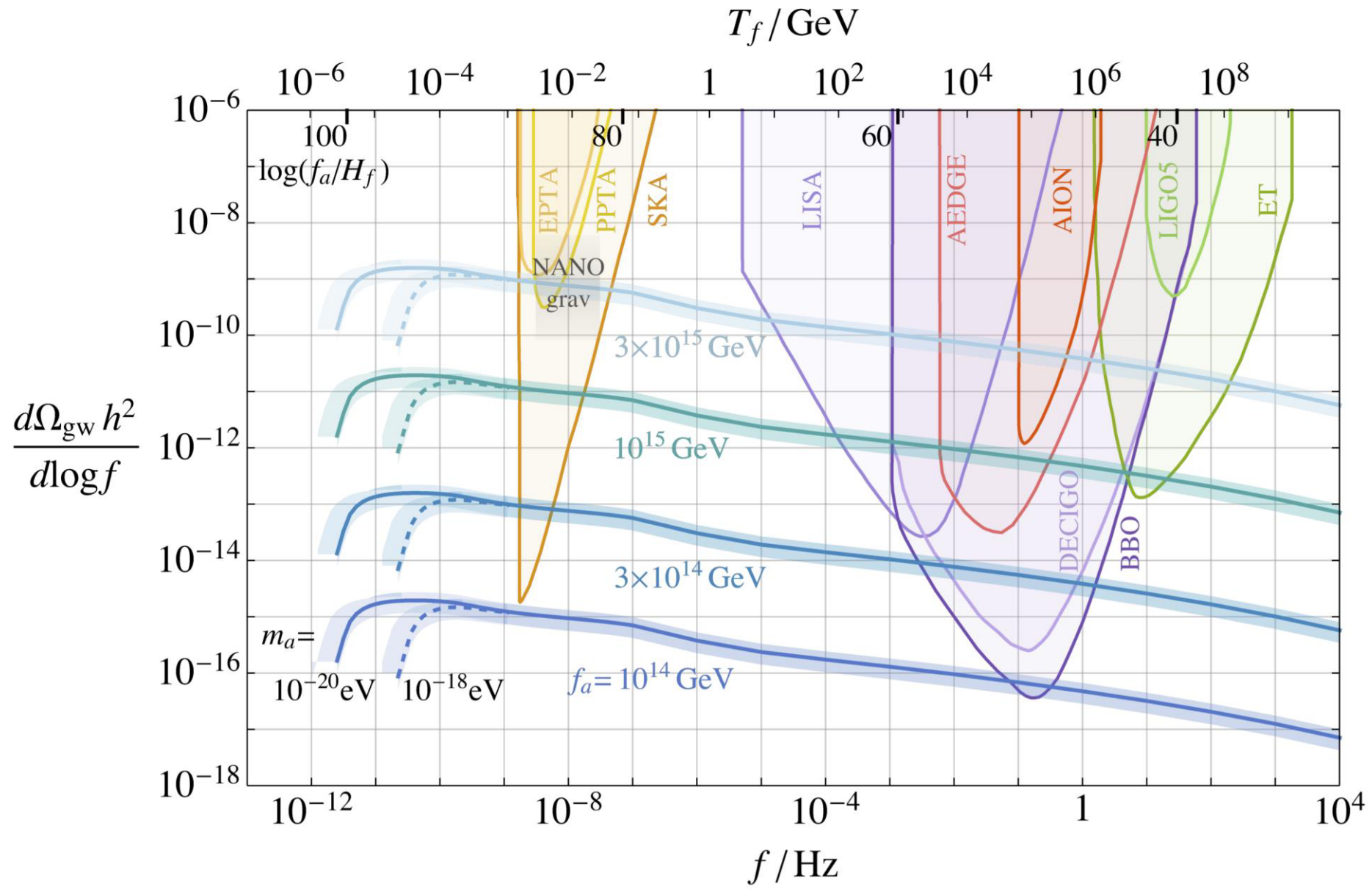
$\propto \frac{\log^4}{t^3}$



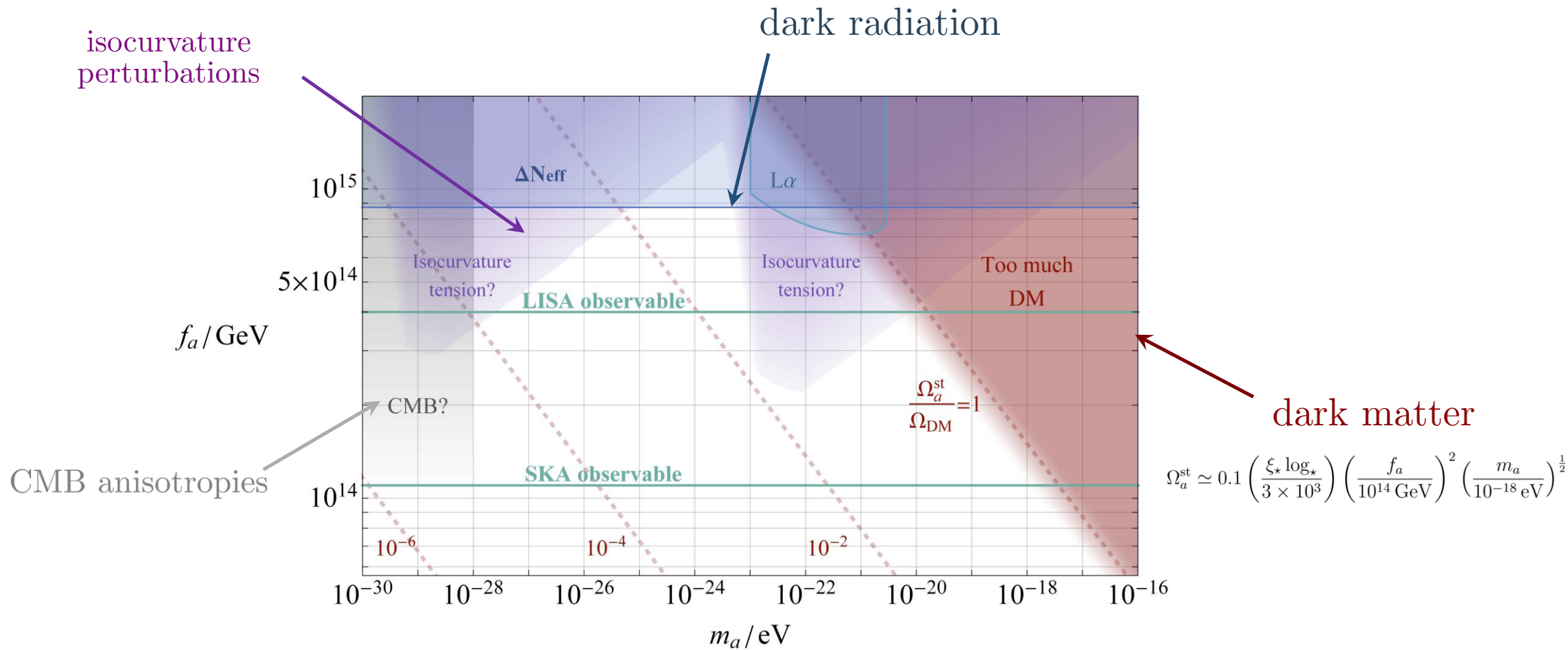
- approximately scale invariant
- $\log^4$  enhancement

$$\frac{d\Omega_{\text{gw}} h^2}{d\log f} \simeq 10^{-15} \left(\frac{r}{0.26}\right) \left(\frac{f_a}{10^{14}\text{GeV}}\right)^4 \left(\frac{10}{g_f}\right)^{\frac{1}{3}} \left\{ 1 + 0.12 \log \left[ \left(\frac{m_r}{10^{14}\text{GeV}}\right) \left(\frac{10^{-8}\text{Hz}}{f}\right)^2 \right] \right\}^4$$

$$\begin{cases} f_a \lesssim 10^{15} \text{ GeV} \\ m_a \lesssim 10^{-18} \text{ eV} \end{cases}$$



# Bounds on the Post-Inflationary Scenario





# Conclusions

- **Axions are motivated BSM candidates**

→ in the post-inflationary scenario, the cosmological evolution is governed by cosmic strings

- **The scaling regime produces an approximately scale invariant GW spectrum**

→  $\Gamma_g \propto \log^4$  (from the increase in  $\mu \propto \log$  and  $\xi \propto \log$ ) leads to logarithmic violations of scale invariance

→ enhances the spectrum at low frequencies

- **The spectrum is visible by multiple experiments for  $f_a > 10^{14}$  GeV**

→ best prospects in PTAs and LISA

- **Constraints on the post-inflationary scenario**

→  $f_a \lesssim 10^{15}$  GeV and  $m_a = 10^{-28} \div 10^{-18}$  eV is viable

## Outlook

- Local strings?
- (Initial conditions for the subsequent evolution?)